

A key step transforming gaz and dust into stellar systems



PROTOSTARS & PROTOSTELLAR DISKS



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From starless cores to protostars: triggering of gravitational instability

Population studies & typical timescales

Transporting material from large to small scales: infall, accretion & the luminosity problem

Observations: statistical properties of protostellar disks: sizes (& masses)

I. DISKS and dust

What sets the properties of the most pristine disks?

Cabinet de curiosités: structures and sneak peaks at individual objects

What do we know of dust evolution while the star-disk system is built ?

Which instruments for the future (if time allows ...)?





« And yet, it rotates » angular momentum problem & the physics of the formation of disks













CALYPSO (Maury+ 2011, 2014, 2019): <25 % of Rdisk > 60 au

CONSTELLATION DE PERSÉE NGC1333

HERSCHEL Gould Belt Survey

Two complete & recent reviews :

https://ui.adsabs.harvard.edu/abs/2022FrASS...9.9223M/abstract https://arxiv.org/pdf/2209.13765.pdf









The envelope masses decreases with increasing T_{bol} (evolution): BUT dust disk radii seem insensitive dust drift drives dust disk sizes ? => kinematics needed



Sheehan+ (2022) from VANDAM survey



Aso et al. (2021) Mix of infall and Keplerian motions

150 au disk in Class ITMC1-A



Gas in Keplerian motions observed in the largest disks (spatially resolved) - few of them => disk radii mostly consistent with dusty radii in the current small samples



Class 0/I L1527 disk resolved from gas kinematics obtained with ALMA

- Keplerian radius 74 au

Don't count your chickens before they are hatched....



=> Measure the angular momentum due to rotational motions inside protostellar envelopes, at smaller radii than values found in the literature



Another possible scenario is that AM in envelopes is not transmitted down to disk scales if envelopes are subject to MAGNETIC BRAKING

Magnetic fields can redistribute angular momentum away from a collapsing region thanks to torsional Alfvèn waves (Mouschovias & Paleologou 1979, 1980)

Timescale for magnetic braking: $t_b = \rho R(2\rho_0 \times V_A)$



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Hennebelle+ 2016 : Analytical non-ideal MHD collapse leads to self-regulation of disks to small (20-50 au radii)

Typical radius is:

$$r_{
m d,AD} \simeq 18 \, {
m AU} \, imes$$

$$\delta^{2/9} \left(\frac{\eta_{\rm AD}}{0.1\,{
m s}}\right)^{2/9} \quad \left(\frac{B_z}{0.1\,{
m G}}\right)^{-4/9} \left(\frac{M_{\rm d}}{0.1\,{
m M}}\right)^{-4/9}$$

By contrast hydro would lead to:

$$r_{\rm d,hydro} \simeq 106 \, {
m AU} \, rac{eta}{0.02} \left(rac{M}{0.1 \, {
m M}_\odot}
ight)$$

=> Early disk formation is magnetically *self-regulated* !



Another possible scenario is that AM in envelopes is not transmitted down to disk scales if envelopes are subject to MAGNETIC BRAKING



See our review in Maury, Hennebelle, Girart 2022

Non-ideal MHD Class 0 disks: sizes and shapes under various local conditions



Masson+ 2015

Don't count your chickens before they are hatched....

=> Measure the magnetic fields in protostellar envelopes, and check predictions of magnetized models



If you want to know more on how to measure B in star-forming environments:

attend Thursday evening short talks !

Dust polarized emission is widely used to trace the B fields in ISM

gneticy sos





dit: Detlev Raveenswaay / Science Photo Library







B-fields are detected in ALL protostars where they are looked for + remain largely organized in low-mass cores (Less in massive cores ? Lack of statistics)

Sadavoy et al. (2019)

Many others in the literature !

Is magnetic field ubiquitous ?

Galametz, Maury + (2018, 2020): observations of polarized dust emission from an unprecedentedly large sample of Class 0 protostars <u>B fields detected in all of them at scales 2000 au !</u>



+ Keep tuned for results from the ALMA BOPS survey (Ian Stephens)

Can magnetic braking be inferred from gas kinematics ?



 µ: ratio of magnetic to gravitational energy (normalized by the critical value)



Protostars with aligned B configurations show : less kinetic energy from rotation in their inner envelopes + less multiple systems

=> an expected outcome of more efficient magnetic braking ? Still a VERY disputed result !





B335 seen by JWST (NIRSpec)









Comparing data to physical MHD models of protostellar disk formation

Numerical simulations performed with the RAMSES code (Teyssier 2002, Fromang+2006)

- Adaptive Mesh refinement (AMR) \bullet
- Non-ideal MHD





Dust assumptions:

- $a_{\rm min}$ = 5 nm to $a_{\rm max}$ = 1 to 50 μ m
- MRN distribution: $dn(a) \propto a^{-3.5} da$ \bullet
- Composition:

Alignment mechanisms:

- Perfect alignment approach (not a physical mechanism!!!)
- Radiative Torques (RAT) (+ Imperfect Internal alignment+ Imperfect Davis-Greenstein) \bullet

3D Monte-Carlo continuum radiative transfer code

- Simulate **intensity** and **polarization** of light
- Dust heating & dust grain alignment

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62.5% astrosilicates + 37.5% carbonaceous grains
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Comparing data to physical MHD models of protostellar disk formation



Our ALMA map of the polarized dust emission at 230 GHz



=> B regulates the early properties of the protostellar disk in B335

A magnetically-regulated disk formation in B335 ?

RAMSES MHD models parameter space: Core: 2.5 Msun Times: 0.07, 0.14 and 0.2 Myrs Mass-to-flux ratio mu : 3, 5, 6, 10 Rotational energy beta 0.1% 1% 10% Turbulent energy: Mach 0.01 0.2 0.5 1.0

MHD model: synthetic N_H & B-field

Only models with dynamically relevant B-field match the data (best model μ ~6)



Coupling of B to star/disk forming material

First measurement of the ionization fraction of the gaz at disk-forming scales Cabedo et al. (2022)

models



Observations consistent with production of cosmic rays LOCALLY => at disk-forming scales (r<500 au), almost quasi ideal MHD conditions?



Ionization from protostellar CRs ?

First measurement of the ionization fraction of the gaz at disk-forming scales

Cabedo, Maury et al. (2022)



Observations consistent with production of cosmic rays LOCALLY => at disk-forming scales (r<500 au), almost quasi ideal MHD conditions?



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Disks in Class I protostars: some show structures



Yamato et al. (2023)





HH212: Class 0 protostar

Gas disk radius : 45 au

M∗+M_{disk} ~ 0.25 M_☉

observations of MHD disk-wind ?



MHD DW dominates the AM extraction out to 40 au, but not inside <4 au (MRI?)

Lee+ 2021











Class I protostar R CrA IRS7B-a

Disk gas mass : 0.41 M_☉ Core mass : I.9 M_{\odot}



Ohashi+ 2023: Keplerian rotation in the disk => central protostellar mass of 2–3 M_{\odot}

Dust disk radius : 62 au Gas disk radius : 80 au

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Challenging the current paradigm leading for disk evolution:

disks observed with ALMA (DSHARP survey, dust seen at 1mm) Andrews+ 2020 T-Tauri



2-Myr old disks : less mass (dust and gas) than required for minimum mass of the solar nebula ? (Cieza+ 2015, Miotello+ 2016, Sturm+ 2022, Manara+ 2022) - highly debated in community

recent observations of 1-5 Myrs T-Tauri disks: more structured & less massive than expected / required ?



SEE TESTI'S LECTURE





Why is dust important beyond planet formation ?

Dust: a key parameter of magnetic field coupling to the star-disk forming material !

See Guillet+ 2020, Zhao+ 2020, Lebreuilly+ 2024



Dust grain sizes affects the MHD resistivities and the role of the magnetic field



SEE LEBREUILLY'S LECTURE

$${
m M}_{
m obj} \propto -$$

Masses of astrophysical structures measured from dust emission depends much on the dust optical constants which vary with dust composition, size etc



Dust is an indirect tracer of mass Used from Galaxies, to clouds, cores, disks etc

Mass deduced from mm observations if optically thin :

Flux $(abs(\nu) B(\nu, T_{obi}))$

SEE TESTI'S & DEMYK'S LECTURE

THERMAL EMISSION

- function constituted by the Planck blackbody B_{V} at temperature T is: $I_{\nu} = B_{\nu}(T)(1 - e^{-\tau_{\nu}}) \quad \text{where} \quad B_{\nu}(T) = \frac{2h\nu^3}{c^2} \frac{1}{e^{h\nu/k_BT} - 1}$
- \Rightarrow Assuming Iv being uniform over the solid angle Ω , the corresponding flux is:

$$F_{\nu} = \Omega(1 -$$

a power law: $\tau_{\nu} = \left(\frac{\nu}{\nu_0}\right)^{\beta}$ where the cut-off frequency $\nu_0 = c/\lambda_0$ is such that $\tau_{\nu_0} = 1$

→ In the limit of v<< vo:
$$(1 - e^{-\tau_{\nu}}) = \tau_{\nu}$$

Observations at three or more wavelengths this simple model can be used to determine both **T** and β

 \Rightarrow Solution of the radiative transfer equation for a medium with optical depth τv and for a source

 $-e^{-\tau_{\nu}}B_{\nu}(T)$

 \Rightarrow Empirical behaviour of τv_{ρ} as a function of v for large interstellar dust grains is modelled as

$$I_{\nu} \approx \left(\frac{\nu}{\nu_{0}}\right)^{\beta} B_{\nu}(T)$$

Measuring dust properties: spectral dependance of dust emissivity from thermal emission

THERMAL EMISSION



Wavelength, µm

Measuring dust properties: spectral dependance of dust emissivity from thermal emission



Dust temperature and spectral index in RJ regime are usually obtained from SEDs of dust thermal emission from 70µm to 3mm

We can measure the spectral index of the SED as α

$$\alpha_{1-2mm} = \frac{\log(F(v_1)/F(v_2))}{\log(v_1/v_2)},$$

Assuming optically thin emission at mm wavelengths, this informs on β as being α -2 in RJ approximation

$$\kappa_{1.2\mathrm{mm}} = \kappa_{250\mu\mathrm{m}} \left(\frac{\nu_{1.2\mathrm{mm}}}{\nu_{250\mu\mathrm{m}}}\right)^{\beta}$$



Protostellar envelopes are **COLD**: the **Rayleigh-Jeans approximation does not hold** because of their low dust temperatures

=> need to know T(dust) to get the Planck law correctly

$$\beta = \frac{\log \frac{S_{1.2\text{mm}}}{S_{2\text{mm}}} - \log \frac{B_{1.2\text{mm}}(T_{\text{d}})}{B_{2\text{mm}}(T_{\text{d}})}}{\log \frac{\nu_{1.2\text{mm}}}{\nu_{2\text{mm}}}}.$$



How do we measure dust properties ?

Dust absorption coefficient (also called dust opacity, and dust emissivity):



wavelength, cm

SEE DEMYK'S LECTURE

Rayleigh regime (a $<< \lambda$):

 $Q_{abs} \sim a => \sigma Q_{abs} \sim a^3 \sim grain$ (total) volume

= when grain size is $<< \lambda$, impossible to derive a from the dust SED of grains emitting at thermal equilibrium

> In the Mie regime (a ~ λ): Q_{abs} strong function of a/λ , so β is affected

> > =>

The dust emissivity spectral index ONLY gives indication on Q_{abs} cross sections when a ~ λ

It can be measured in absorption, or in emission





DENSE ISM

$R_V = A_V / E_{B-V}$ the color excess

A measure of wavelength dependence of the interstellar extinction



- Extinction = Absorption + Scattering
- Scattering efficiency is a strong function of a/λ , and peaks when $2\pi a \sim \lambda$.
- A dust model is needed: $a_{max} \simeq 0.3 0.5 \ \mu m$ for the diffuse ISM (Rv = 3.1)
- The polarization curve in extinction can also be used.

Kandori et al. (2003) in L1251

the non-star-forming part of the cloud has a typical interstellar value $R_V=3.1$,

while in the denser part of the cloud experiencing star formation has $R_V = 6.5$



(b)

DENSE ISM

$R_V = A_V / E_{B-V}$ color excess changes in regions experiencing star formation



The abundance of small grains is very different !!!

Important for disk formation ! What do you start with ?

- MRN from diffuse ISM ($R_V=3.1$)?
- Already evolved ($R_V = 5.3$)?

PRESTELLAR CORES



See K. Demyk's lecture for more clues on dust evolution in clouds/cores with **coreshine** effect and **scattered light excess** in JWST spectra !

Coreshine: grains in low-mass cores scatter mid-infrared (3.6–4.5 µm) light efficiently

Pagani et al. (2010b) reported on the detection of **coreshine** for about half of 110 investigated prestellar cores

- To fit simultaneously the near-IR extinction and FIR Herschel emission profiles: the grain opacity has to increase
- Requires a = 0.5 1.5 µm (Steinaecker+2015) with astrosilicates.
- Requires a ~ 0.7 µm and fluffy grains (Ysard+2015, THEMIS) through the accretion of aliphatic-rich C-H mantles.



PRESTELLAR CORES

 \Rightarrow Solution of the radiative transfer equation for a medium with optical depth au_V and for a source

function constituted by the Planck blackbody Bv at temperature T is:

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⇒ In the limit of v<< v₀:
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 $I_{\nu} \approx \left(\frac{\nu}{\nu_0}\right)^{\rho} B_{\nu}(T)$

Observations at three or more wavelengths this simple model can be used to determine both **T** and **B**

Using longer wavelengths allows to probe for larger grains

$$\kappa_{1.2mm} = \kappa_{250\mu m} \left(\frac{\nu_{1.2mm}}{\nu_{250\mu m}}\right)^{\beta}.$$
$$\beta = \frac{\log \frac{S_{1.2mm}}{S_{2mm}} - \log \frac{B_{1.2mm}(T_{d})}{B_{2mm}(T_{d})}}{\log \frac{\nu_{1.2mm}}{\nu_{2mm}}}$$

No indication of change in the mm dust opacities of pre stellar cores ? => no grains larger than > 20 μ m



PROTOSTELLAR ENVIRONMENTS

Protostars embedded in the Perseus L1448 clump: observations with IRAM/NIKA2





Observations show flattening of **α** when probing larger column densities

NGC1333 star-forming region in Perseus cloud



β map: ranges from 1 to 3

- ➡ Low values tend to correlate with T_d peaks
- $\Rightarrow \beta$ values can be intrinsically

temperature dependent (Boudet et al. 2005)

ightarrow Grain growth can cause $oldsymbol{eta}$

values to decrease significantly



PROTOSTARS

L1527 IRS in the sky





L1527 IRS in the uv-space



Envelope-only $\alpha(r)$, $\beta(r)$



PROTOSTARS



PROTOSTARS



Millimeter spectral index of dust emissivity: Synthetic observations at envelope scales (in function of max temperature on the line-of-sight)

Using standard dust models:

Only models with large grains can reproduce observed spectral indices





What should we do to dust to fit observations?



Emissivity of irregular shaped grains can increase up to a few at mm (Lefevre et al. 2019, Min et al. 2016, Ysard et al. 2023) But not down to $\beta < 1.0$?



What should we do to dust to fit observations?





TO KNOW MORE: ATTEND K. DEMYK'S LECTURE !