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Stellar feedback

Image credit: NASA / ESA / CSA / Judy Schmidt

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What is stellar feedback?

- **Stellar feedback = injection of** *matter*, *momentum* and *energy* by stars into the circumstellar medium

- **Stellar feedback processes:**

- Photons & ionising radiation
- Winds
- Accretion, jets & outbursts (proto-/pre-MS stars)
- Supernovae
- Cosmic rays
- Binary stars
- X-ray binaries
- Runaway stars
- ...

Suggested literature

- [Hartquist, Dyson, Ruffle 2004](#), *Blowing Bubbles In The Cosmos: Astronomical Winds, Jets, and Explosions* (book)
- [Krumholz et al. 2014](#), *Star Cluster Formation and Feedback* (review; Protostars and Planets VI)
- [Dale 2015](#), *The modelling of feedback in star formation simulations* (review)
- [Zhang 2018](#), *A Review of the Theory of Galactic Winds Driven by Stellar Feedback* (review)
- [Hodges-Kluck et al. 2019](#), *Astro2020 Science White Paper: Hot Drivers of Stellar Feedback from 10 to 10,000 pc*
- [Geen et al. 2023](#), *Bringing Stellar Evolution and Feedback Together: Summary from a Lorentz Center Workshop*

Why does stellar feedback matter?

Stellar feedback...

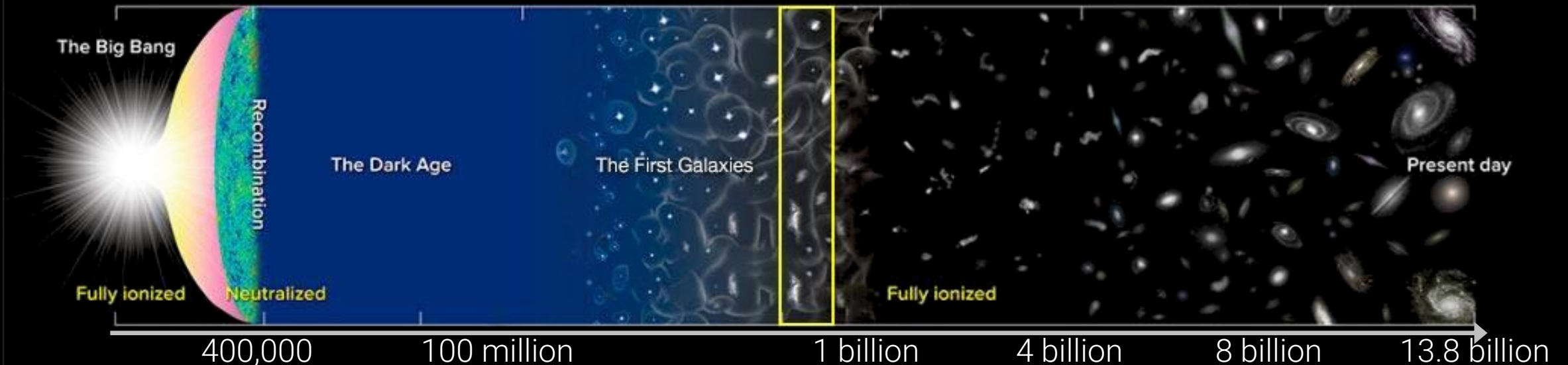
- ...enriches Universe with chemical elements



Why does stellar feedback matter?

Stellar feedback...

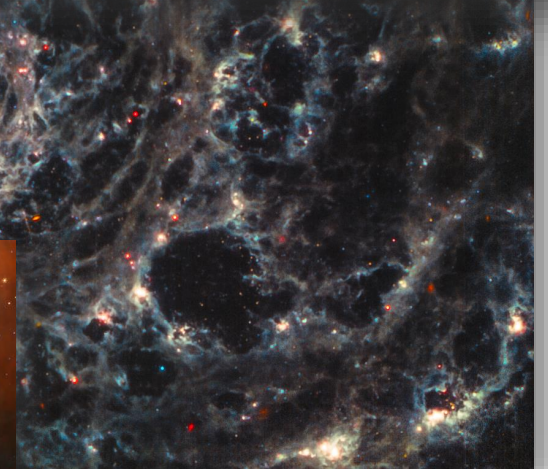
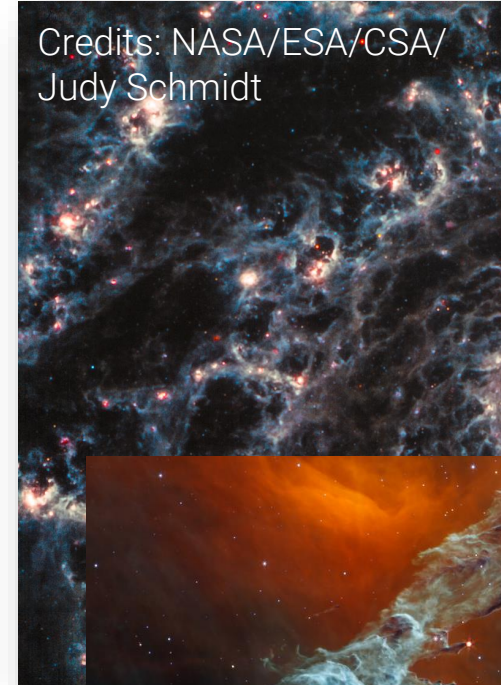
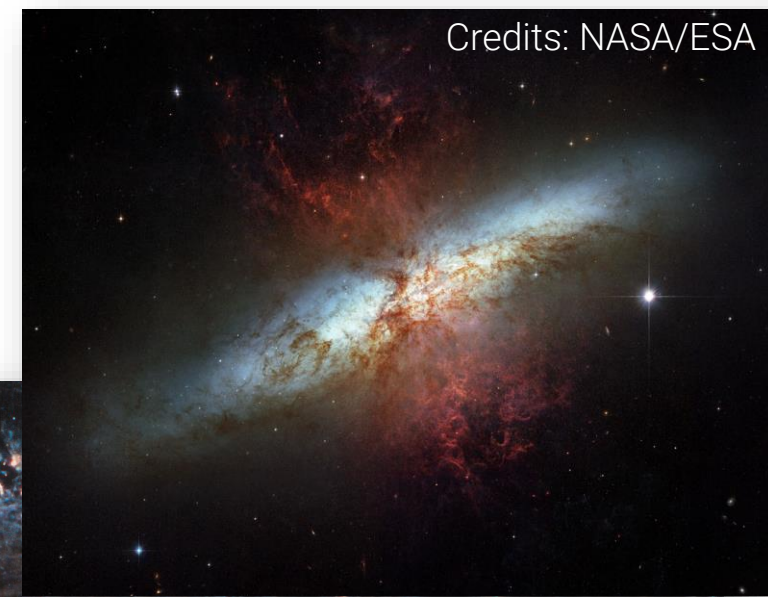
- ...enriches Universe with chemical elements
- ...helped to reionise Universe after dark ages



Why does stellar feedback matter?

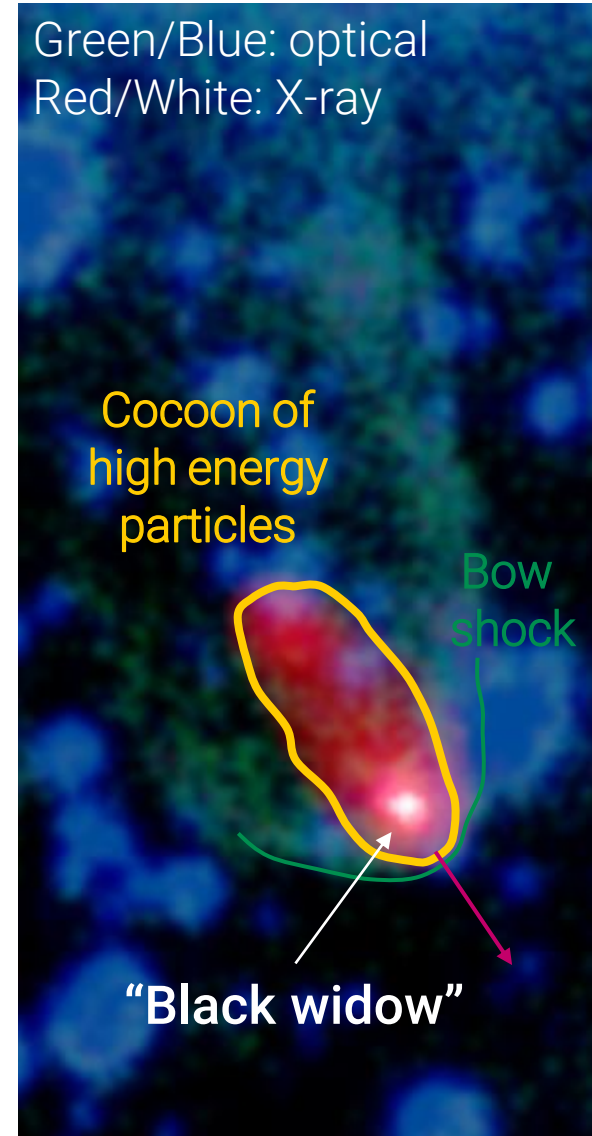
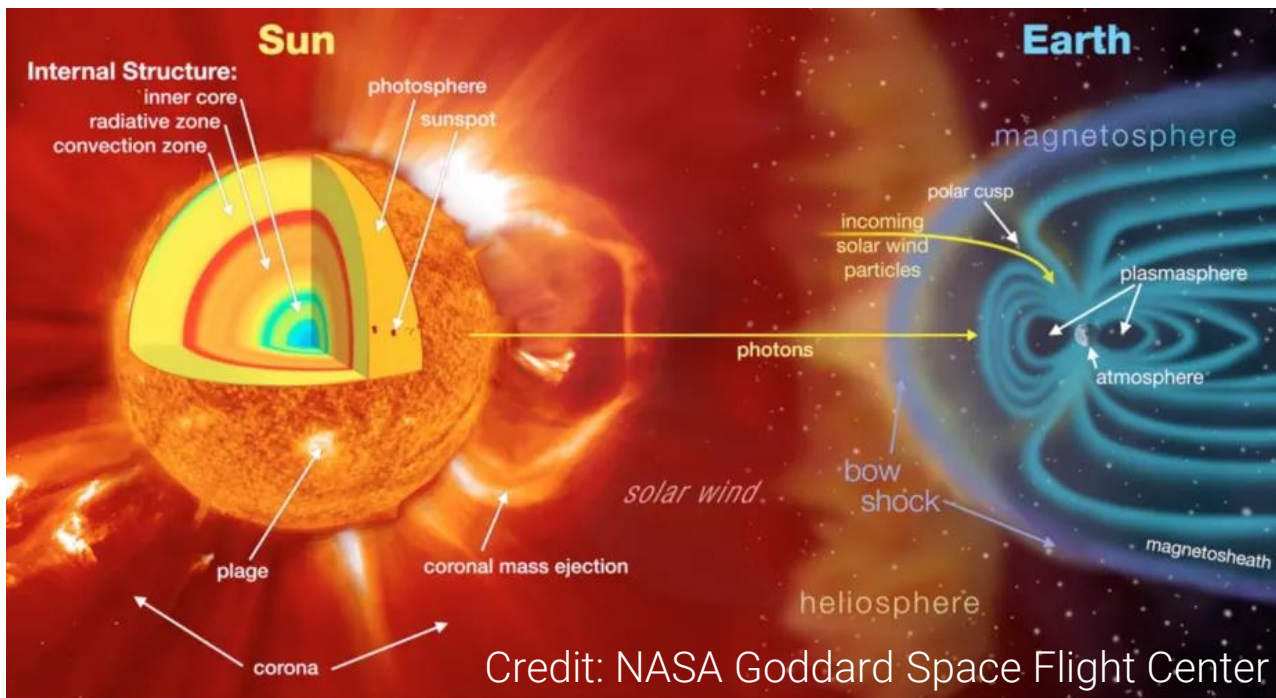
Stellar feedback...

- ...enriches Universe with chemical elements
- ...helped to reionise Universe after dark ages
- ...dictates appearance & evolution of galaxies
- ...regulates star formation (inhibits & triggers SF)
→ (in-)efficiency of star formation
- ...drives galactic winds and affects intergalactic medium (chemical enrichment, temperature, ...)
- ...sets life-cycle of giant molecular clouds
- ...is necessary for forming realistic galaxies in simulations (e.g. properties at $z=0$)
→ also limits formation of dwarf galaxies
- ...



Disclaimer

- **Stellar feedback relevant on many scales**
 - depends on application
 - not all scales covered in lecture (e.g. planetary scale)
- **Stellar feedback intrinsically linked to environment** – feedback from stars in vacuum does not do much...
 - effective stellar feedback



A vibrant nebula with swirling blue and orange gas clouds, set against a dark background filled with numerous stars. The nebula's structure is complex, with filaments and clumps of gas. The blue regions are more prominent in the center, while orange and yellow hues are seen towards the edges and in some filaments. The stars are scattered throughout, with some appearing as bright points and others as faint specks.

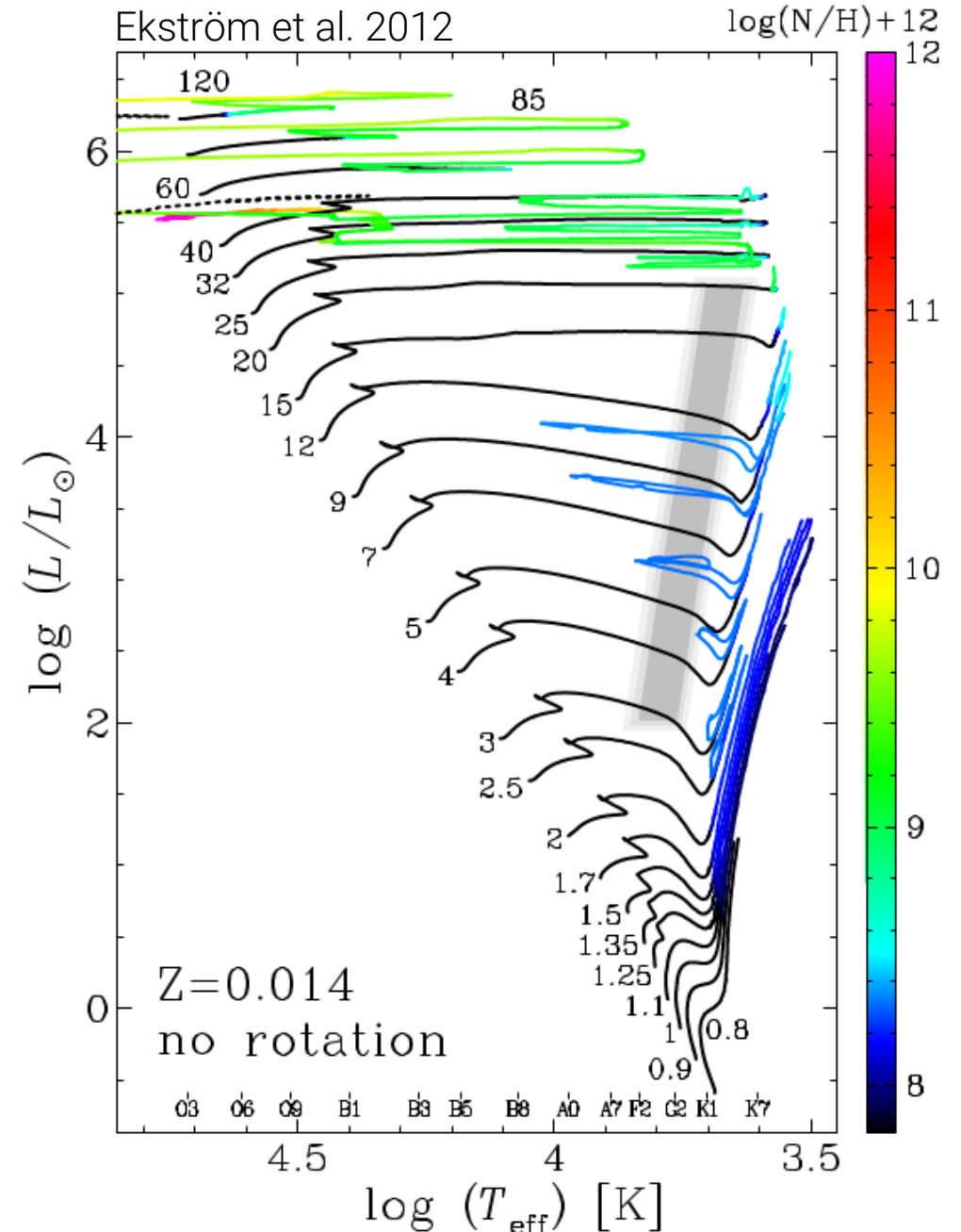
Radiation feedback

Radiation feedback

- **Main effects:** ionisation, heating and radiation pressure
- Depends on **stellar source** and **environment**, i.e. its optical depth/ability to absorb photons

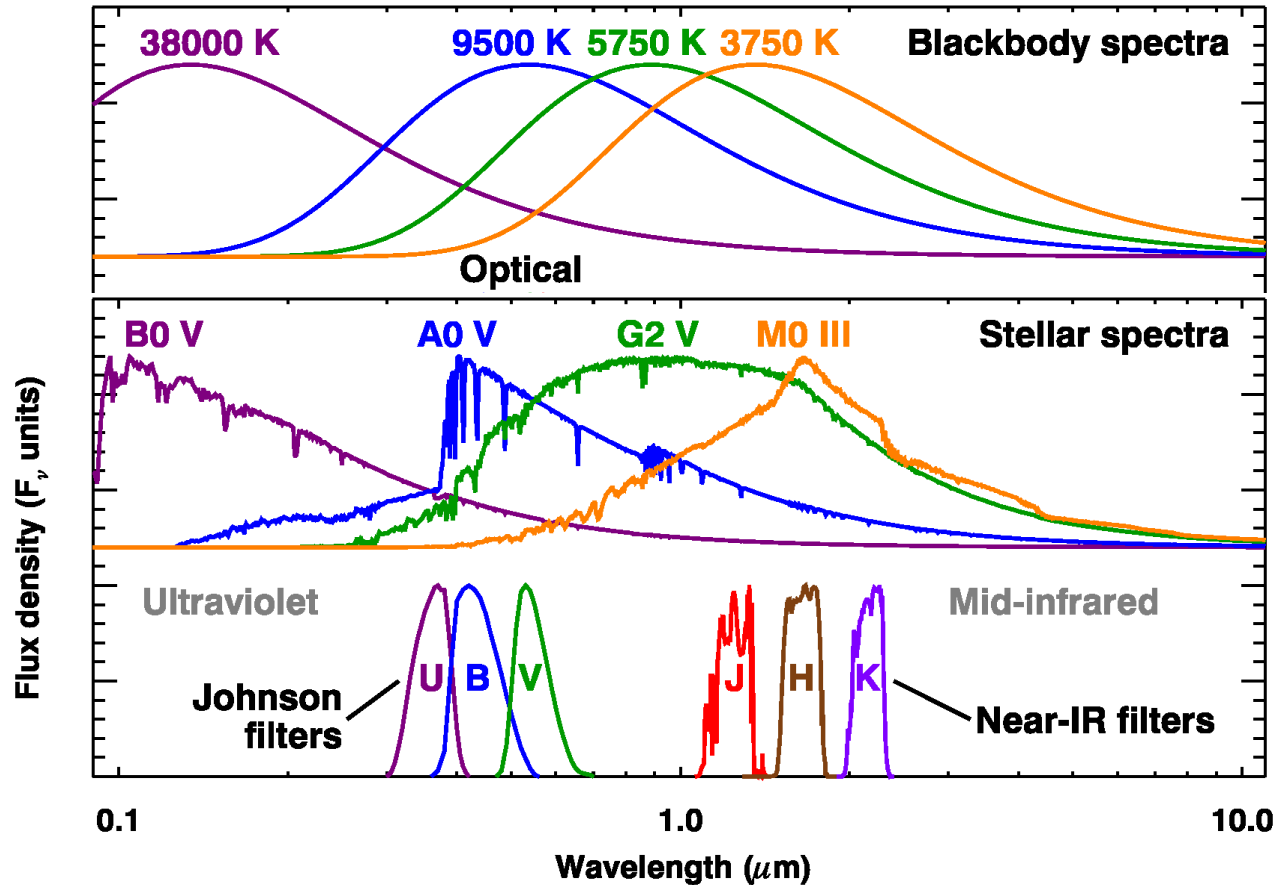
Example: Radiation pressure/radiative momentum flux: $\rightarrow P_{\text{rad}}(\mathbf{r}) = f_{\text{trap}} L / 4\pi r^2 c$
(f_{trap} : encapsulates uncertain photon absorption, e.g. $f_{\text{trap}} = 1$ all photons absorbed once; complex radiative transfer problem)

- Stellar sources: **steep mass-luminosity relation**
 - One very massive star can be worth millions of Suns in terms of radiative luminosity
 - Effective temperature/spectral energy distribution sets individual photon energies (e.g. UV vs IR)



Spectral energy distribution of stars ($\rightarrow T_{\text{eff}}$)

Credits: Gregory Sloan



- Traditionally, treat stars as black bodies
 $\rightarrow T_{\text{eff}}$ sets spectral energy distribution
- **But:** photons interact with gas in atmosphere, modifying the SED (e.g. Balmer break, absorption/emission lines)
 \rightarrow exact stellar ionising & heating radiation
- **Reminder:** T_{eff} of models such that *stellar flux* $L/4\pi R^2 = \text{black body flux of } \sigma T_{\text{eff}}^4$
 - Eddington grey atmosphere (T - τ relation)
 $\rightarrow T_{\text{eff}}^4 = 4/3 T^4 (\tau + 2/3) \rightarrow$ get T
 - Atmosphere **defined** at $\tau = 2/3 \rightarrow T_{\text{eff}} = T$
 \rightarrow half of photons escape freely ($e^{-2/3} \sim 0.5$)

The 30 Doradus starburst

Record holders:

- Most massive stars (up to $\sim 300 M_{\odot}$)
- Fastest rotating star (near break-up)
- $\sim 90 M_{\odot}$ runaway
- $\sim 150 M_{\odot}$ walkaway

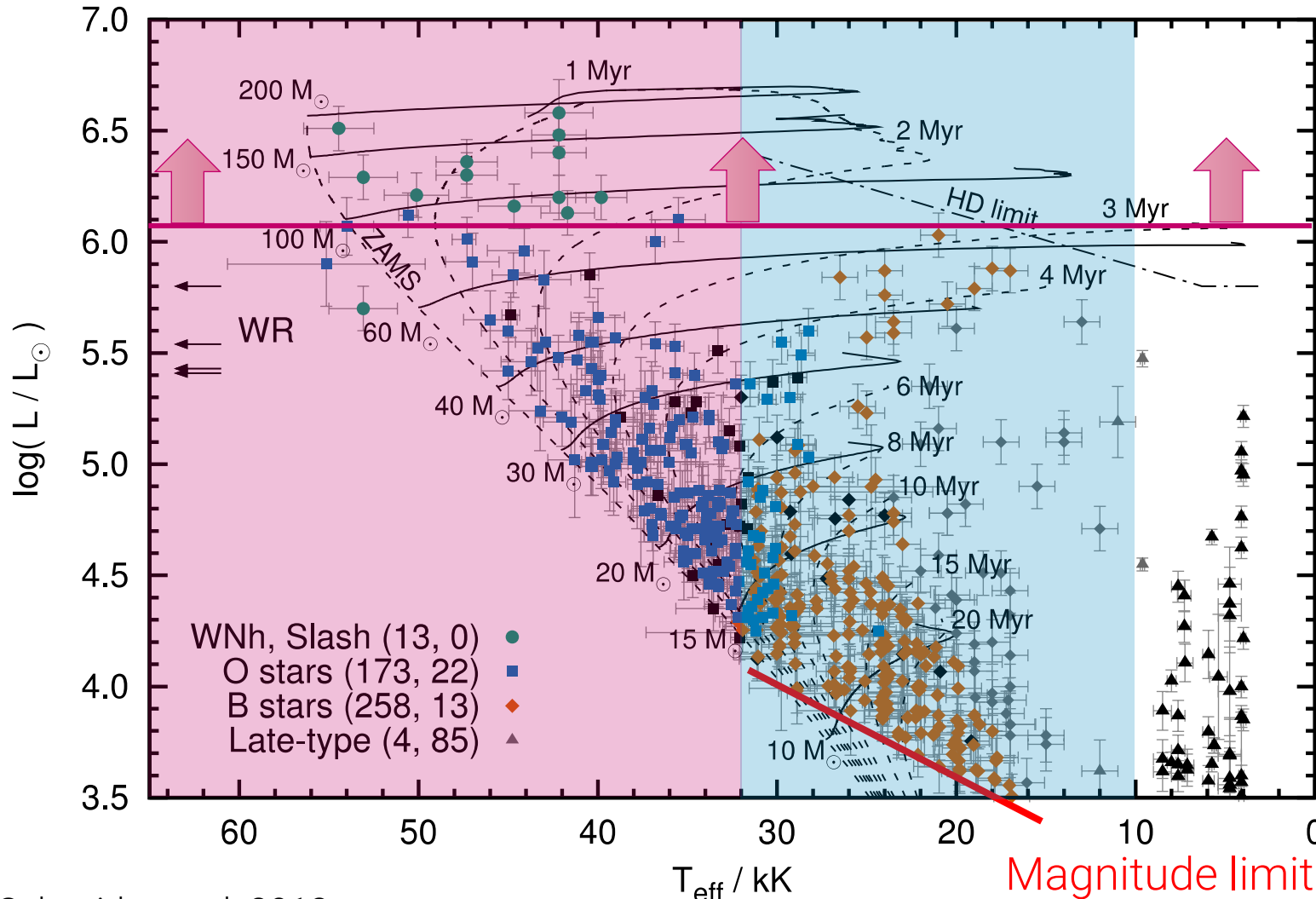
- **Brightest HII region** in Local Group (Kennicutt 1984)
- **Well known distance** (50 kpc; Pietrynski+2013)
- Prototype of more **distant starbursts**
 - similar to so-called **"Green Pea" galaxies** (→ re-ionisation)
- **>1000 OB stars** incl. several record holders
 - **ideal laboratory for massive star formation and evolution**

200 pc

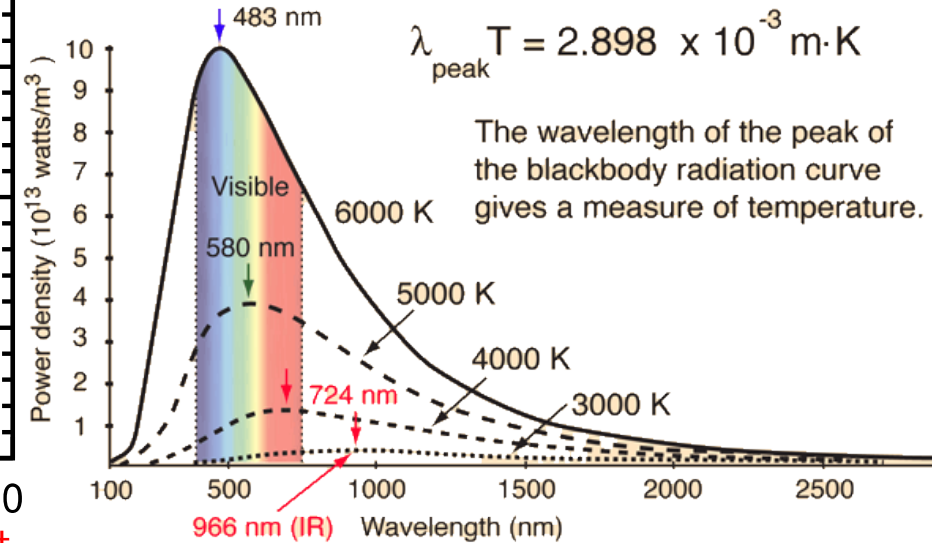


Massive stars in 30 Doradus

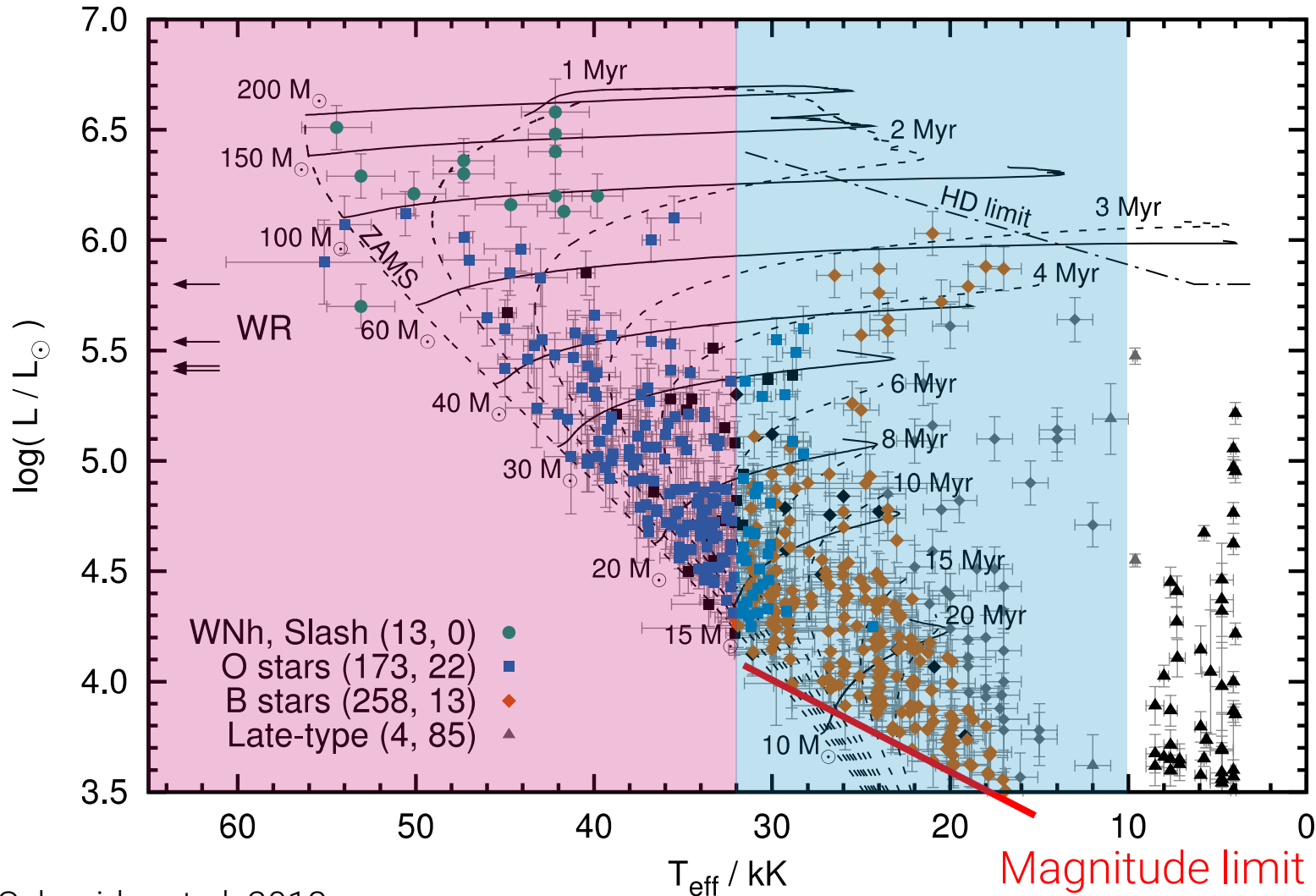
~25% of ionising radiation
from stars with $M_{\text{ini}} > 100 M_{\odot}$
(Doran+2013)



- One very massive star can be worth millions of Suns in terms of luminosity
- $T_{\text{BB}} \sim 32,000 \text{ K} * E_{\text{ion}}/13.6 \text{ eV}$
- Tail of BB: still H-ionising photons for $32 \text{ kK} < T_{\text{eff}} < 10 \text{ kK}$
- HII regions from OB stars



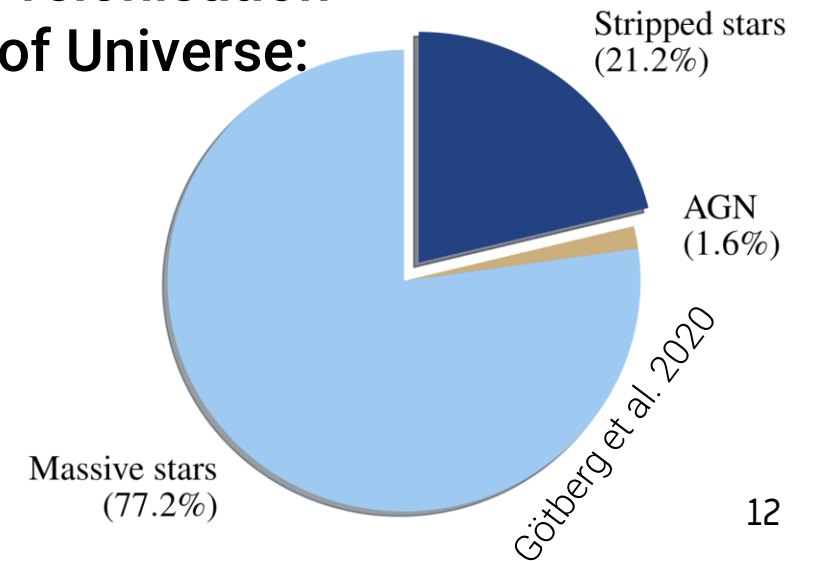
Massive stars in 30 Doradus



Schneider et al. 2018

- One very massive star can be worth millions of Suns in terms of luminosity
- $\rightarrow T_{\text{BB}} \sim 32,000 \text{ K} * E_{\text{ion}} / 13.6 \text{ eV}$
- Tail of BB: still many H-ionising photons for $32 \text{ kK} < T_{\text{eff}} < 10 \text{ kK}$
 \rightarrow HII regions from OB stars

H-reionisation of Universe:





Stellar winds

Stellar winds

- **Stellar wind:** mechanism that accelerates stellar surface layers to beyond *escape velocity* and thereby drives a (spherically-symmetric) *outflow*

- **Example: photon-driven winds:**

$$a_{\text{rad}} = \frac{1}{c} \int_0^{\infty} \kappa_{\nu}(r) F_{\nu}(r) d\nu \equiv \kappa_F(r) \frac{L}{4\pi cr^2}$$

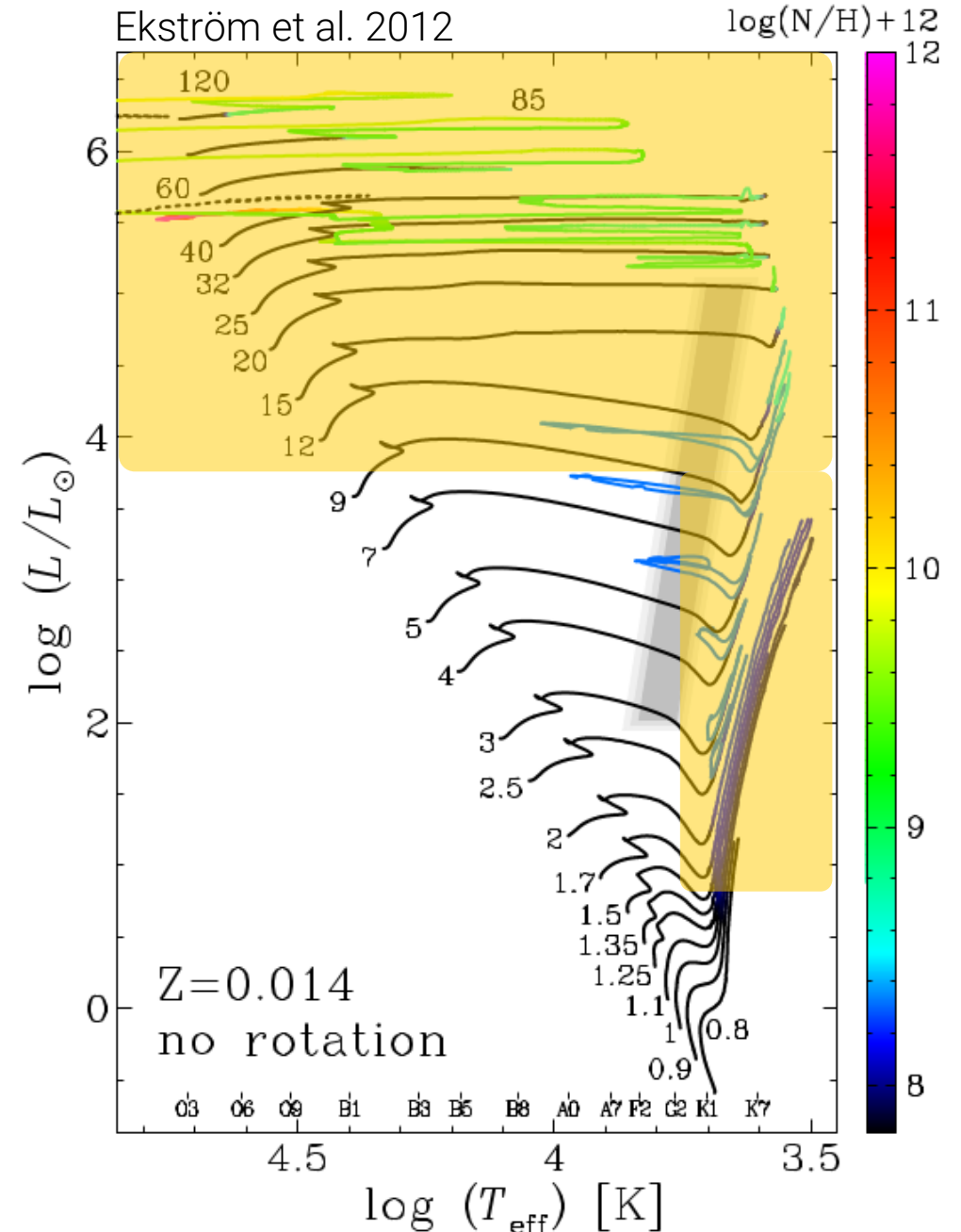
with the flux-weighted mean opacity $\kappa_F(r)$

- W.r.t. grav. acceleration, obtain **Eddington factor**

$$\Gamma_{\text{rad}}(r) = \frac{a_{\text{rad}}(r)}{g(r)} = \kappa_F(r) \frac{L}{4\pi cGM} \propto \kappa_F \frac{L}{M}$$

- $\Gamma_{\text{rad}} = 1$ & electron scattering, $\kappa_{\text{es}} = 0.2 (1+X) \text{ cm}^2 \text{ g}^{-1}$,

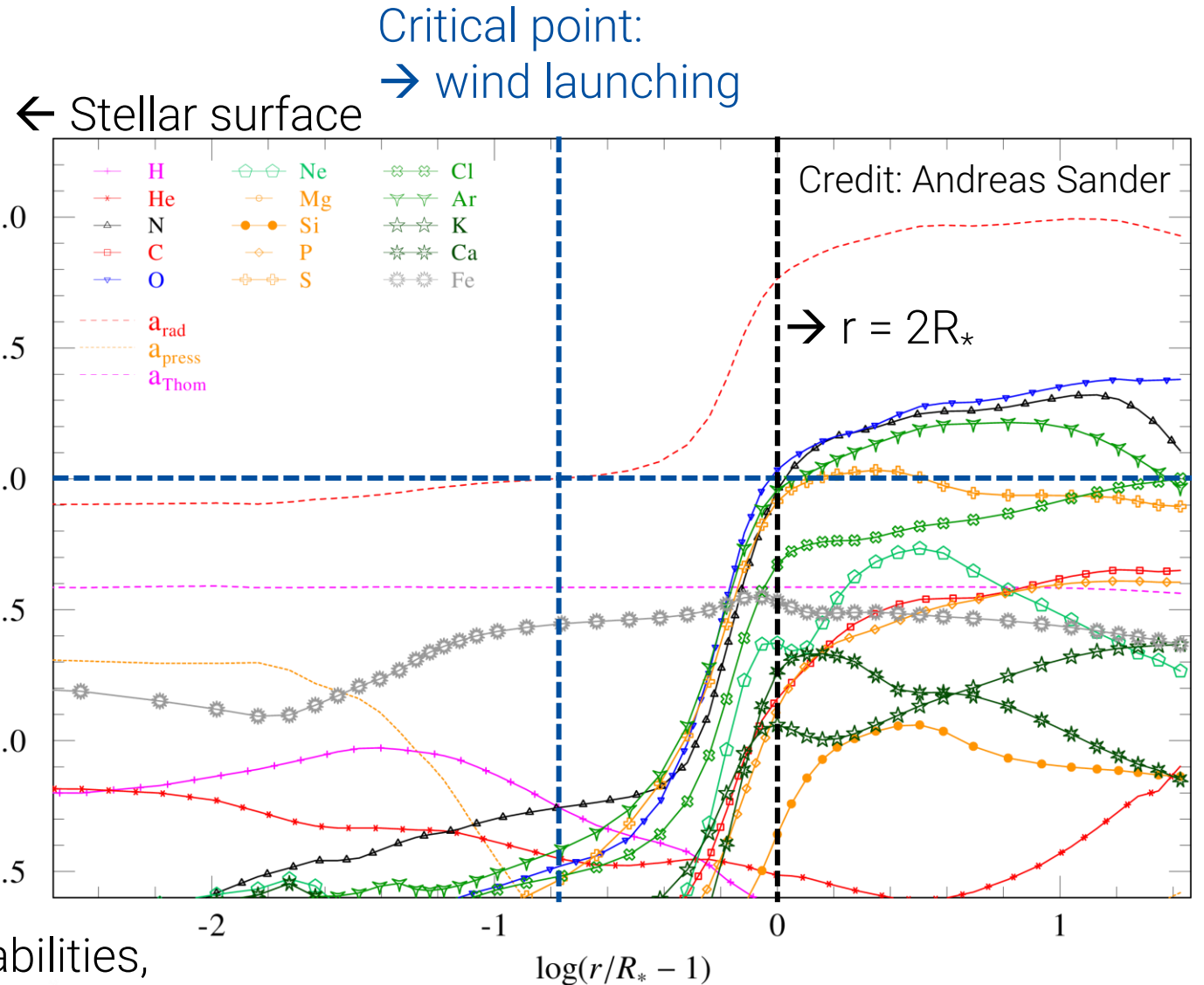
$$L_{\text{edd}} = \frac{4\pi cGM}{\kappa_{\text{es}}} \approx 3 \times 10^4 \left(\frac{M}{M_{\odot}} \right) L_{\odot} \quad \text{Eddington luminosity}$$



Stellar winds

Line driving

- **Massive-star winds are line-driven**
 - ζ Pup (O supergiant, $T_{\text{eff}} \sim 42$ kK)
- **Wind launching** (i.e. wind mass loss rate)
 - iron (Fe) most relevant
- **Terminal wind velocity**
 - other atoms dominant (e.g. O, N, C)
- **→ strong iron/metallicity dependence**
 - $\dot{M} \propto Z^{0.5 \dots 0.8}$
- **Winds in low-Z environments/early Universe much weaker**
 - less wind feedback
- Other driving mechanisms: pulsations, instabilities, photon driving on molecules & dust, ...



Stellar winds

Feedback by stellar winds

- Integrated wind momentum

$$p_{\text{wind}} = \int_0^{\tau} \dot{M}(t) v_{\infty}(t) dt$$

- Integrated wind energy

$$E_{\text{wind}} = \int_0^{\tau} \frac{1}{2} \dot{M}(t) v_{\infty}^2(t) dt$$

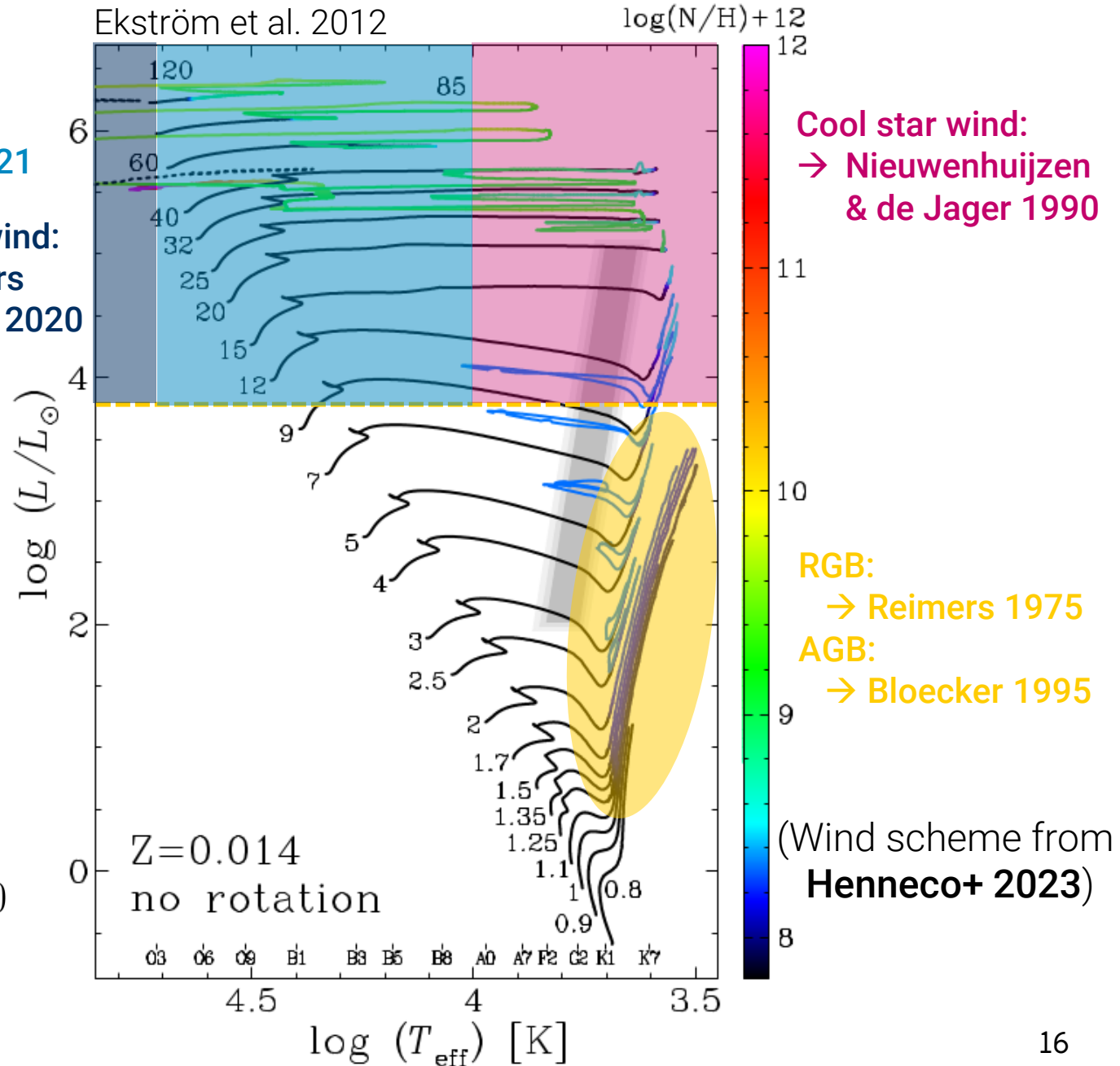
- Needs: wind velocity & mass loss rates

$$v_{\infty} = \underbrace{\sqrt{\frac{2GM}{R}}}_{=v_{\text{esc}}} \begin{cases} 0.7 & T_{\text{eff}} < 12.5 \text{ kK} \\ 1.3 & 12.5 \leq T_{\text{eff}}/\text{kK} < 21.0 \\ 2.6 & T_{\text{eff}} \geq 21 \text{ kK} \end{cases}$$

(Lamers et al. 1995)

Hot star wind:
 → Vink+ 2000
 → Björklund+ 2021

Wolf-Rayet star wind:
 → Nugis & Lamers
 → Sander & Vink 2020



Typical winds of massive stars

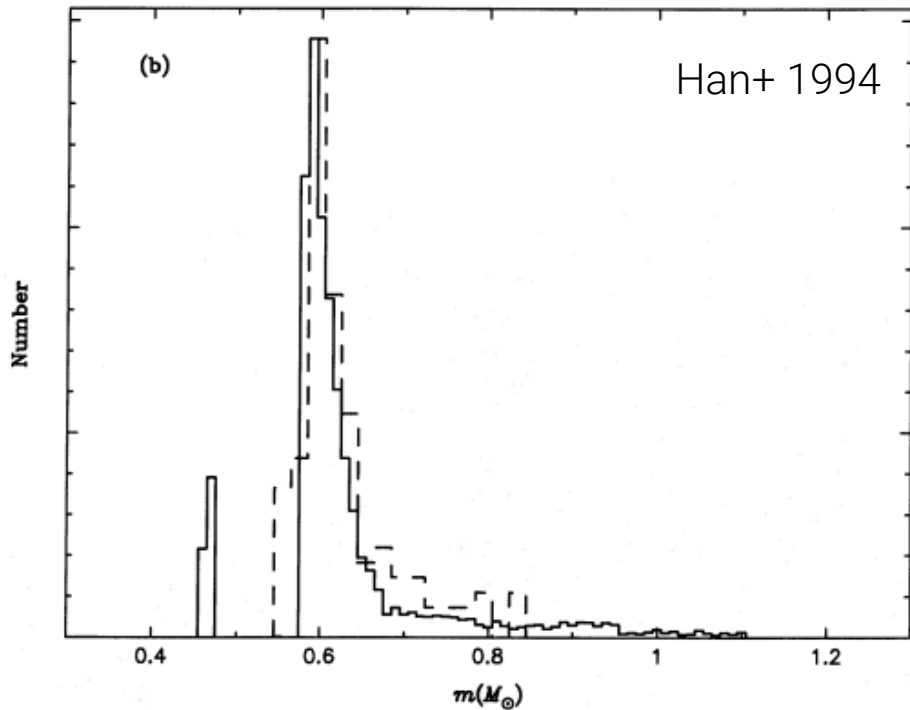
Table 1 Typical wind parameters for different stellar types^a

Vink 2022

Type	T_{eff} (kK)	M (M_{\odot})	v_{∞} (km s ⁻¹)	\dot{M} (M_{\odot} year ⁻¹)
O	30–45	20–60	2,000–3,500	10^{-7} – 10^{-5}
WNh	35–50	80–300	1,500–3,000	10^{-4}
BSG	15–25	15–30	500–1,500	10^{-7} – 10^{-5}
YSG	5–10	10–25	50–200	10^{-6} – 10^{-4}
RSG	3–5	10–25	10–30	10^{-7} – 10^{-4}
LBV low-L	10–15	15–25	100–200	10^{-5}
LBV high-L	10–30	40–	200–500	10^{-4} – 10^{-3}
cWR	90–200	10–30	1,500–6,000	10^{-5} – 10^{-4}
Stripped He	50–80	1–5	1,000	10^{-8}

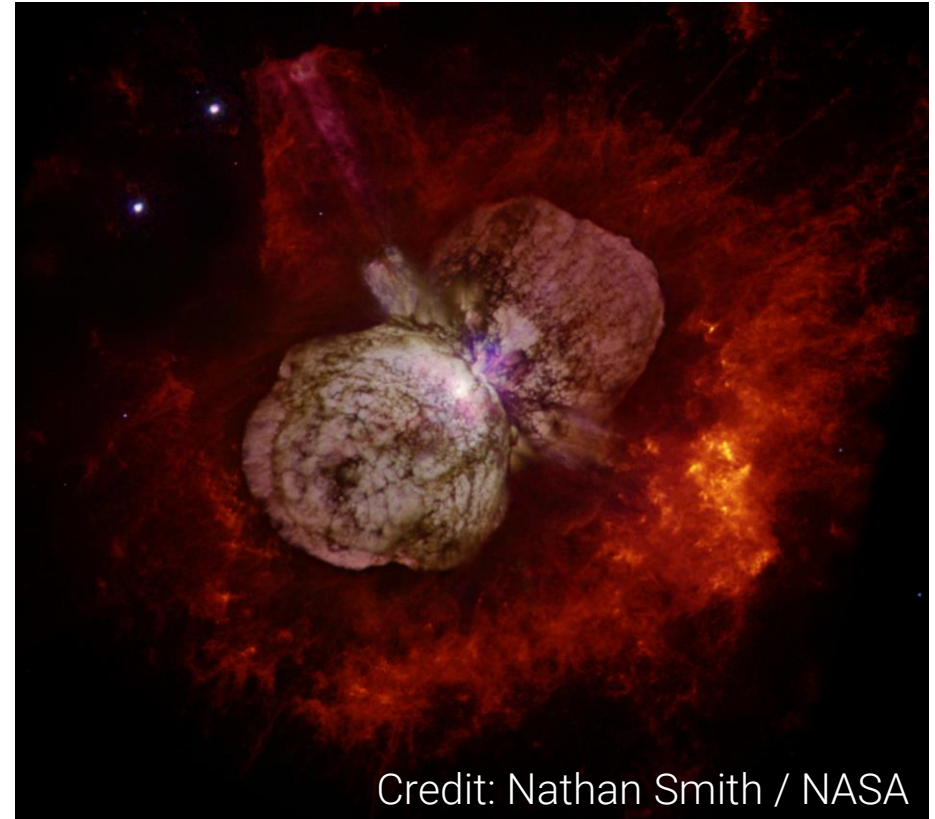
Envelope instabilities, eruptions & outflows

- Why do most WDs have $\sim 0.6 M_{\odot}$?



- Envelope instability once $E_{\text{bind}} > 0$
- Significant mass loss/feedback

- Luminous blue variables (e.g. S Dor stars)



- Envelope instability? Giant eruptions with up to near SN energies ($\sim 10^{50}$ erg)



Accretion, jets & outbursts

Accretion

- **Trivial statement:** to form stars, need accretion
- **Maybe less obvious:** accretion releases energy, hence produces feedback:
 - mass ΔM accreting on star of mass M and radius R releases grav. potential energy $GM\Delta M/R$
 - For mass accretion rate \dot{M} , produce luminosity

$$L_{\text{acc}} \approx \frac{GM\dot{M}}{R}$$

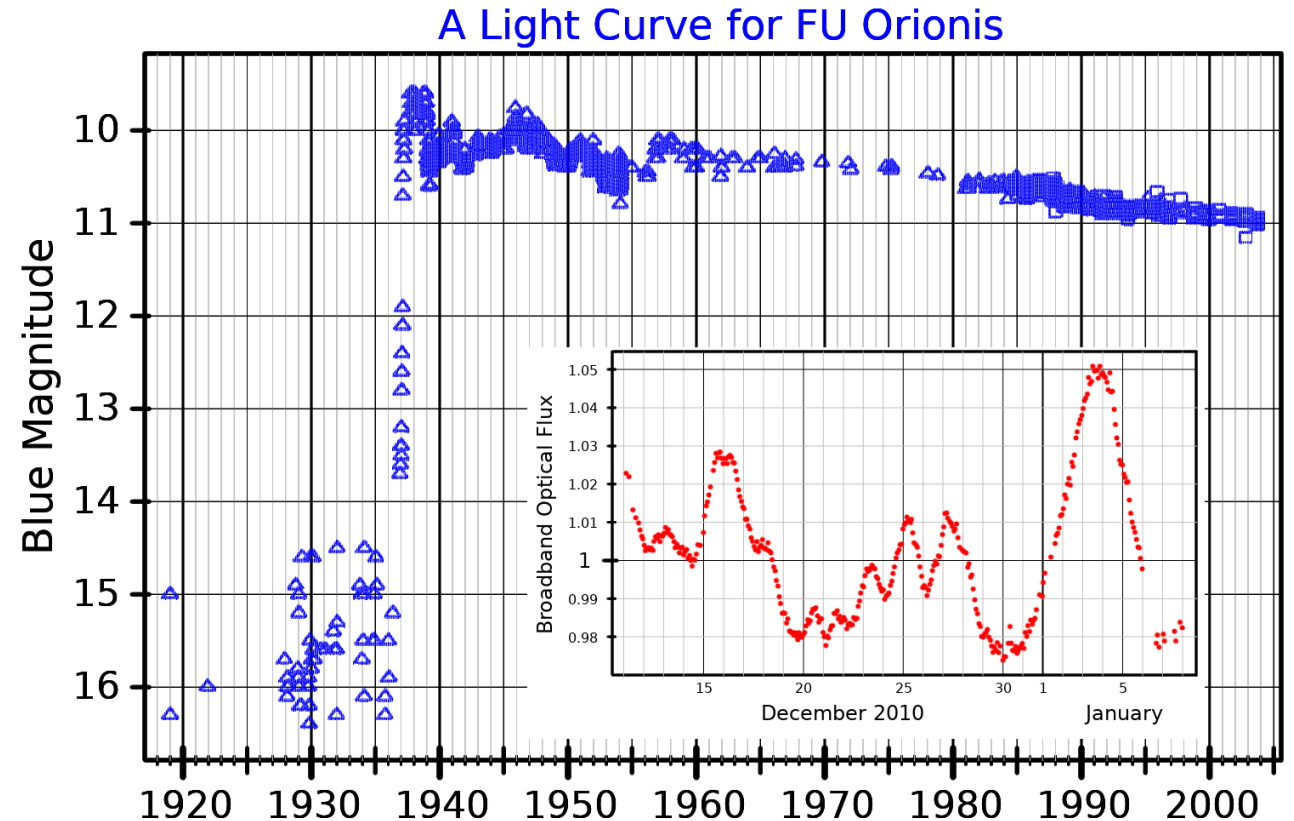
- Not all of released grav. potential energy radiated away (e.g. kinetic/rotational energy, magnetic fields, driving of outflows, ...)

Comparison

- 1 M_{\odot} star accreting at $10^{-6} M_{\odot} \text{ yr}^{-1}$:
 - $L_{\text{acc}} \sim 15 L_{\odot}$
 - $L_{\text{acc}} > L_{*}$
- $>9 M_{\odot}$ stars accreting at $<10^{-3} M_{\odot} \text{ yr}^{-1}$:
 - $L_{\text{acc}} < L_{*}$
- *Rule of thumb:* once there are OB stars, they dominate the radiative feedback

Episodic accretion and outbursts

- Accretion not steady but episodic & variable
- Occasionally, have **enhanced accretion rates**
 - **Outburst-like increase of L_{acc}**
 - Example: **FU Orionis-like outbursts**
- What causes enhanced accretion? Unknown, there are several ideas (e.g. Audard+2014, Vorobyov+2021, Borchert+2022a,b):
 - MRI triggered by ionisation increase
 - Clump accretion by disk fragmentation
 - Close encounters with other stars
 - Planet-disk interaction
 - ...

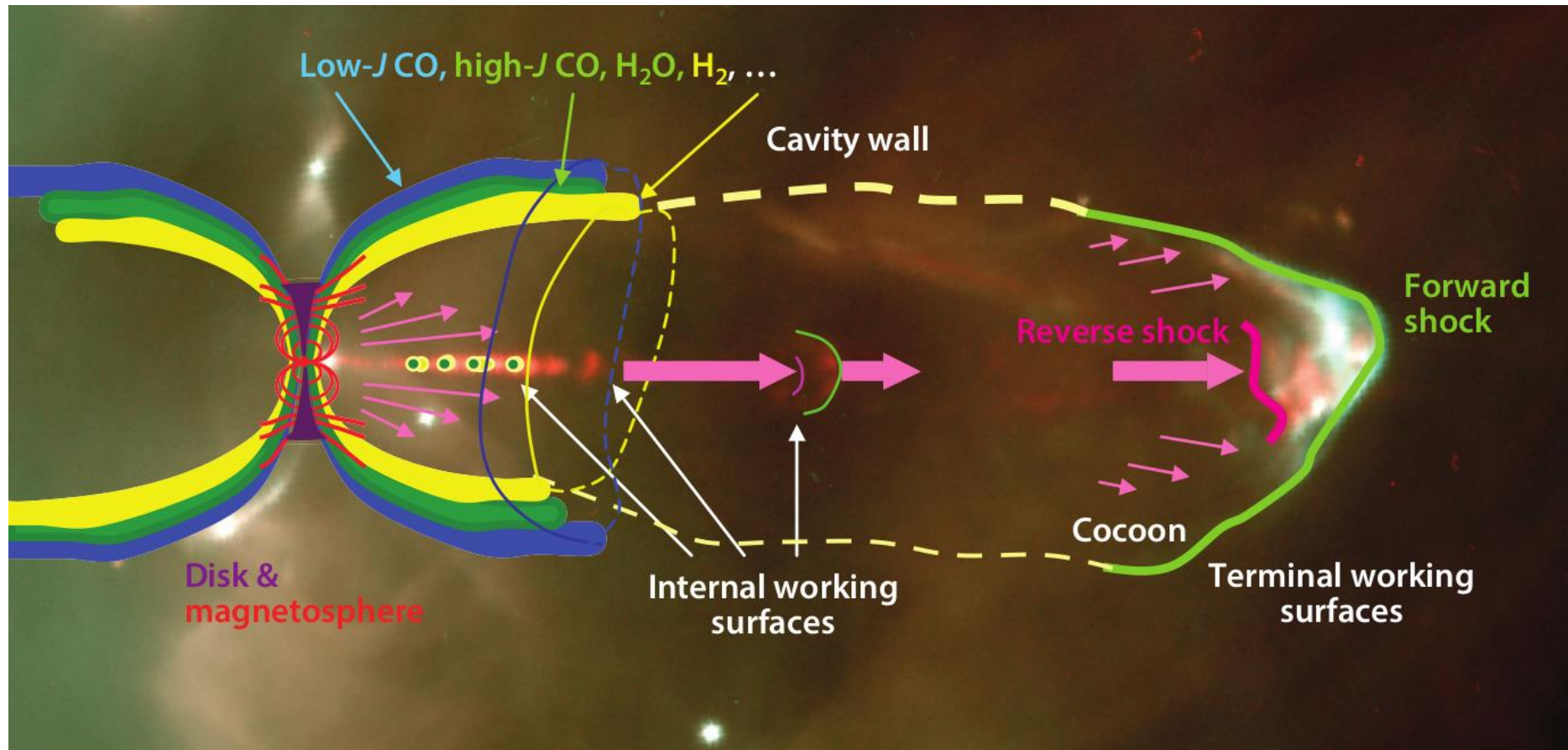


Long-term data from Clarke+2005 and short timescale variations from Siwak+2013 (inset).

Figure Credit: Wikipedia

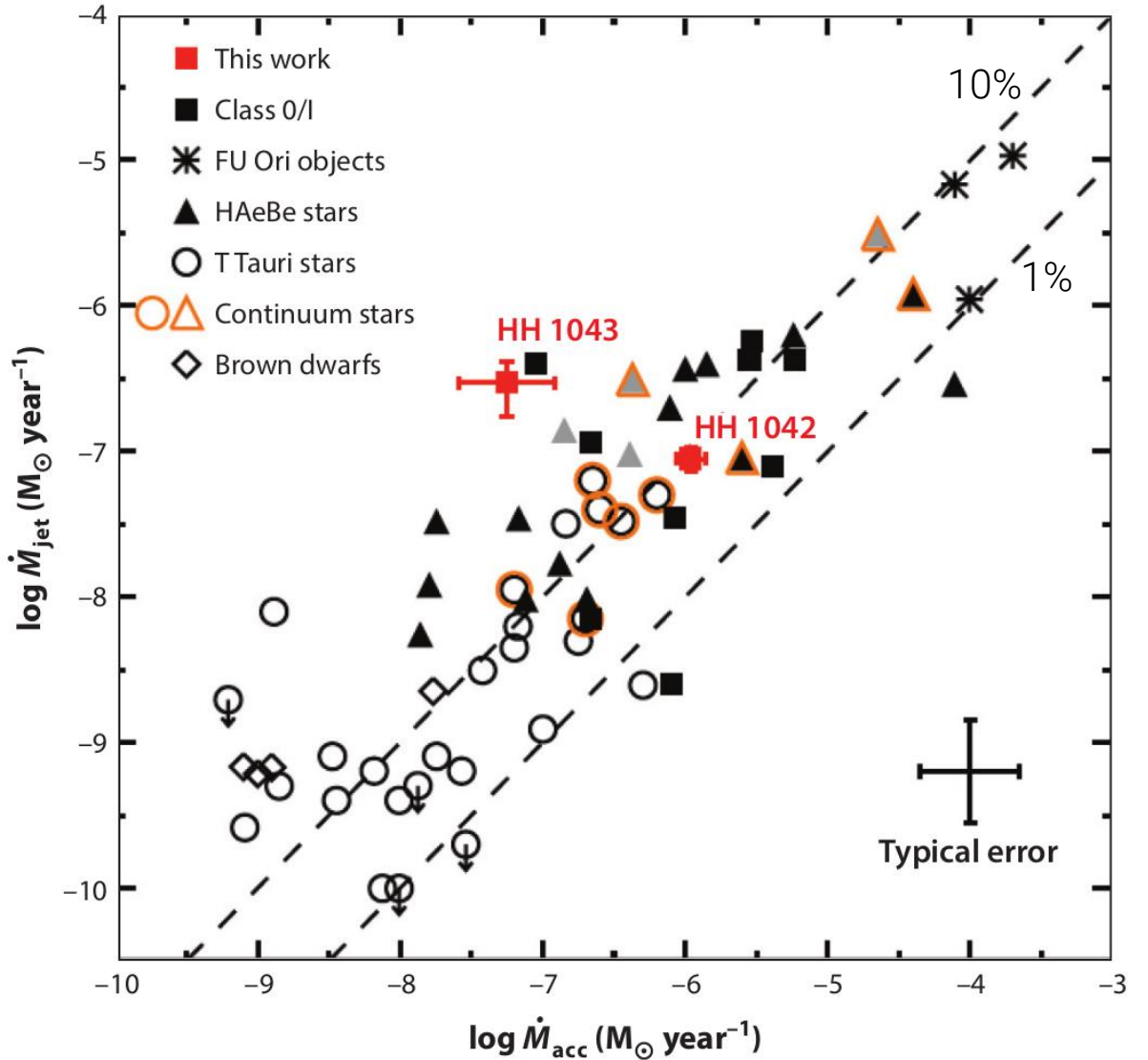
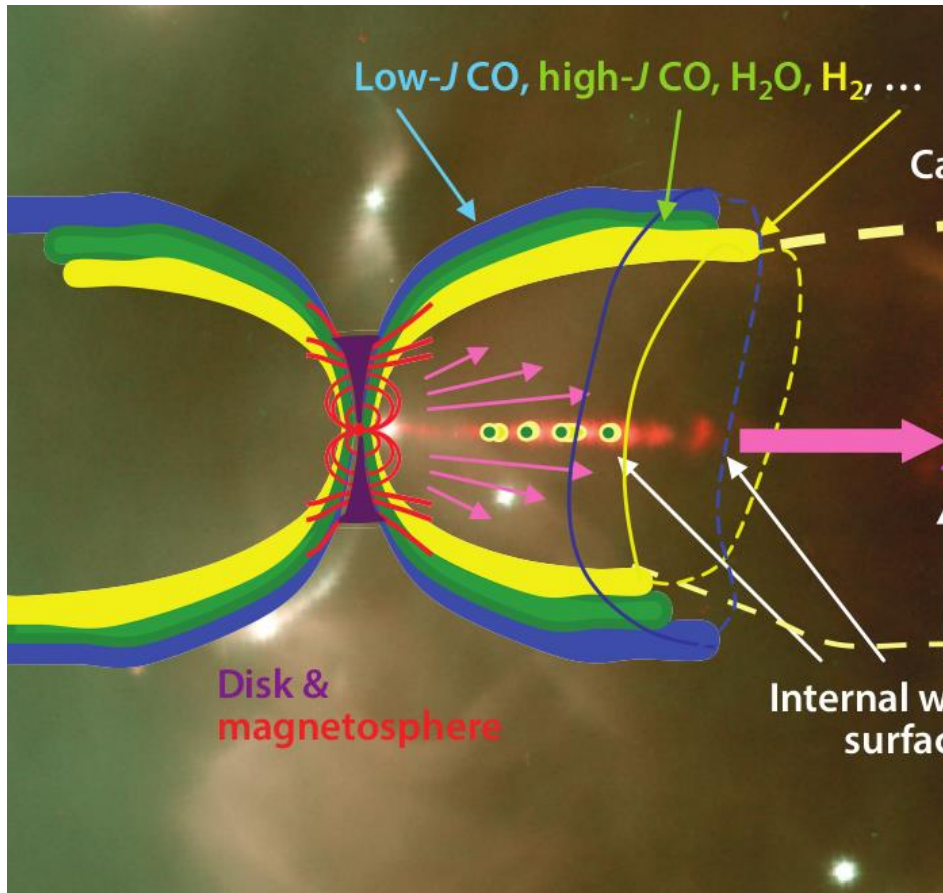
Jets and bipolar outflows

- Review on protostellar outflows: **Bally 2016**
→ bipolar outflows/jets are magnetically-driven



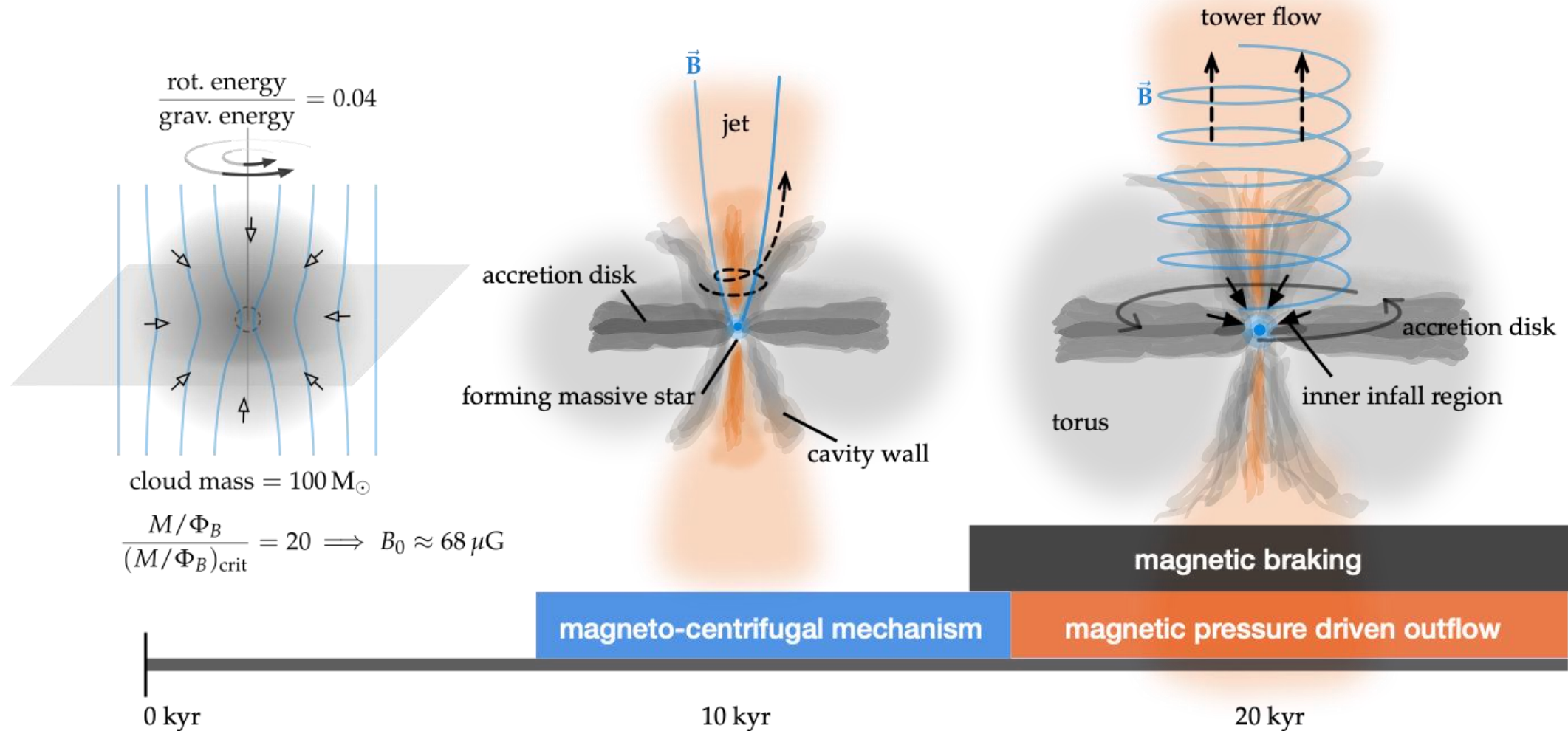
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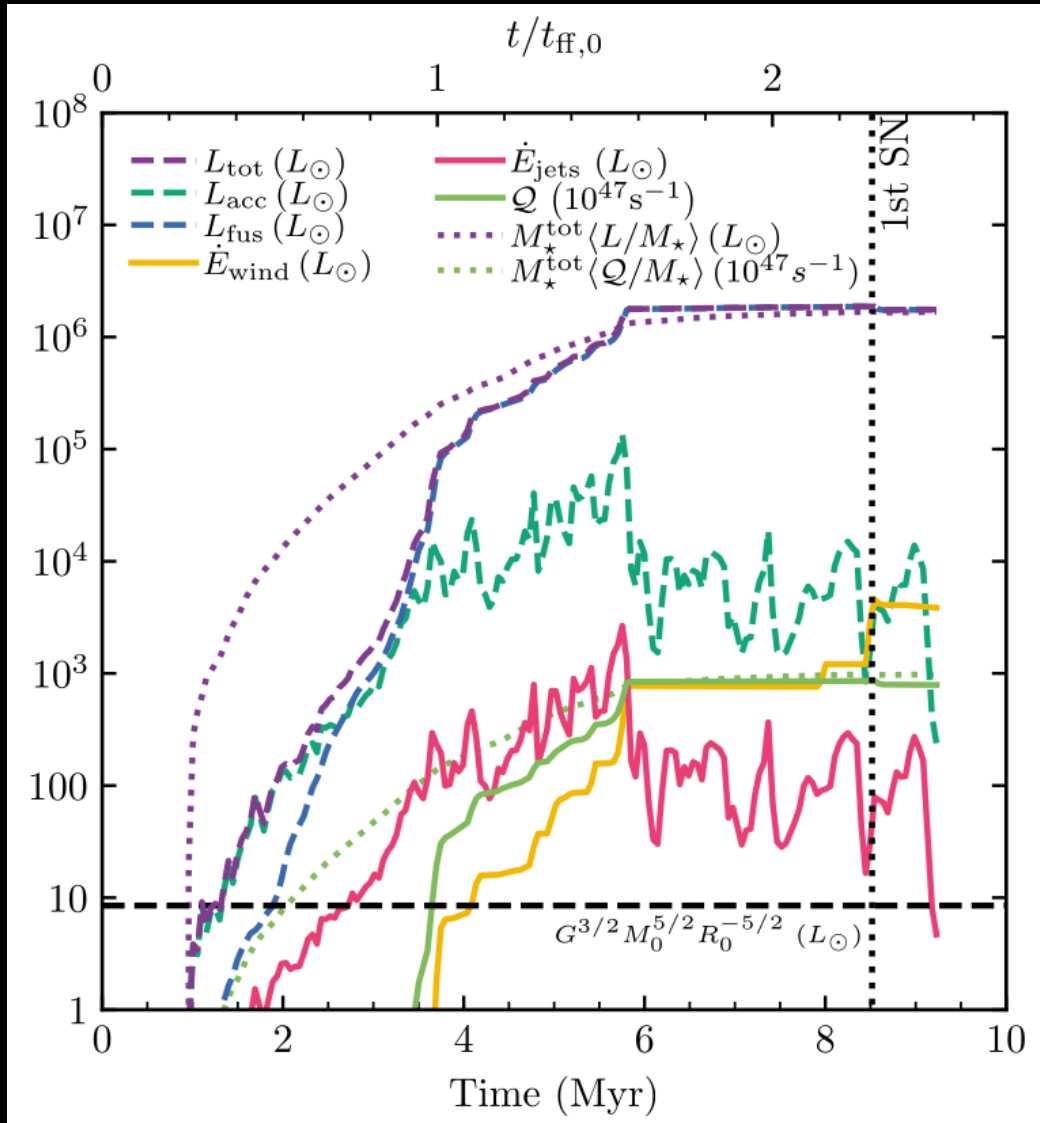


Jets and bipolar outflows

- Outflows: **magneto-centrifugal** (Blandford & Payne 1982) & **magnetic pressure** driven (Lynden-Bell 2003)



Jets and bipolar outflows



Grudic+2022; initial cloud mass $2 \times 10^4 M_{\odot}$, radius 10 pc

- Exact launching mechanism of jets unclear (so-called X-winds also possible, see Shu+1994) → see reviews by Pudritz & Ray 2019 and Krumholz & Federrath 2019
- Effective modelling (as in Cunningham+2011): fraction f_m of accreted mass is ejected in bipolar outflow of velocity

$$v_{\text{jet}} = f_K \sqrt{\frac{GM}{R}}$$

into cone of angular size θ_0

$$\xi(\theta, \theta_0) = \left(\ln \left(\frac{2}{\theta_0} \right) \sin^2 \theta + \theta_0^2 \right)^{-1}$$

- Initially, jets dominate feedback momentum, but not energy (Grudic+2022)

Supernovae


A multi-wavelength image of a supernova remnant, showing a complex, filamentary structure of glowing gas in various colors (red, orange, yellow, green, blue, purple) against a dark background. The central region is the most intense, with a bright purple and blue core. The outer layers are more diffuse and show a mix of red, orange, and green. The overall shape is roughly circular but with irregular, filamentary edges.

Credit: X-ray: NASA/CXC/J.Hester (ASU); Optical: NASA/ESA/J.Hester & A.Loll;
Infrared: NASA/JPL-Caltech/R.Gehrz

Supernovae

Rates:

- about 5 per second in the Universe
- several 100 discovered per year
- about 2–3 per century in Milky Way



Supernova 1987A
discovered 24 Feb 1987
in LMC

This image shows a bright, glowing supernova remnant in the Large Magellanic Cloud. The central region is a brilliant white-yellow color, surrounded by a diffuse, reddish-pink nebula. The background is filled with numerous stars of various colors.




exploding star:
Sanduleak -69° 202a
blue supergiant of ~17 Msun

This image shows a field of stars in the Large Magellanic Cloud. A white arrow points to a specific star, Sanduleak -69° 202a, which was the progenitor of Supernova 1987A. The star is a blue supergiant with a mass of approximately 17 solar masses.

© Anglo-Australian Observatory/David Malin Images

Type II SN



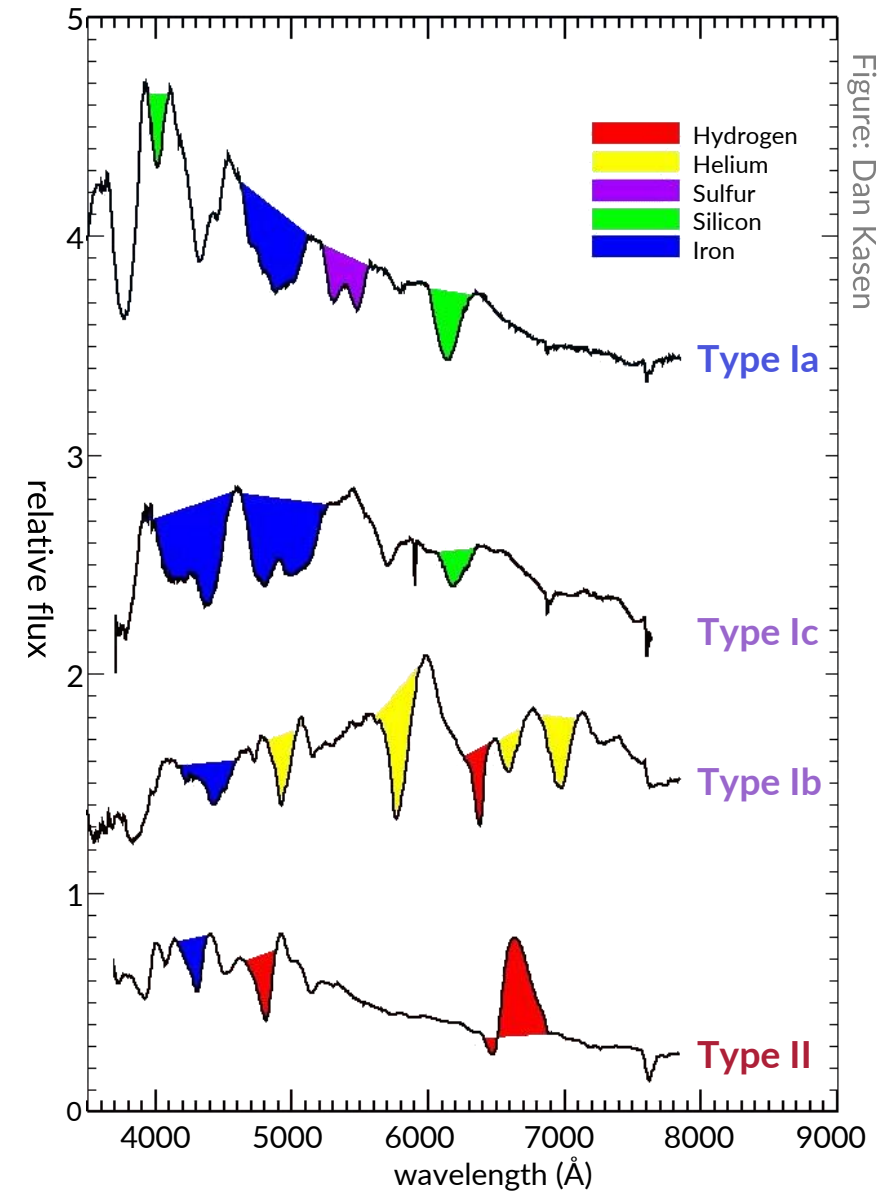
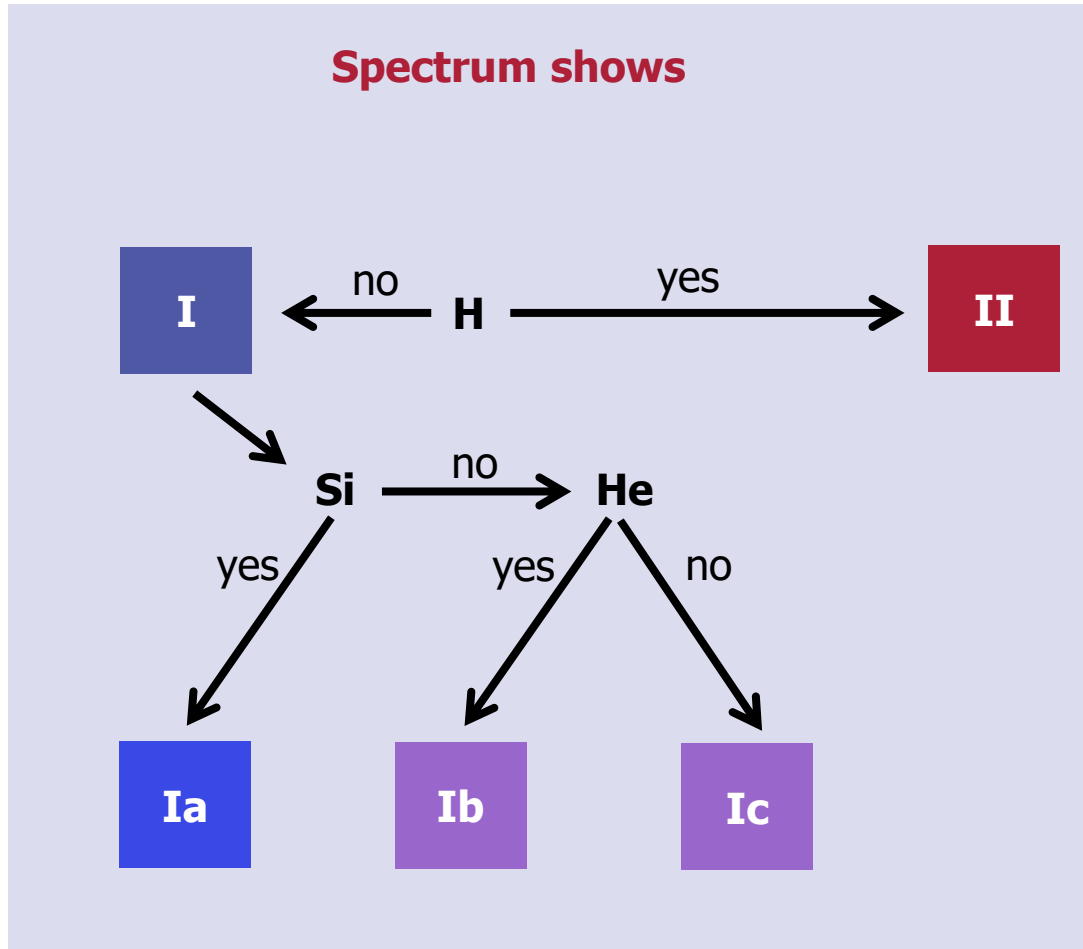
Supernova 1994D
at the edge of NGC 4526
image: HST NASA/ESA

This image shows a galaxy, NGC 4526, with a bright supernova remnant at its edge. The galaxy is a barred spiral, and the supernova remnant is a bright, glowing object. The image is a composite of Hubble Space Telescope observations.

Type Ia SN

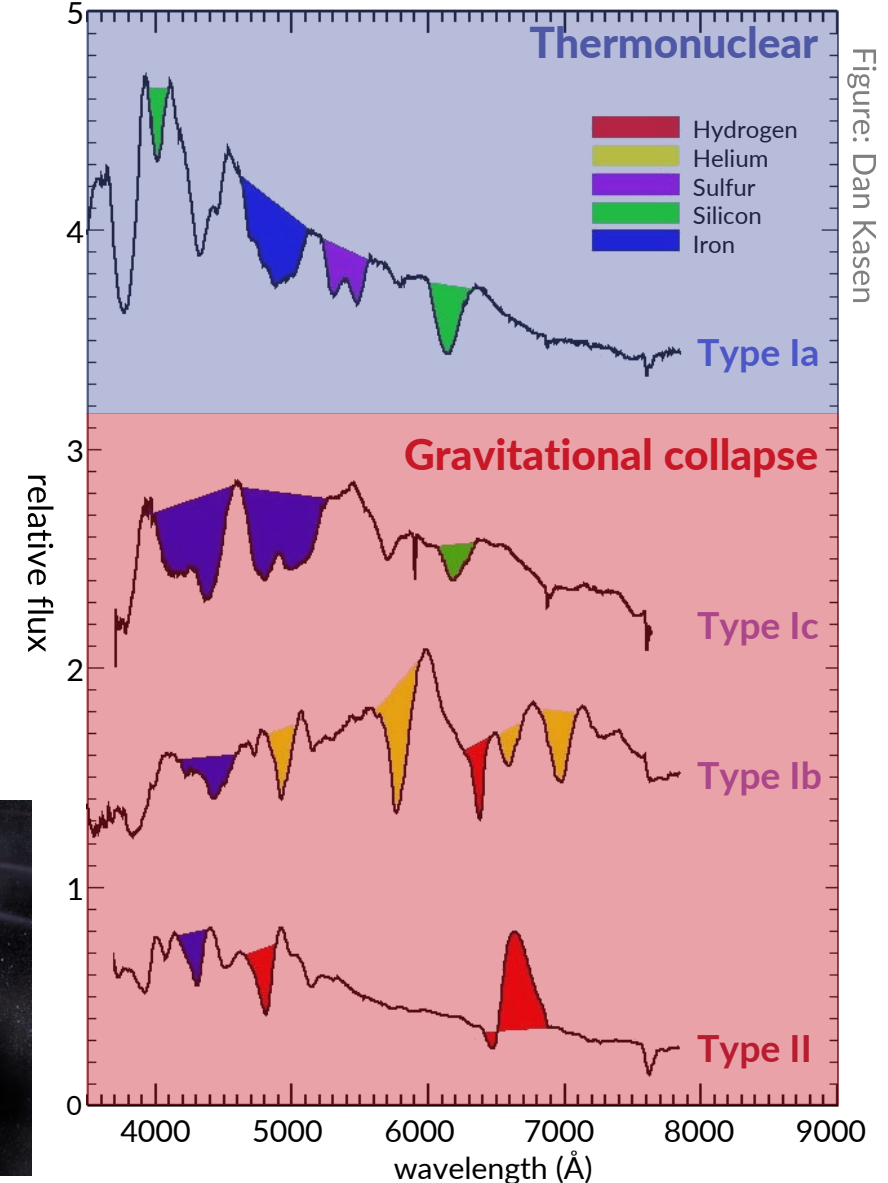
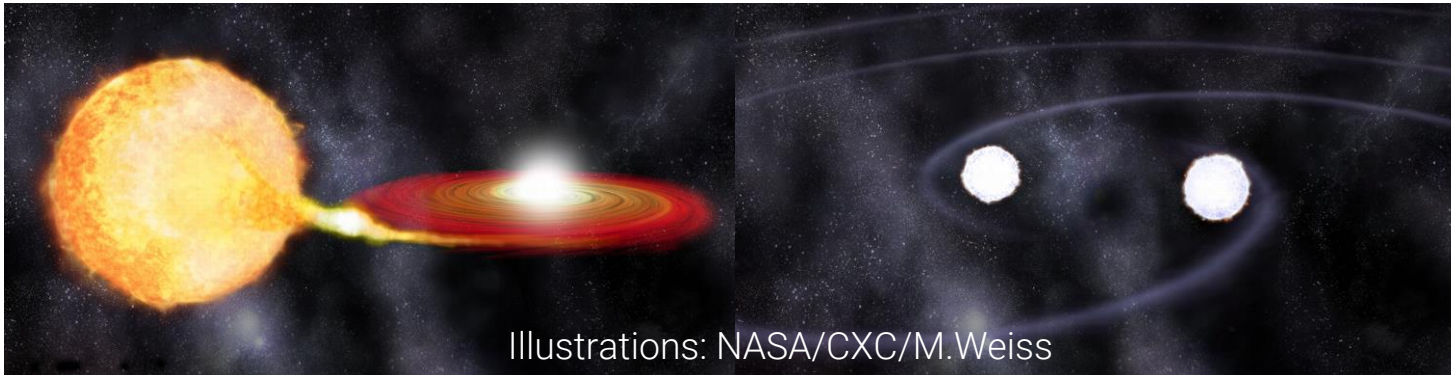
Supernova classification

- Classification according to spectral features



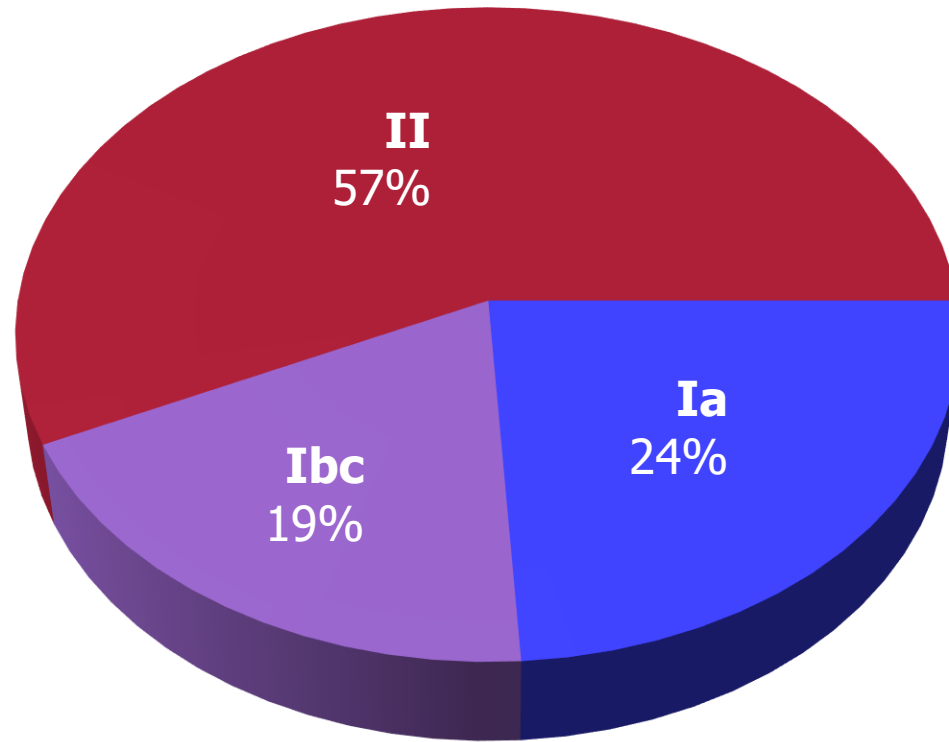
Supernova classification

- **Classification according to explosion physics**
- **SN Ia: exploding white dwarfs**
(single or double degenerate channel)
 - long delay time (easily 10^9 yr after star birth)
 - mostly irrelevant for mechanical feedback
 - not for chemical enrichment (e.g. iron!)
- **Here: focus on SNe from massive stars (short delay)**



Supernova rates

- Volume-limited (Li et al. 2011)



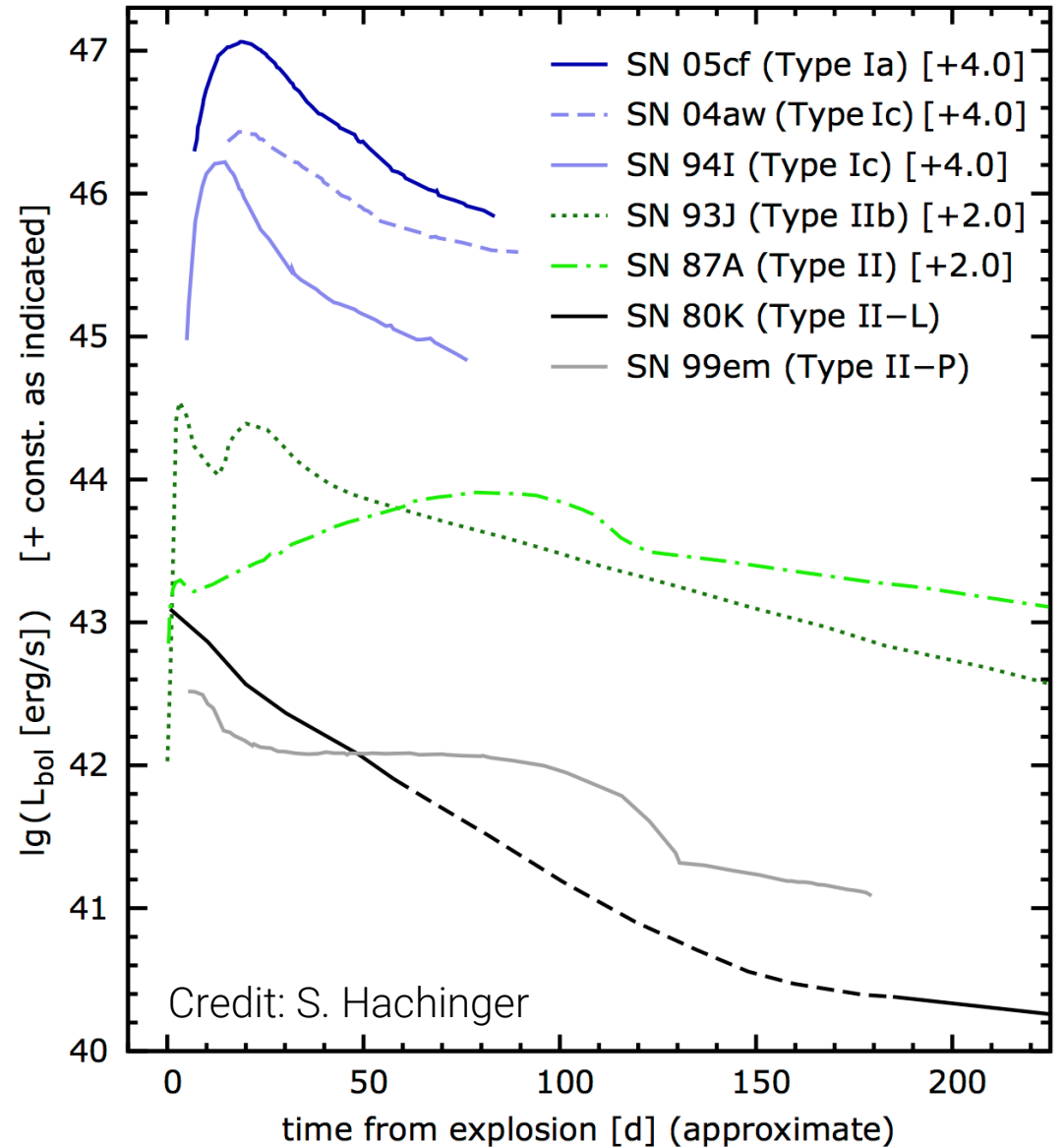
Supernova energetics

- **Threefold feedback:**

1. **Supernova lightcurve** → radiation
2. **Supernova explosion energy** → mechanical
3. **Cosmic rays** → energetic particles (later)

SN lightcurve (usually neglected)

- **SN II: $\sim 10^{49}$ erg; SN Ia: $\sim 10^{49} - 10^{50}$ erg**
- Peak luminosity: $10^{41} - 10^{43}$ erg/s
→ up to 100 times that of $10^5 M_{\odot}$ SF region
(e.g. 30 Doradus → $10^{41} - 10^{42}$ erg/s)
- Photon luminosity of one massive star over entire life: $10^{52} - 10^{53}$ erg
- **Energy sources:** thermal emission from hot plasma (+ H & He recombination), interaction with mass shell, nuclear (Ni decay)



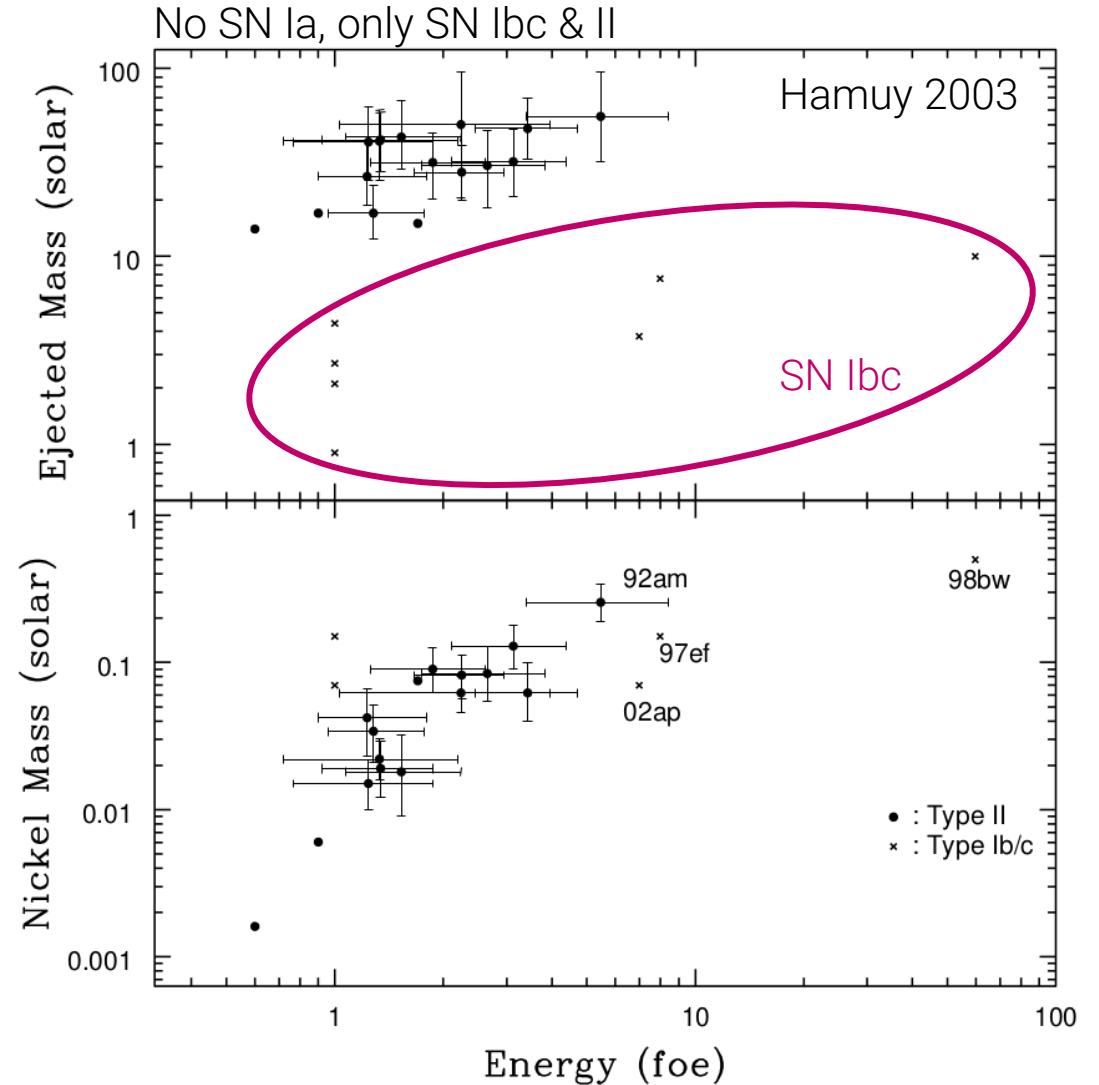
Supernova energetics

- **Threefold feedback:**

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Supernova explosion energy

- **SN II & Ia: $\sim 0.5 - 5 \times 10^{51}$ erg** (10^{51} erg = 1 foe = 1 B [Bethe])
- Ejecta velocities: $10^3 - 10^4$ km/s
- Nickel masses: $0.001 - 0.3 M_{\odot}$
- Ejecta masses: $\sim 0.1 - 60 M_{\odot}$



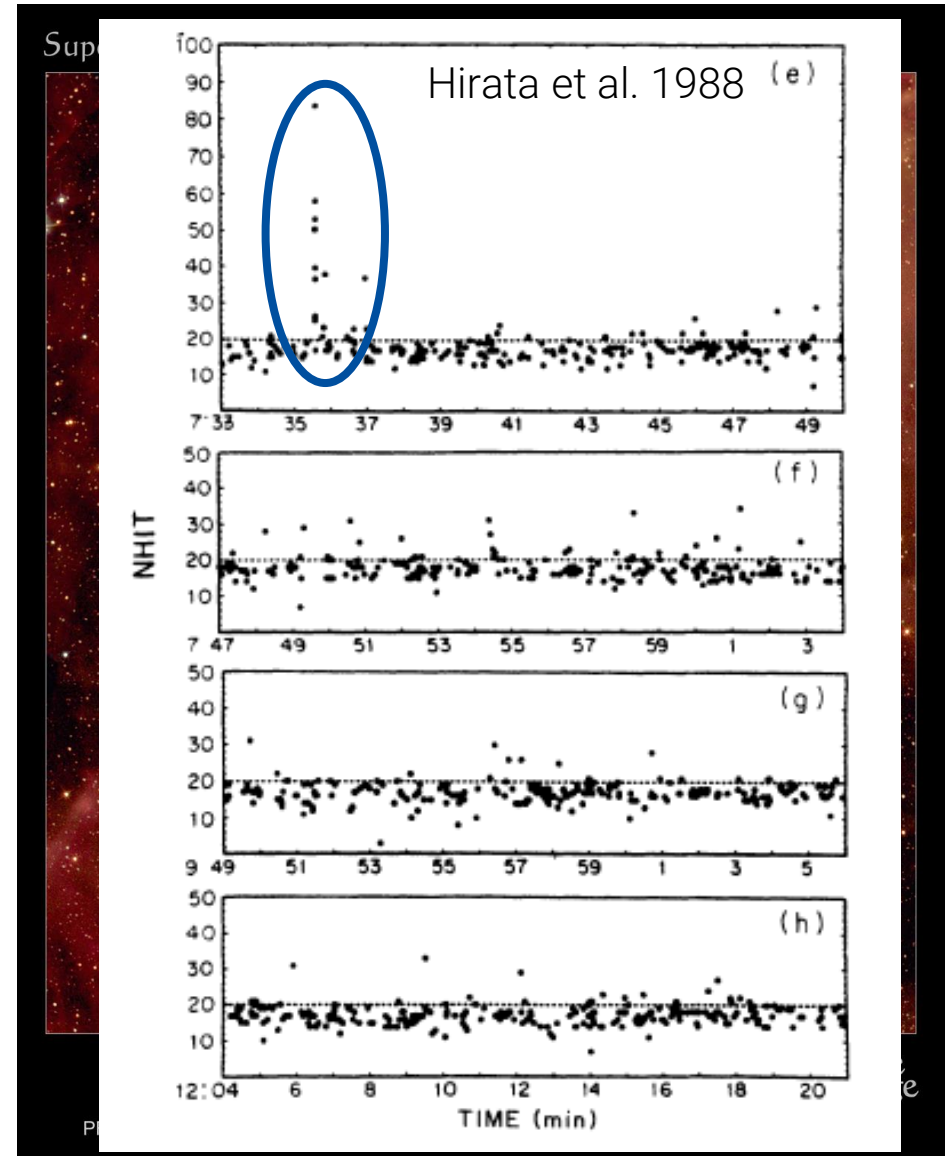
Supernova energy source

Core collapse SNe

- **Observed kinetic energy: $\sim 10^{51}$ erg**
- **Energy source:** collapse to neutron star?
→ “gravity bomb”
- Release of gravitational binding energy: Earth-size iron core ($R_{\text{Fe}} \sim 3000$ km) collapses to proto-NS ($R_{\text{NS}} \sim 12$ km)

$$\Delta E_{\text{bind}} = \frac{GM^2}{R_{\text{NS}}} - \frac{GM^2}{R_{\text{Fe}}} \approx \frac{GM^2}{R_{\text{NS}}} \approx 10^{53} \text{ erg}$$

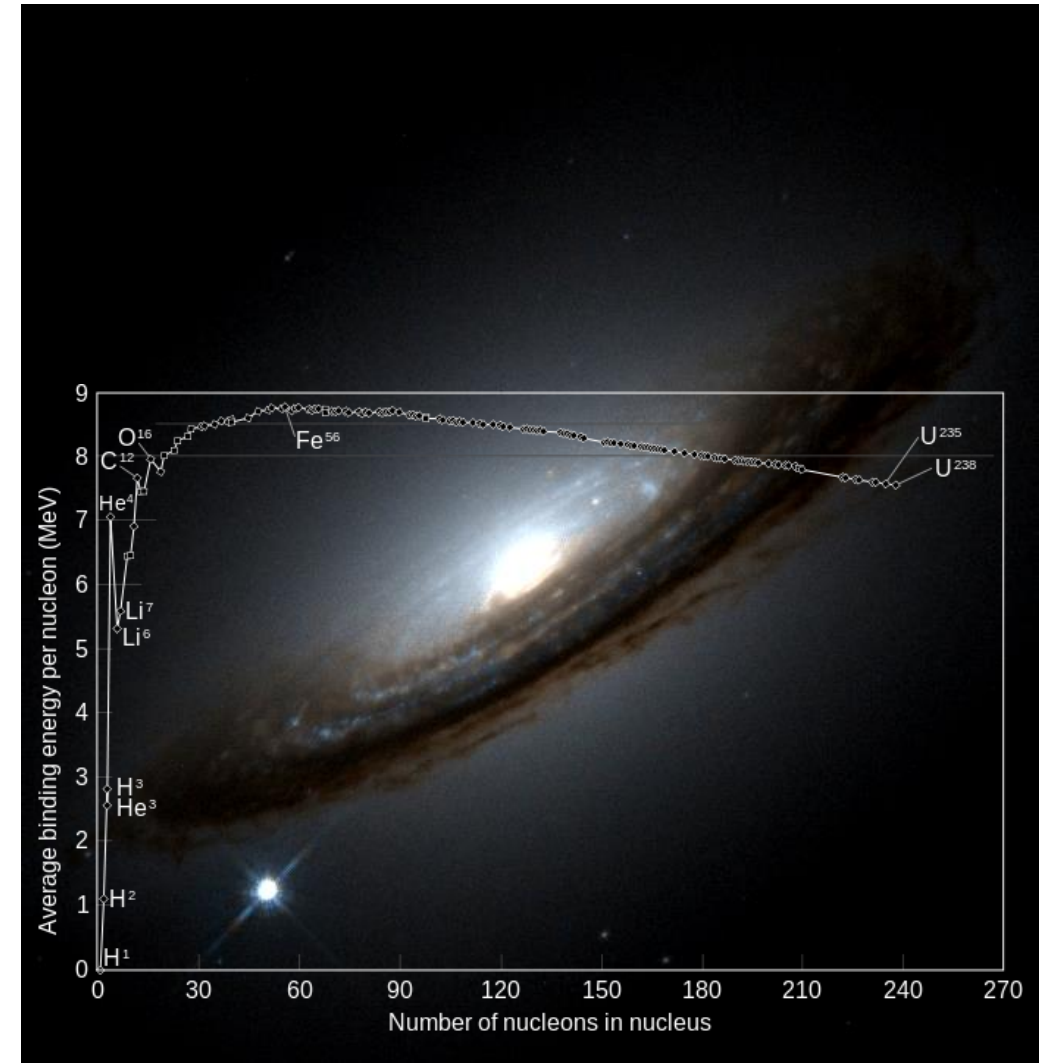
- → **Sufficient to power supernova!**
- What happens to majority of energy? → neutrinos!



Supernova energy source

Thermonuclear SNe

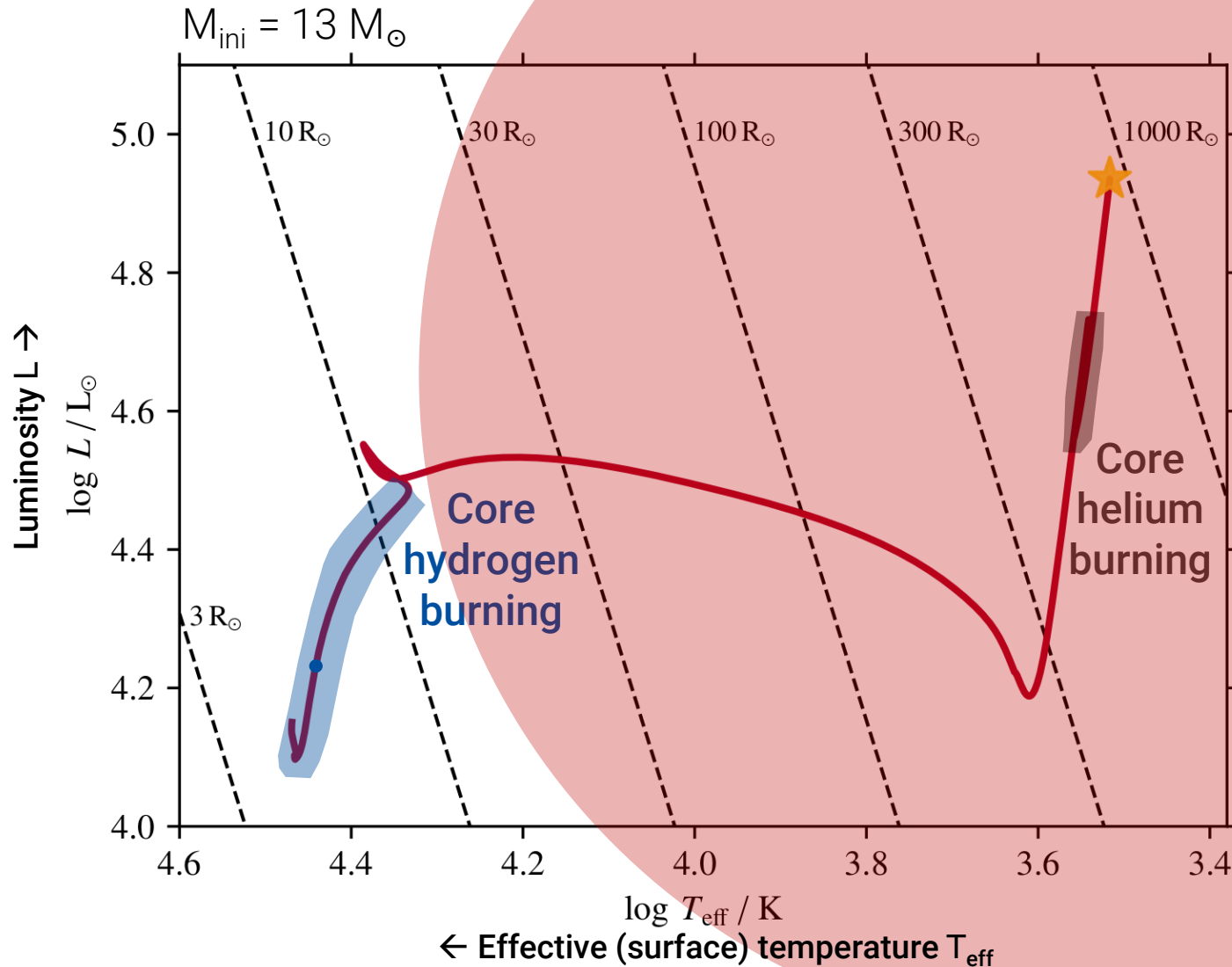
- **Observed kinetic energy: $\sim 10^{51}$ erg**
- **Energy source:** nuclear energy?
- no H, He in SNe Ia spectra
→ exploding star: C+O white dwarf
- Energy release due to burning of C+O material to ^{56}Ni :
 7.86×10^{17} erg/g
- Chandrasekhar-mass ($1.4 M_{\odot}$) WD: → 2×10^{51} erg
- → **Sufficient to power supernova!**



Stars 101:

Surface evolution

$1M_{\odot} = 1$ solar mass
 $1R_{\odot} = 1$ solar radius
 $1L_{\odot} = 1$ solar luminosity

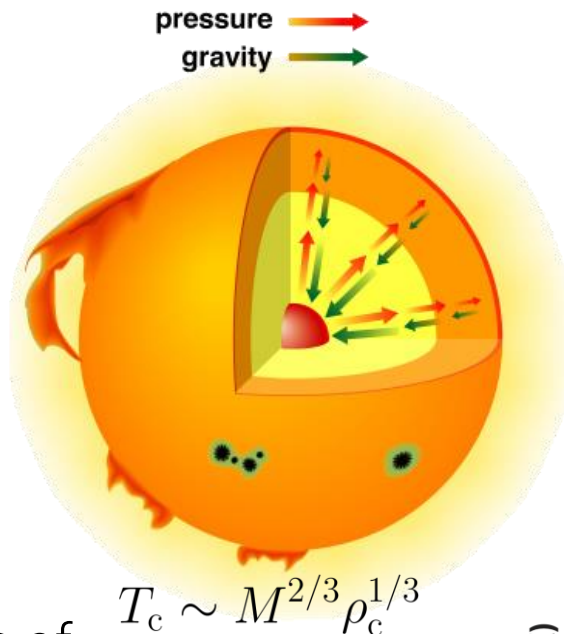


- Evolution of a $13 M_{\odot}$ star in **Hertzsprung–Russell diagram**
- Lines of constant radii from Stefan-Boltzmann law of black body

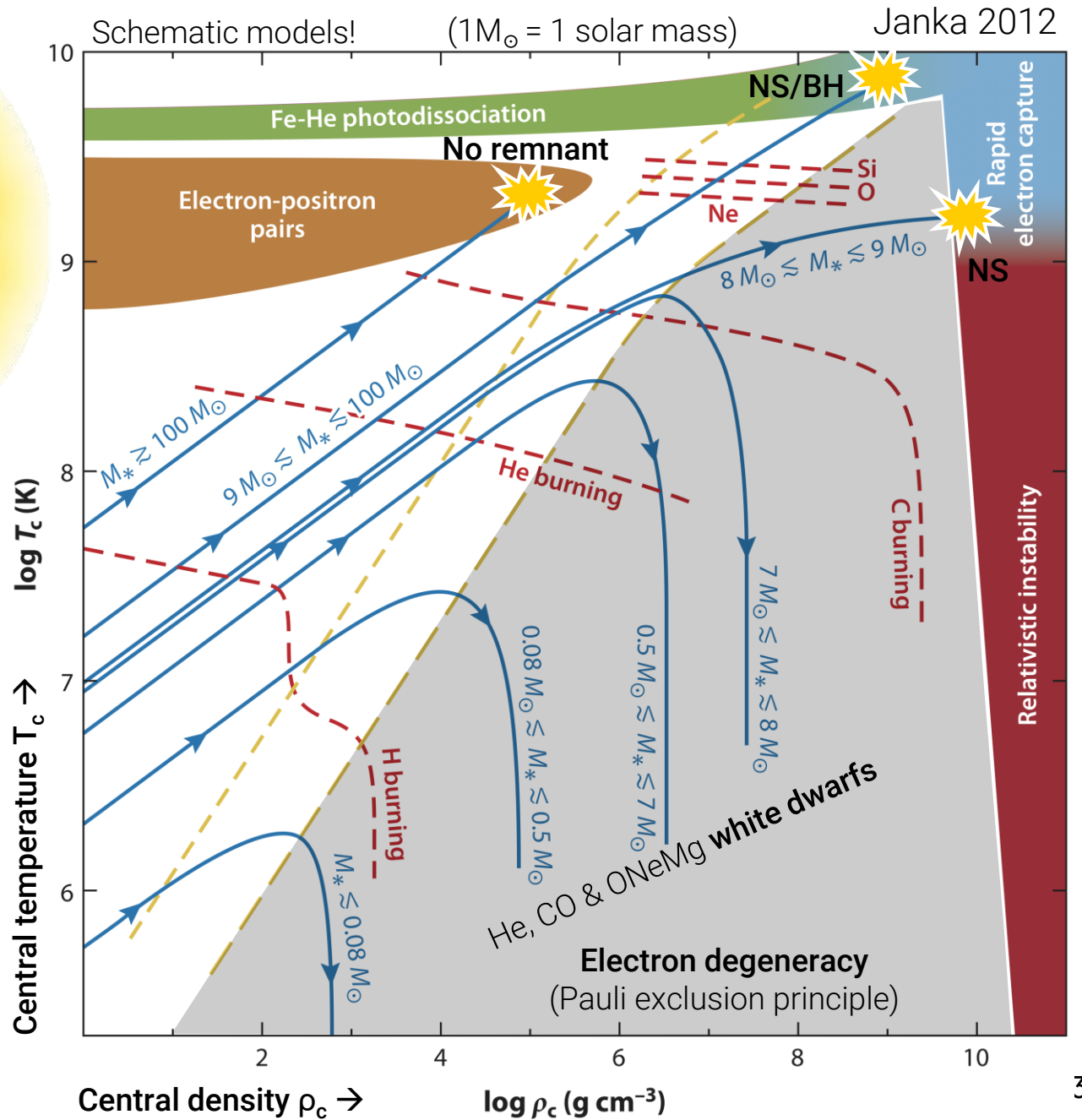
$$L = 4\pi R^2 \sigma T_{eff}^4$$
- While **core contracts**, the **entire star expands**
- Star becomes red supergiant and then explodes in supernova, leaving behind a neutron star
- More massive stars evolve faster

Stars 101: Core evolution

- Life of a star: eternal fight against gravity → Contraction
- Contraction halted temporarily by episodes of nuclear burning, e.g.
 - Hydrogen burning: $H \rightarrow He$
 - Helium burning: $He \rightarrow C, O$
- End of life: reaching death zones ($>8 M_{\odot}$; supernova, formation of neutron star or black hole) or gravity permanently balanced by degenerate electrons in white dwarfs ($<8 M_{\odot}$)



$$T_c \sim M^{2/3} \rho_c^{1/3}$$



Stars 101

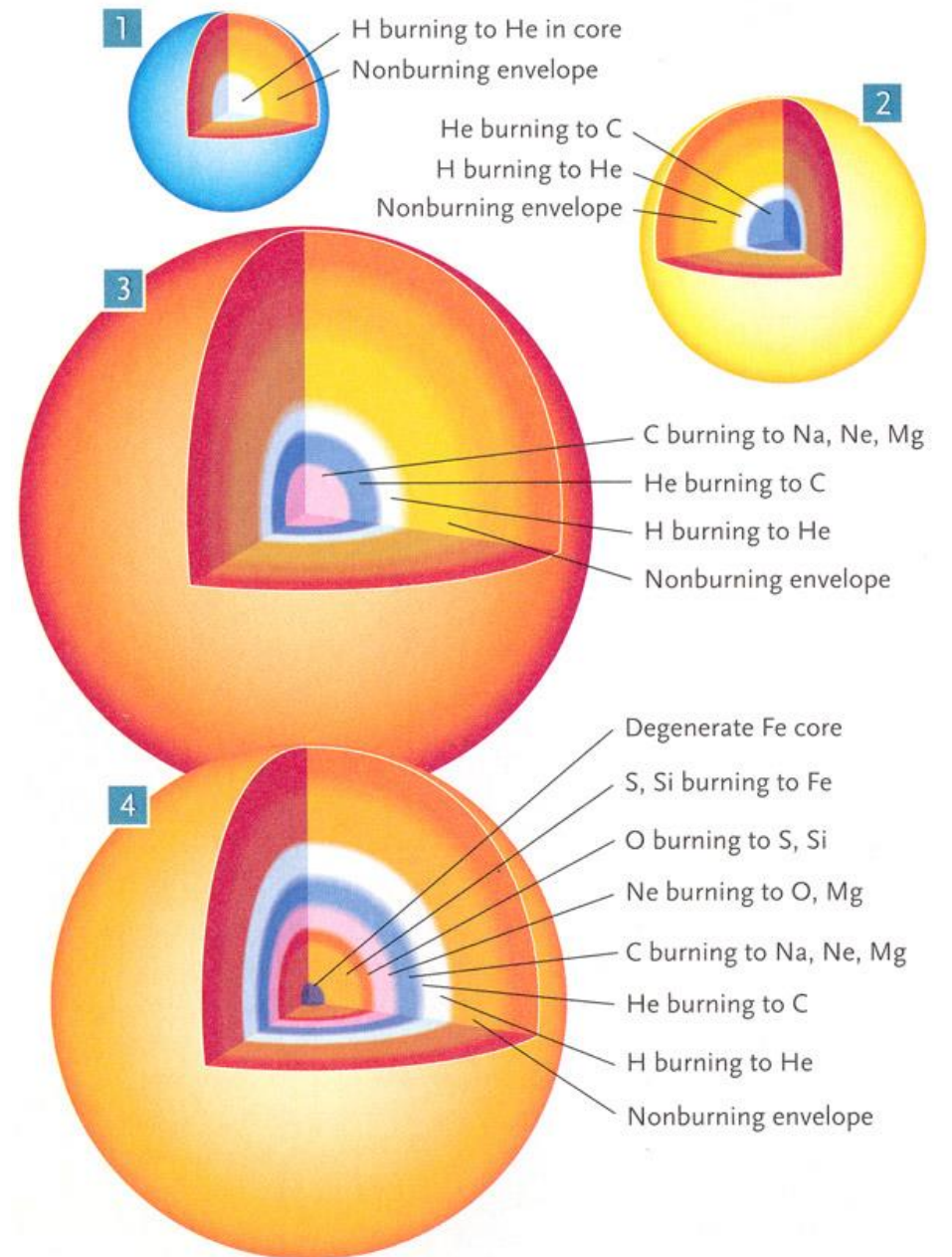
Chandrasekhar mass M_{Ch}

- Nuclear burning until iron core formation
→ nuclear fusion of iron-group elements impossible
→ only electron degeneracy left to balance gravity
- Once core mass exceeds effective **Chandrasekhar mass**, core starts to collapse (+ additional instabilities)

$$M_{Ch} \sim 1.43(2Y_e)^2 M_{\odot} \quad Y_e \equiv \frac{Z}{A}$$

- Two ways to reach M_{Ch} :
 - 1) add mass (shell burning, [accretion in SN Ia])
 - 2) reduce Y_e (electron capture reactions)

Image credit: Casey Reed

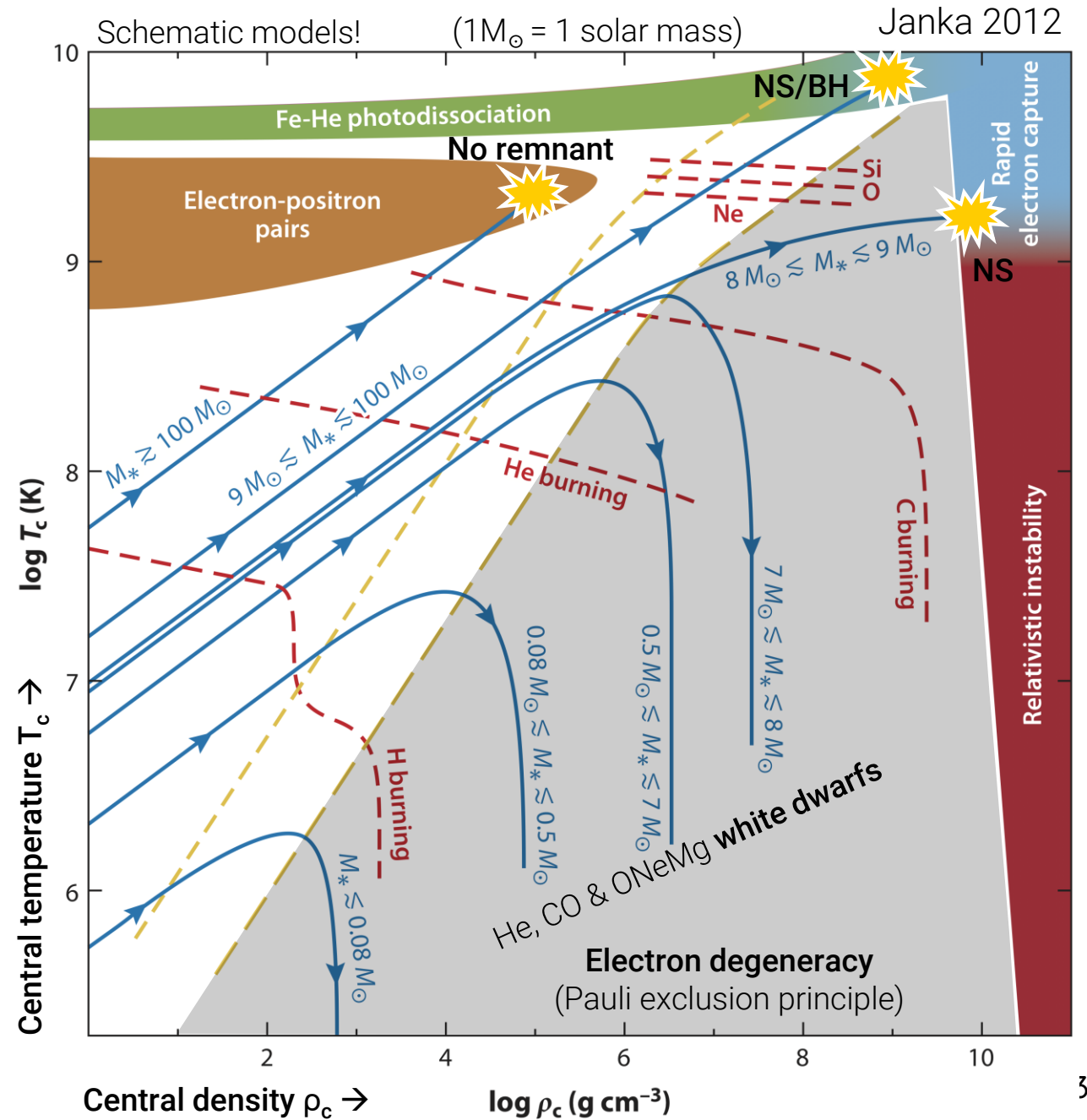


Three death zones

- Electron capture supernovae

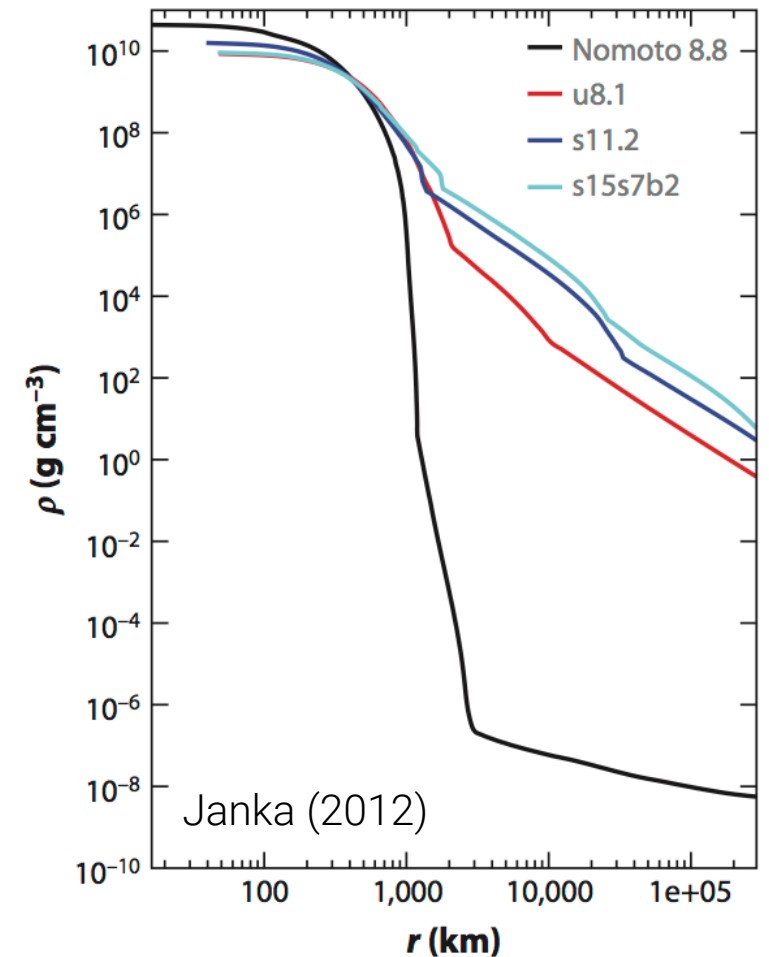
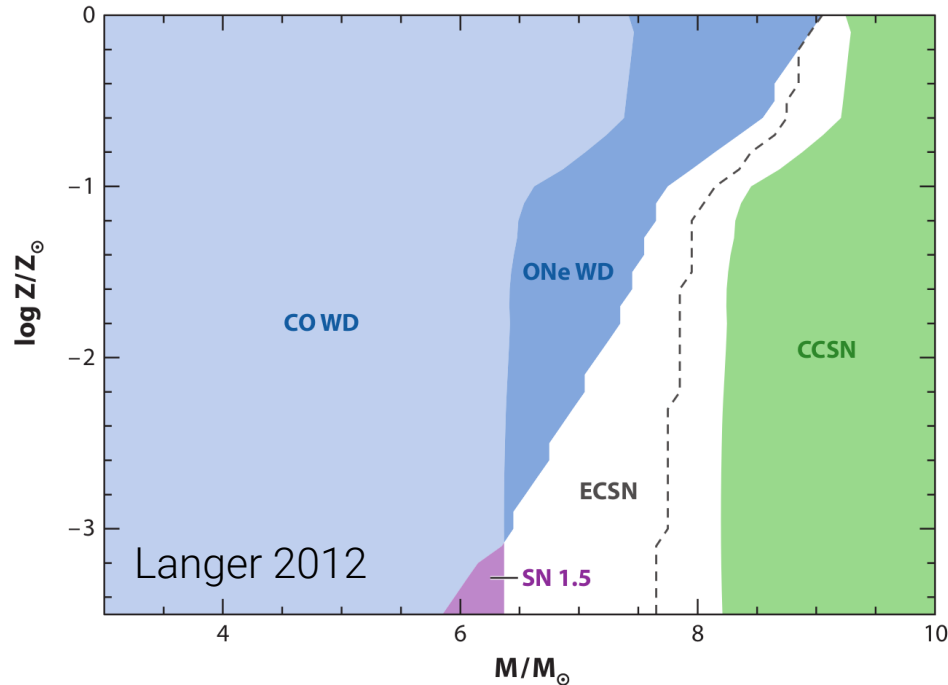
- Core collapse supernovae

- Pair instability supernovae
+ pulsational pair instability SNe



Electron capture supernovae

- applies to solar metallicity stars in range $\sim 8/9 \dots 10 M_{\odot}$
→ about 20 to 30% of all CCSNe may be ECSNe
- Core reaches electron-degeneracy before Ne burning
- Electron captures on ^{20}Ne and ^{24}Mg lower Y_e
→ M_{Ch} is reduced & core collapses
- Mass window can shift with other parameters (e.g. Z , binarity)



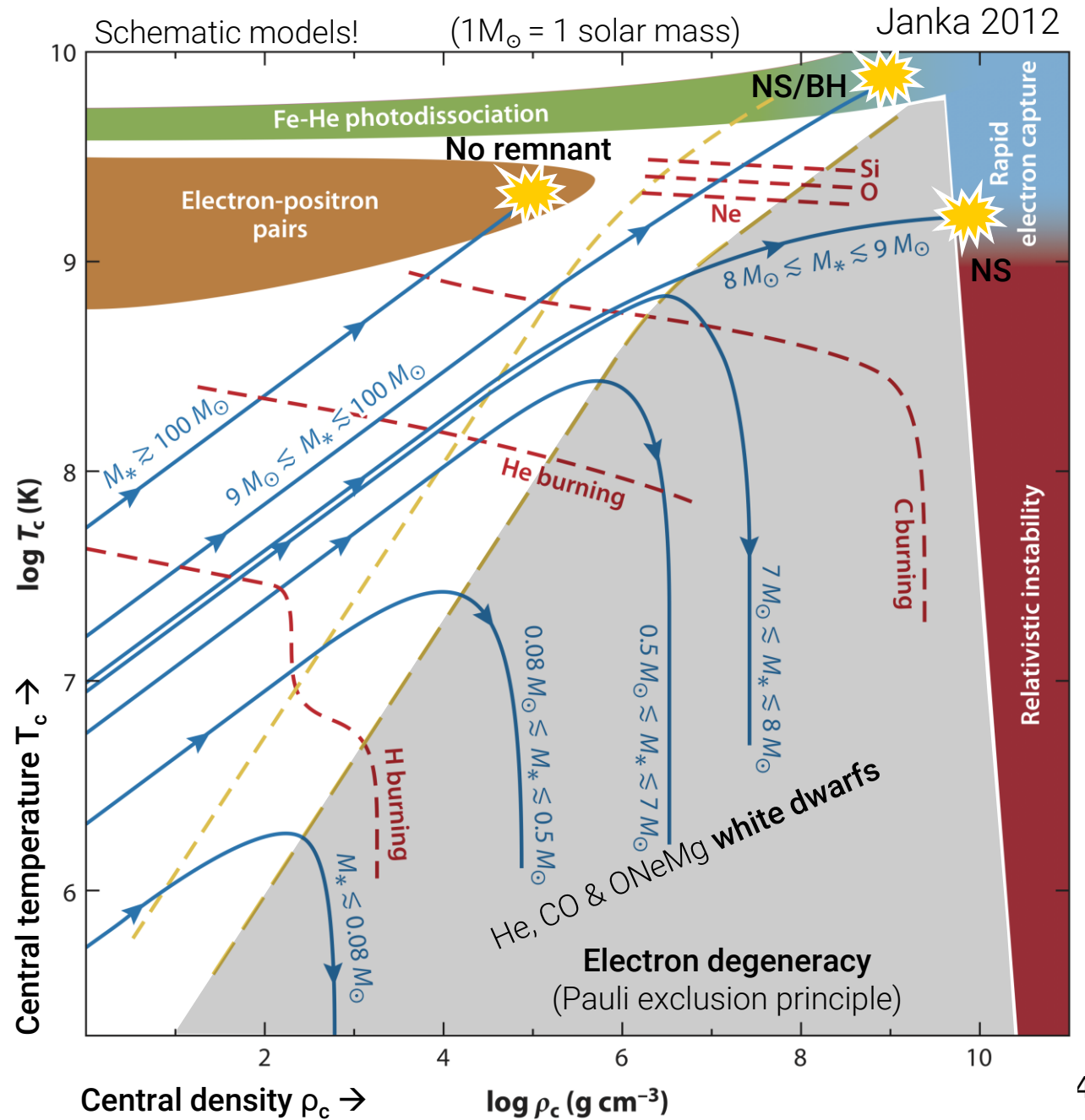
- Steep density decline at edge of O-Ne core:
→ special explosion characteristics compared to SNe from higher-mass stars (little mass ejection, very little Ni → faint, **smaller explosion energies $\sim 10^{50}$ erg**)

Three death zones

- Electron capture supernovae

- Core collapse supernovae

- Pair instability supernovae
+ pulsational pair instability SNe



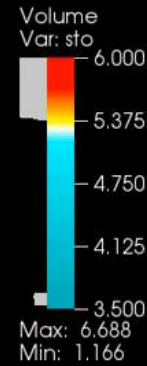
Core collapse supernovae

The “gravity bomb”

10 ms

Melson, Janka et al. 2015

- **Core collapse** triggered if **iron core mass** > **effective Chandrasekhar mass**
- Energy source: **gravitational energy**
→ **~99% released in neutrinos (SN1987A)**
→ **how to use neutrinos to explode star?**
- **Core bounce** and **shock formation** at nuclear densities ($\rho \sim 10^{14} \text{ g cm}^{-3}$)
- **Shock stalls** (→ dissociation of nuclei)
- **Proto-NS forms** and the **released neutrinos** **heat gain region behind shock** → **convection**
- **Threshold process**: if enough energy deposited, shock is revived, and **star explodes**
- **Simulations**: consistent explosions among different groups, yet too little energy (e.g. O'Connor+2018, Burrows & Vartanyan 2021)



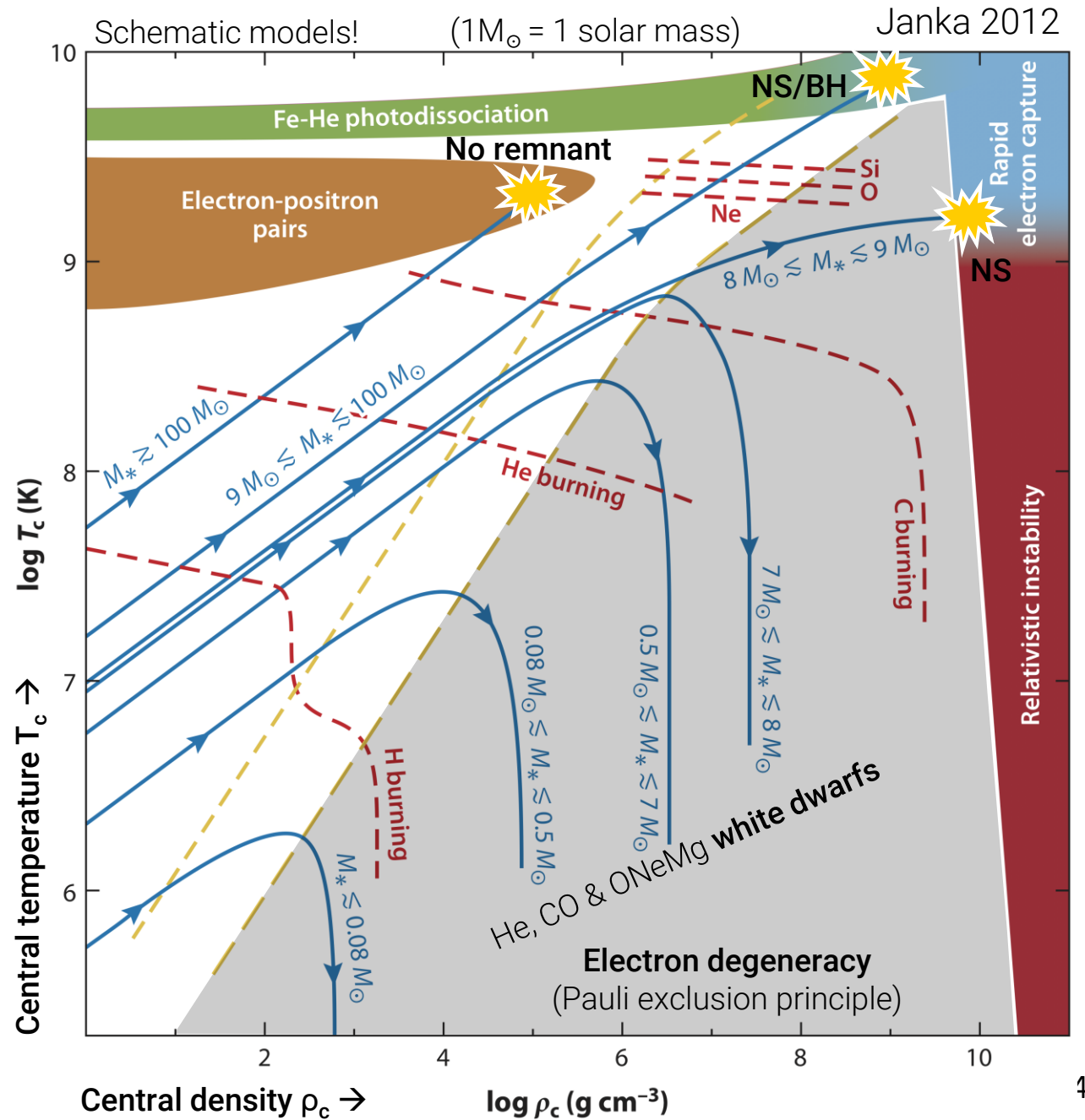
192 km

Three death zones

- Electron capture supernovae

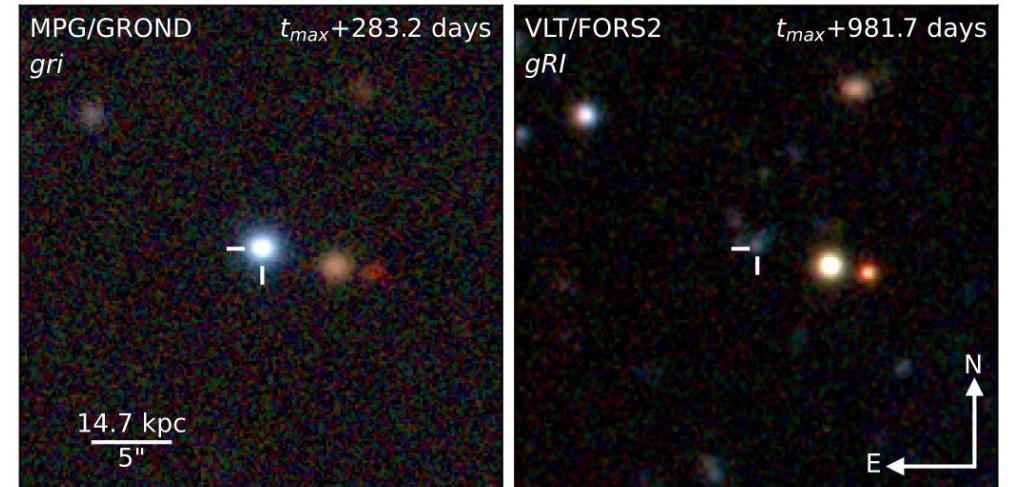
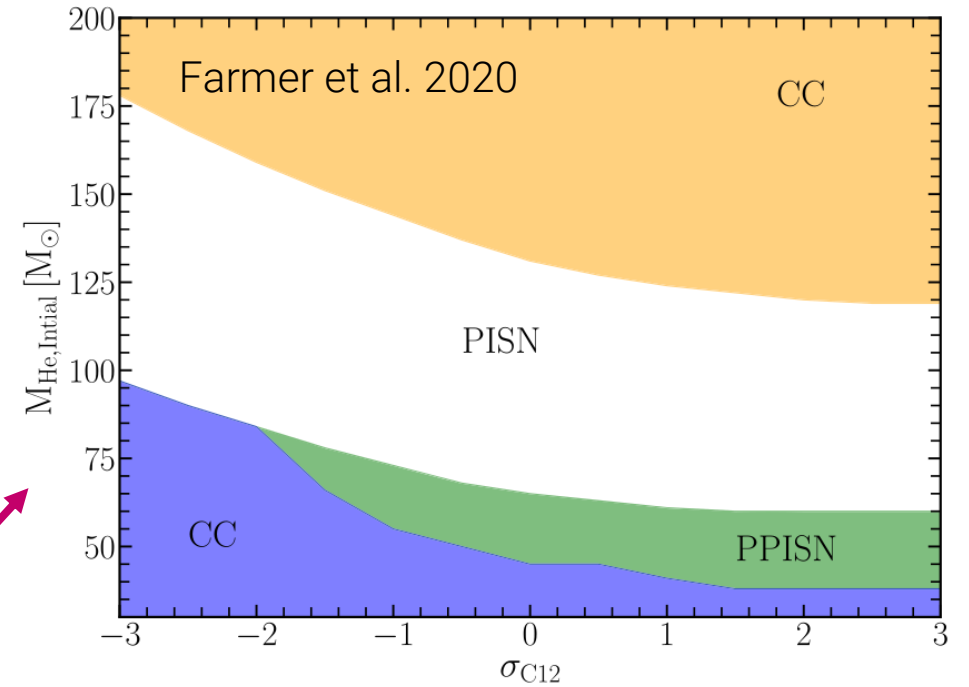
- Core collapse supernovae

- Pair instability supernovae
+ pulsational pair instability SNe



Pair instability supernovae

- Stars with initial mass above $\sim 100 M_{\odot}$ \rightarrow very hot
- Pair instability after C burning** at $T \sim 1 \text{ GK}$: formation of e^+e^- -pairs from high-energy photons converts thermal energy into rest-mass energy \rightarrow reduces adiabatic index of E.o.S. below $4/3$
- Thermonuclear explosion** of left nuclear fuel \rightarrow enough to completely disrupt the star? \rightarrow mass-dep.
- Mass range uncertain, e.g., nuclear reaction rates
- Occur only at $Z < Z_{\odot} / 10$ (Langer 2007)
- Explosion energy: $\sim 10^{53}$ erg**, can produce more than $50 M_{\odot}$ of ^{56}Ni , but most events will produce “usual” CCSN Ni masses
- Best candidate so far: SN 2018ibb (Schulze+2023)



SN 2023ibb; Schulze et al. 2023

Stars 101: How massive single stars end their life

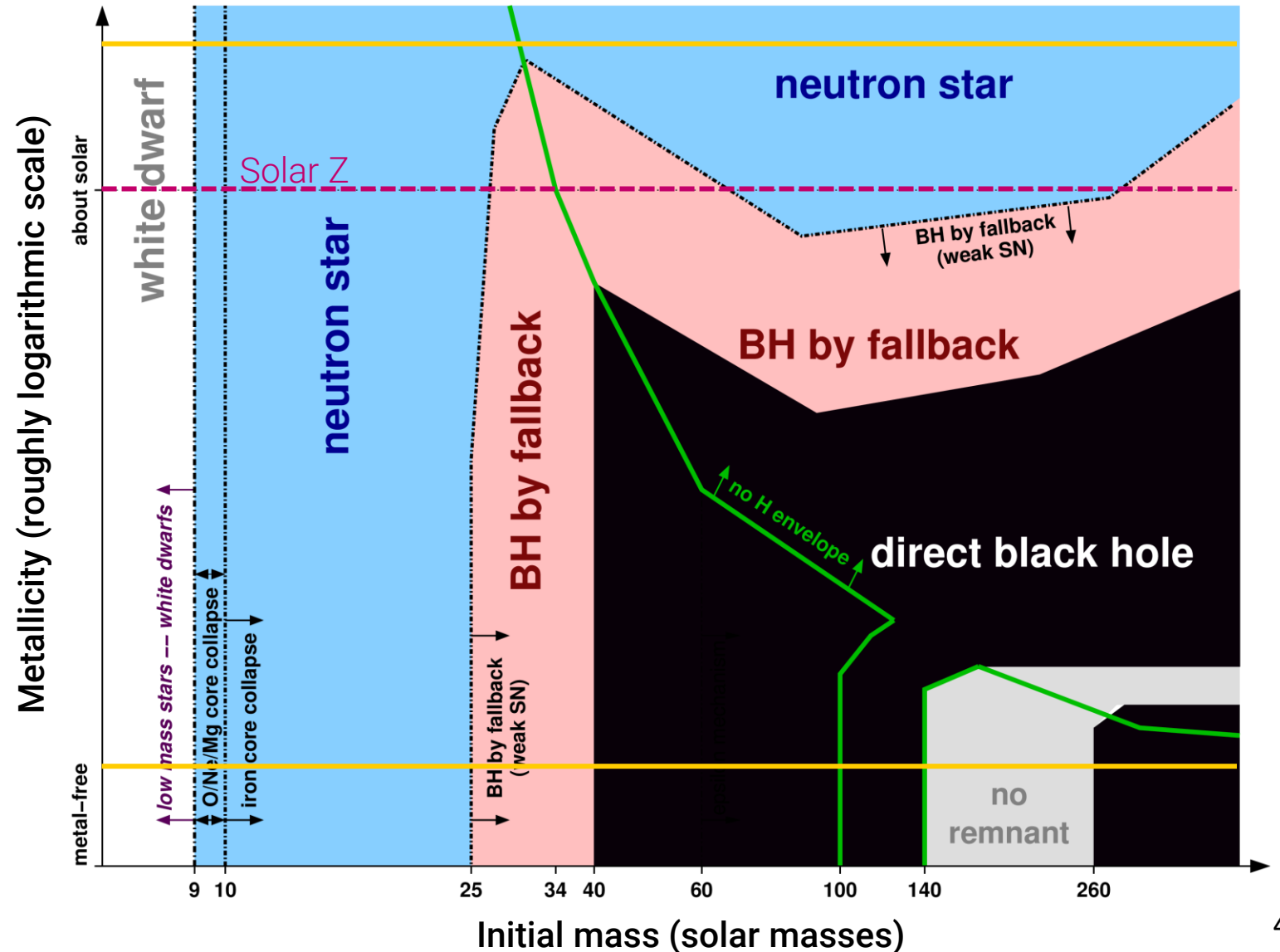
What we teach in lectures:

- $8/9-10 M_{\odot}$: electron capture supernova (core collapse)
- $10-25 M_{\odot}$: iron core collapse supernova → **neutron star**
- $25-40 M_{\odot}$: weak core collapse supernova → **black hole** by fallback
- $>40 M_{\odot}$: no supernova → collapse to **black hole**

Special cases: → stellar winds

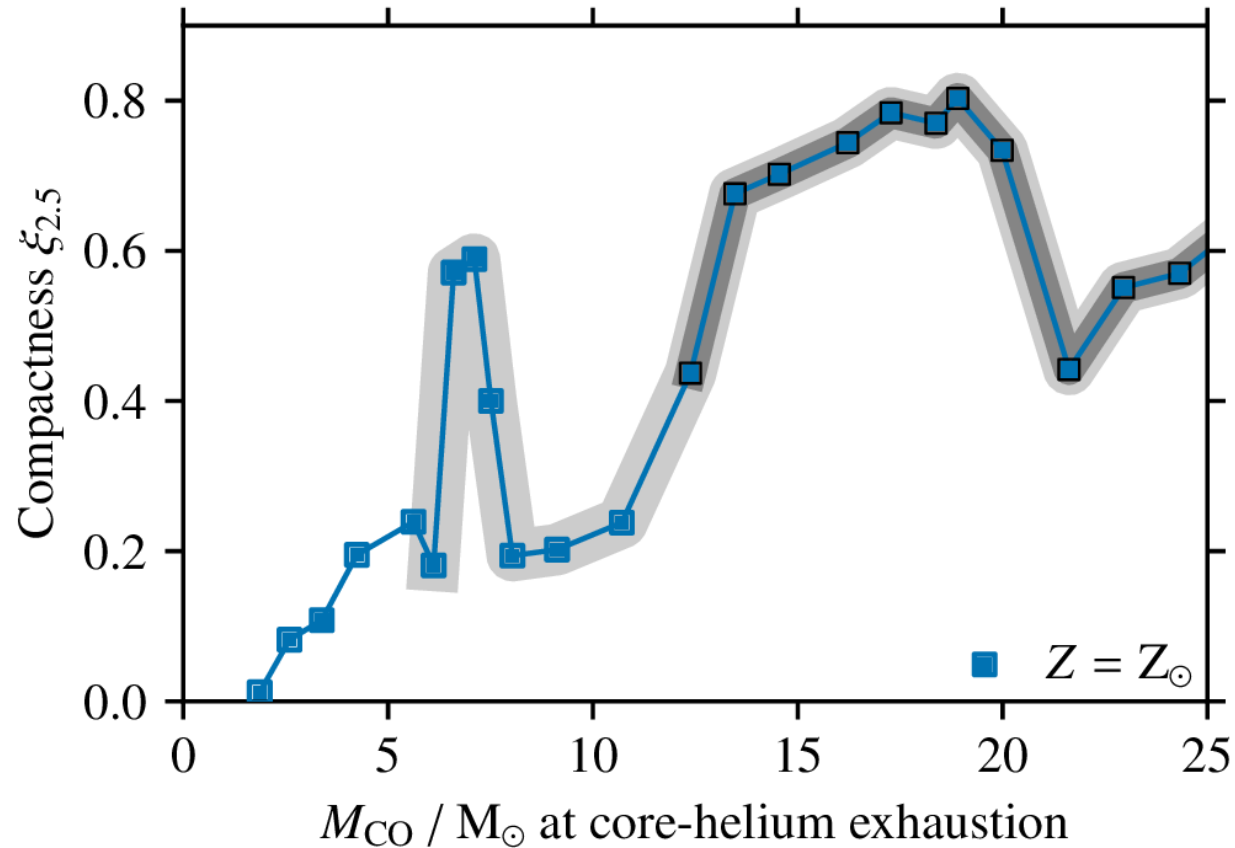
- **High-Z**: no black holes at all
- **Low-Z**: pair-instability supernovae, no remnant → **PISN BH gap**: $\sim 45-120 M_{\odot}$

Heger et al. 2003; Heger, Müller & Mandel 2023



Final fate of (single) stars (CCSNe)

A closer look at the pre-SN stellar structure



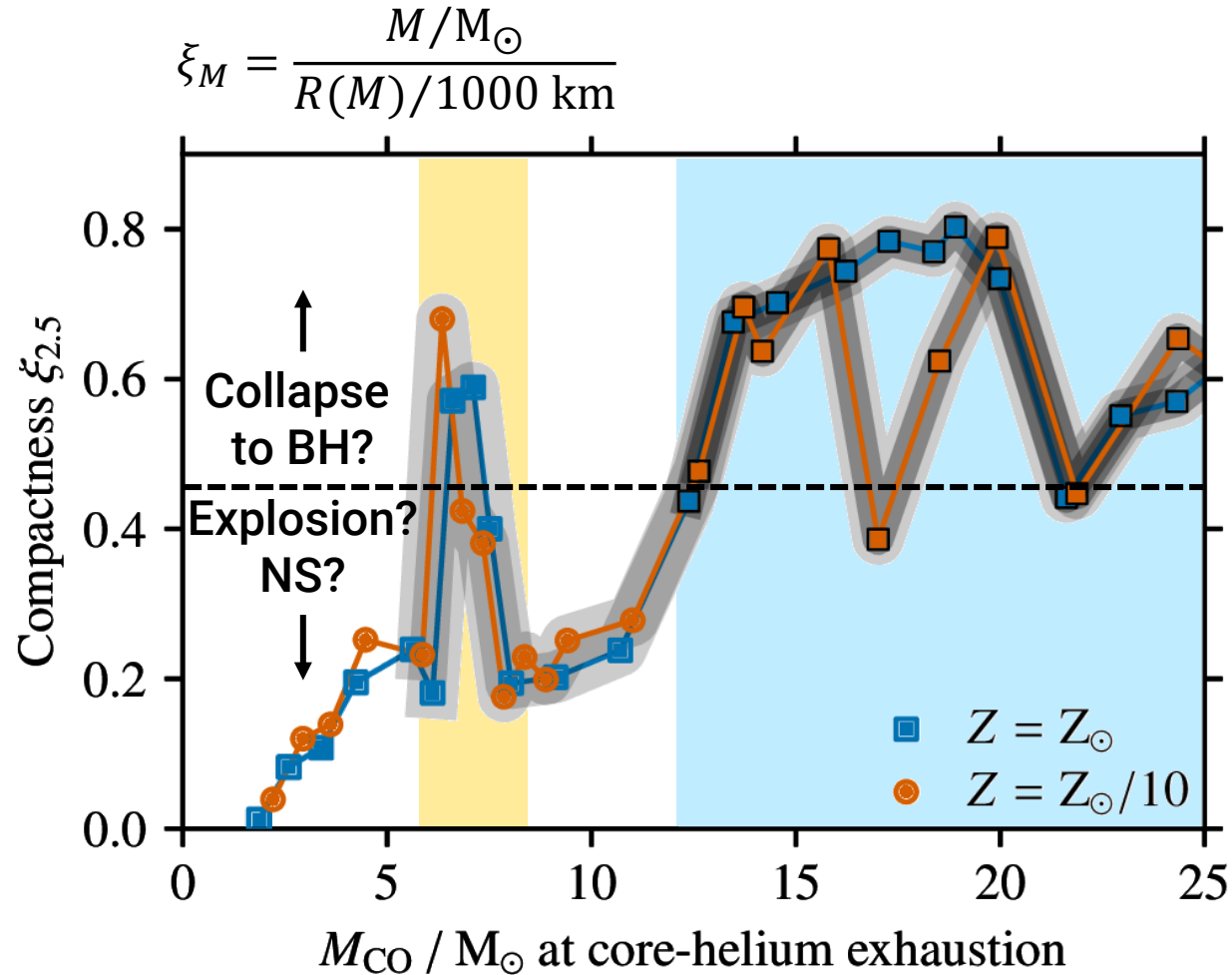
- Pre-SN stellar structure summarised by compactness parameter ξ_M

$$\xi_M = \frac{M/M_{\odot}}{R(M)/1000 \text{ km}}$$

- Non-monotonic landscape** with core mass of stars (here CO core mass M_{CO})
- Stars with large ξ_M difficult to explode

Final fate of (single) stars (CCSNe)

A closer look at the pre-SN stellar structure



- **Compactness as explosion criterion:**
 - $\xi_{2.5} > 0.45 \rightarrow$ Collapse to BH?
 - $\xi_{2.5} < 0.45 \rightarrow$ Explosion and NS formation?

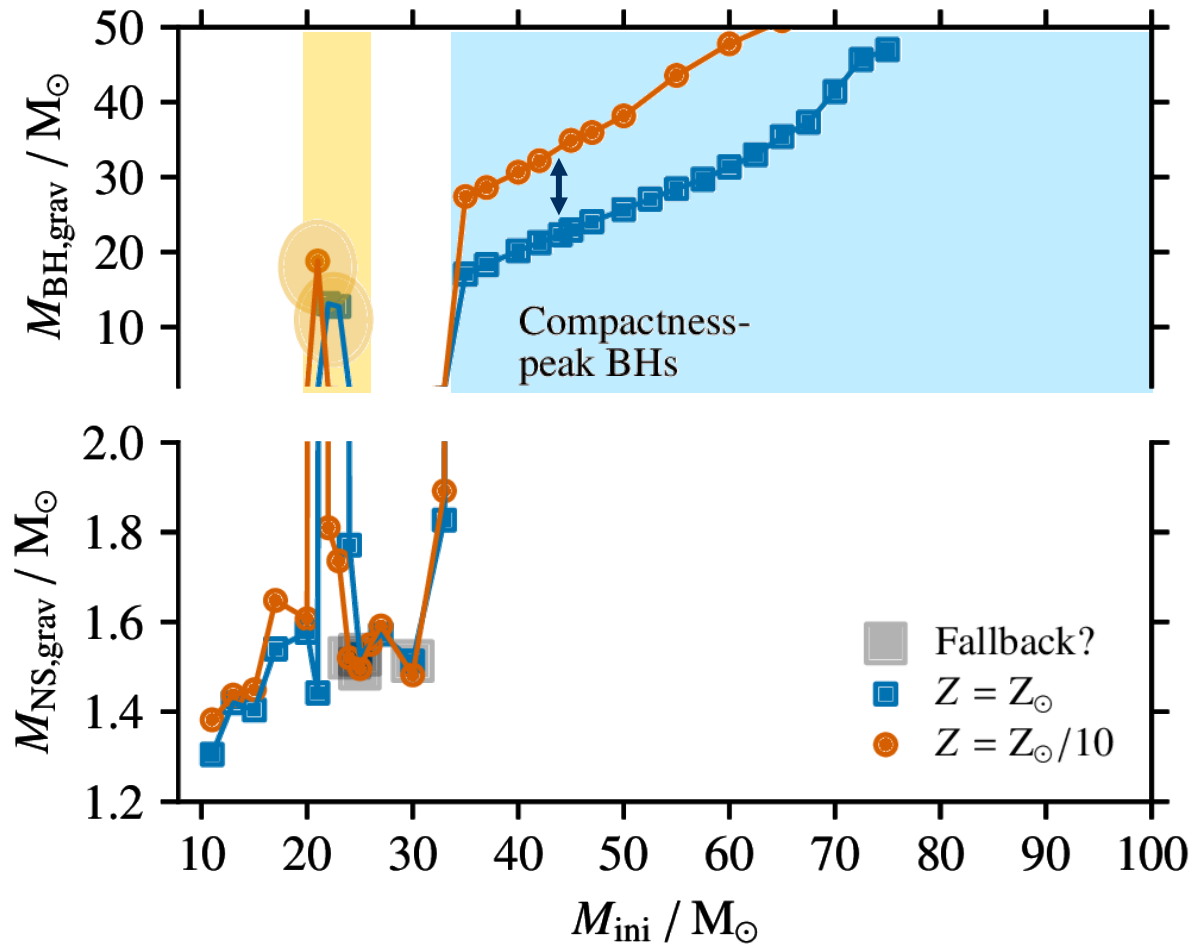
- **High compactness at**
 - $\rightarrow M_{\text{CO}} \sim 7.5 M_\odot$ (“compactness peak”) and
 - $\rightarrow M_{\text{CO}} > 13 M_\odot$

- Driven by **neutrino-dominated carbon and neon burning**

- **Compactness pattern not strongly metallicity dependent**
(unless winds erode helium and CO cores)

Final fate of (single) stars (CCSNe)

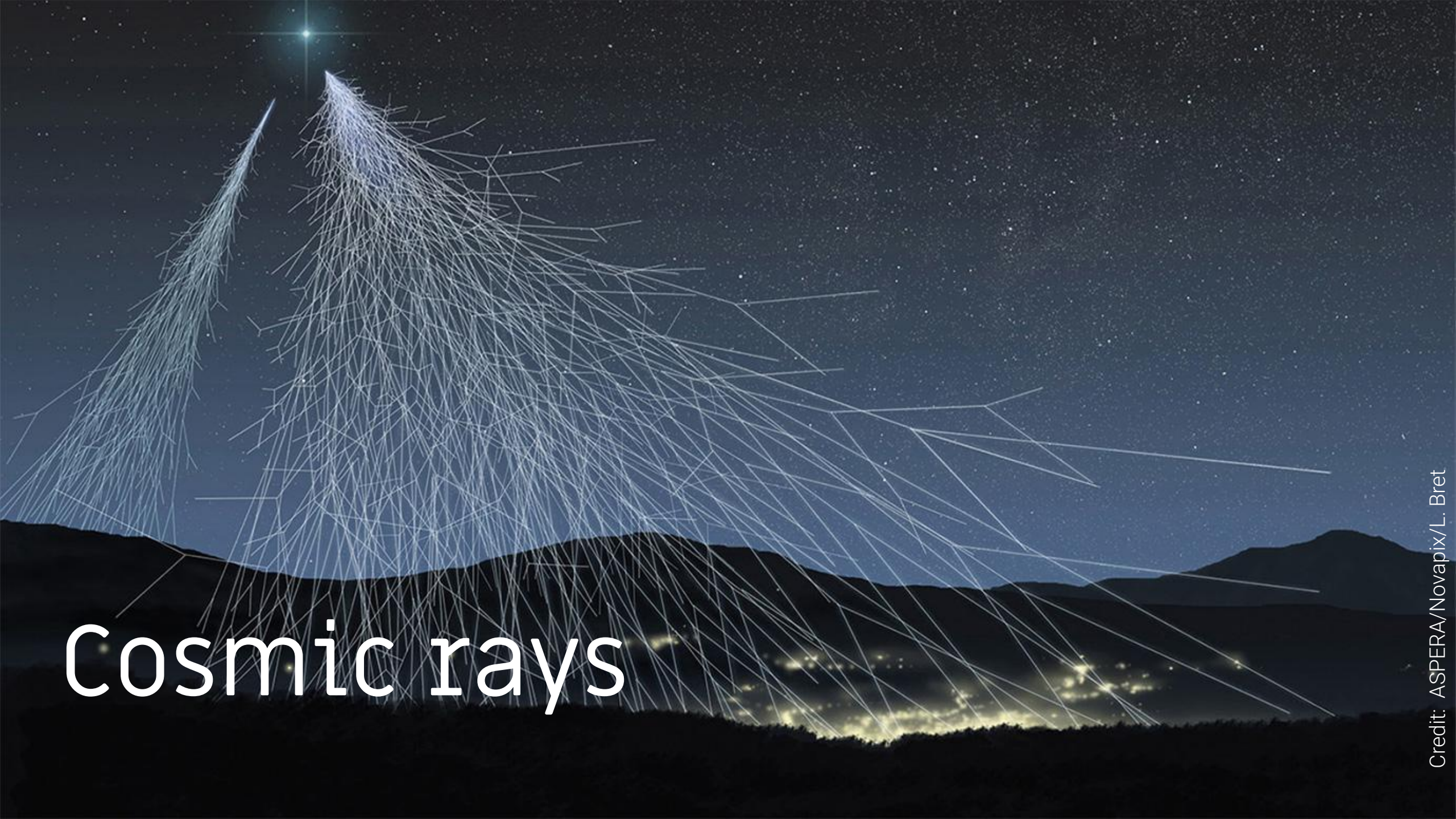
Compact remnant masses



- Models exploded with **semi-analytic, v-driven SN model of Müller et al. 2016**
- Explodability largely follows compactness (cf. O'Connor & Ott 2011, Ugliano+2012, Sukhbold & Woosley 2014, Ertl+2016, Müller+2016, ...)

- BH formation at**
 - $M_{\text{ini}} \sim 20-25 M_{\odot}$ ("compactness peak")
 - $M_{\text{ini}} \sim 25-35 M_{\odot}$: some fallback BHs
 - $M_{\text{ini}} > 35 M_{\odot}$

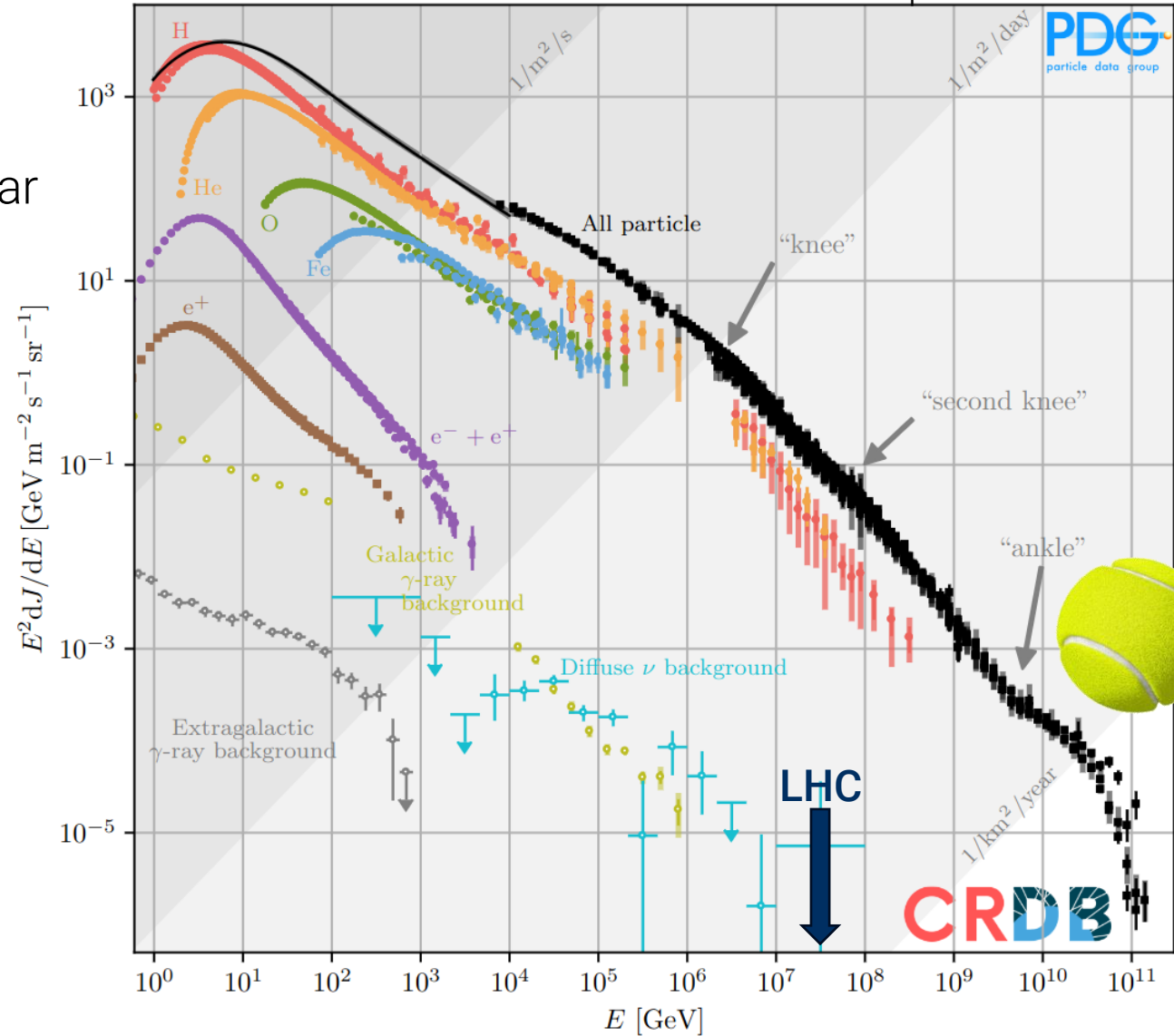
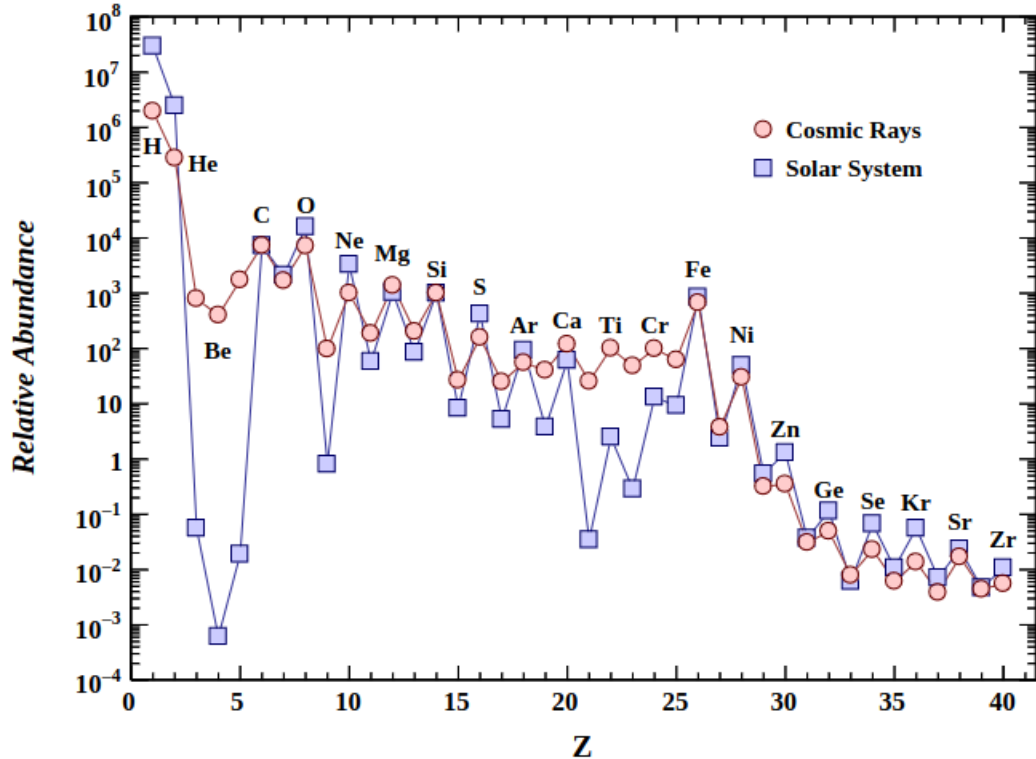
- At lower $Z = Z_{\odot}/10$, essentially **same landscape**, but **BHs more massive**
 - weaker winds, hence less mass loss
- Warning: highest BH masses unrealistic because of missing enhanced mass loss from LBVs etc.



Cosmic rays

What are cosmic rays (CRs)?

- **Relativistic particles** originating outside the solar system (Galactic & extra-galactic)
 → ~99% atomic nuclei (90% p or H), ~1% e^-



Where do cosmic rays come from?

Credit: NASA, ESA, CSA, STScI, Danny Milisavljevic, Ilse De Looze, Tea Temim

- Mostly: **particles accelerated in SN shock fronts**
 - up to knee (protons) and beyond (heavier atoms)
 - **diffusive shock acceleration** (Fermi acceleration)
 - highest energies: extra-galactic (AGNs?)
- Particles repeatedly cross shock front and each time gain energy (Krymskii 1977, Axford+1978, Bell 1978, Blandford & Ostriker 1978)

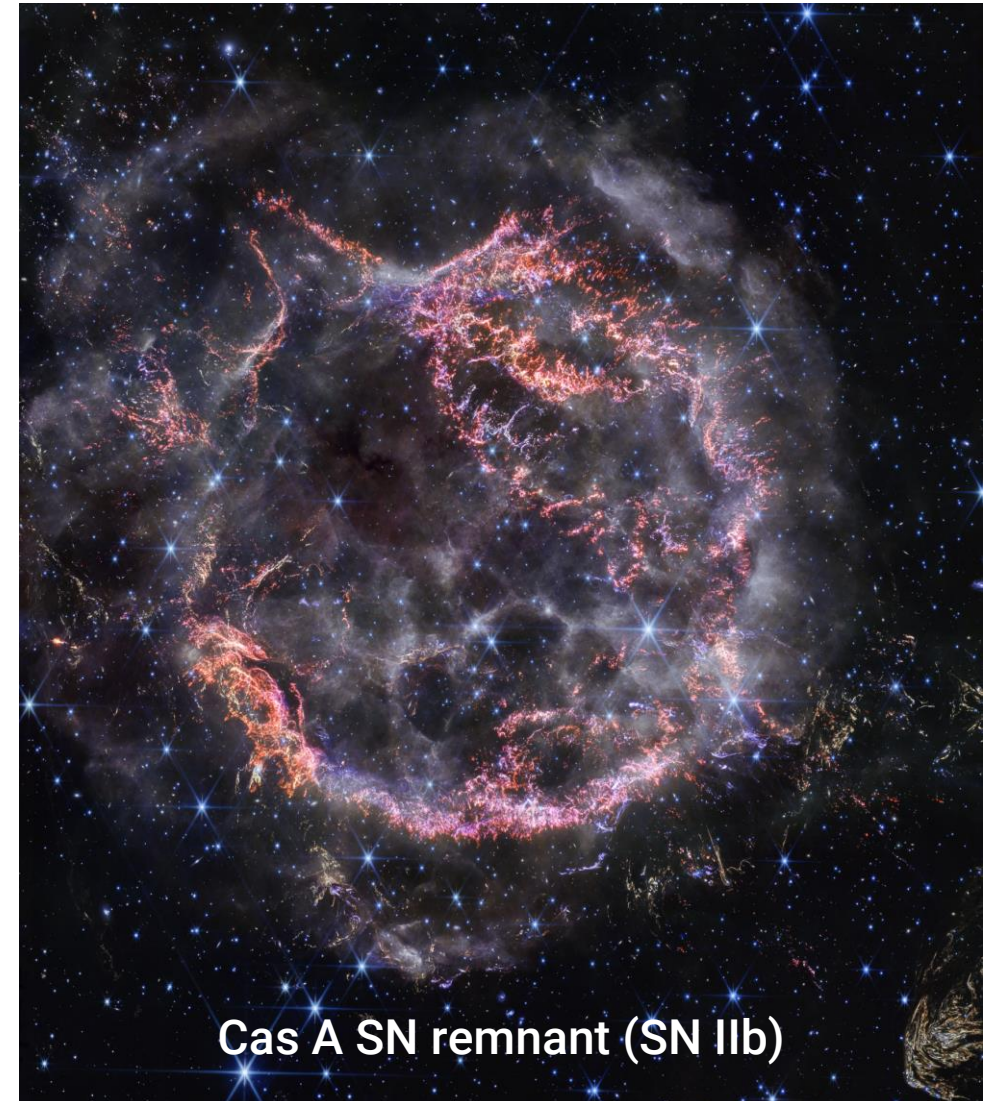
$$\frac{\Delta\epsilon}{\epsilon} = \frac{v_{\text{shock}}}{c}$$

- Results in power-law energy distribution
- Fits energetics in MW: CRs diffusive escape losses

$$L_{\text{CR}} = \frac{V_{\text{MW}}\epsilon_{\text{CR}}}{\tau_{\text{esc}}} \sim 3 \times 10^{40} \text{ erg s}^{-1}$$

- Can transfer ~10% of SN shock energy to CRs

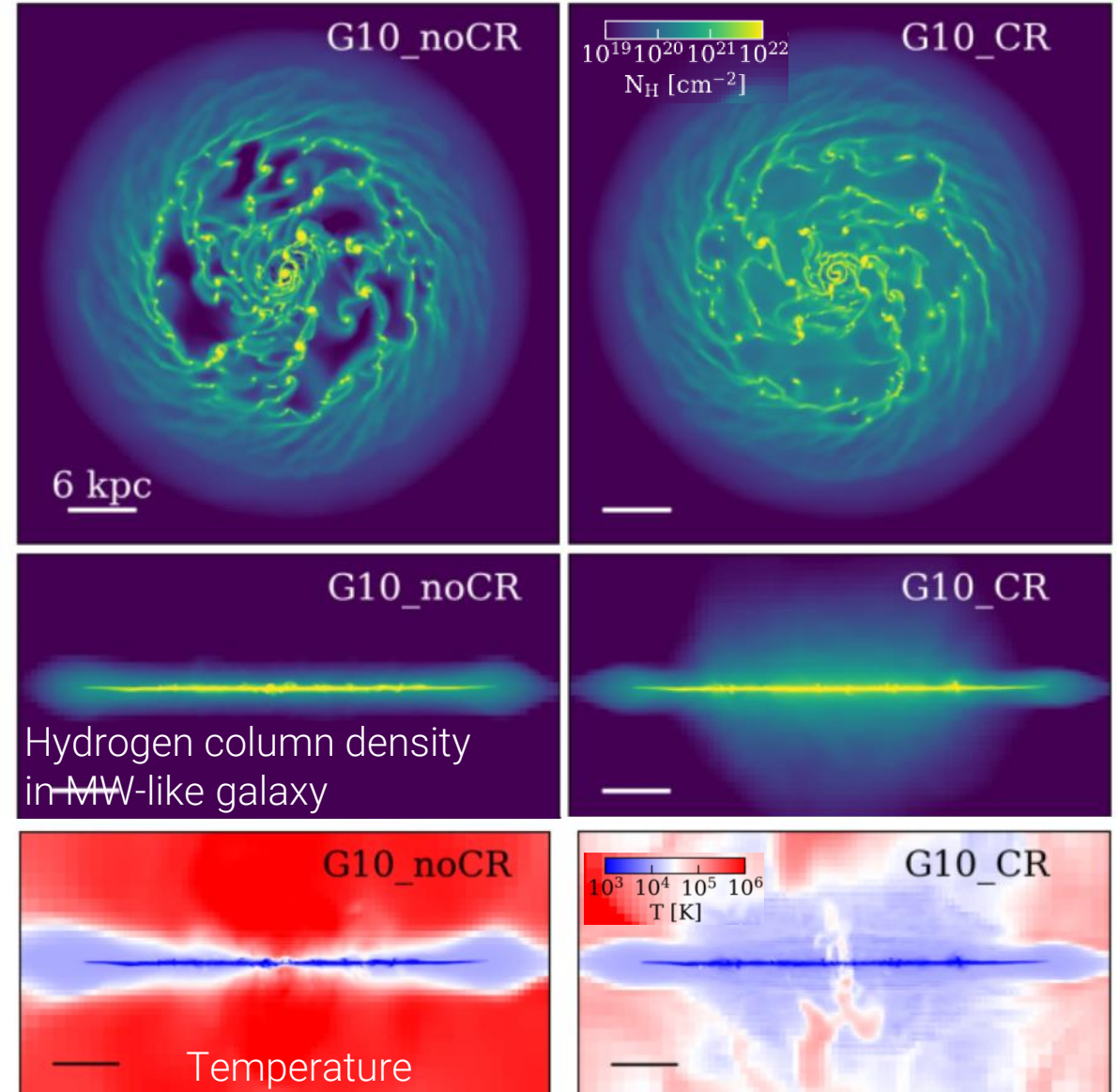
$$L_{\text{CR}} \sim 0.1L_{\text{SN}} \sim \frac{0.1E_{\text{SN}}}{\tau_{\text{SN}}} \sim \frac{0.1 \times 10^{51} \text{ erg}}{100 \text{ yr}} \sim 3 \times 10^{40} \text{ erg s}^{-1}$$



What is the feedback of cosmic rays?

Farcy et al. 2022

- Review: **Ruszkowski & Pfrommer 2023**
- CRs meander along B-fields of galaxies or escape galaxy depending on energy
- CRs ionization (most have few GeV & can penetrate optically thick regions)
- CR-driven galactic winds → CR pressure gradient
- CR feedback reduces SF rate
→ helps solve missing satellites problem?
- CRs lead to cooler ISM and galactic outflows
- Nucleosynthesis: production of Li, Be and B



Binary stars



Binary stars

- **Most massive stars are in binaries** (Sana+2012, *Science*)
→ huge consequences for evolution and final fate

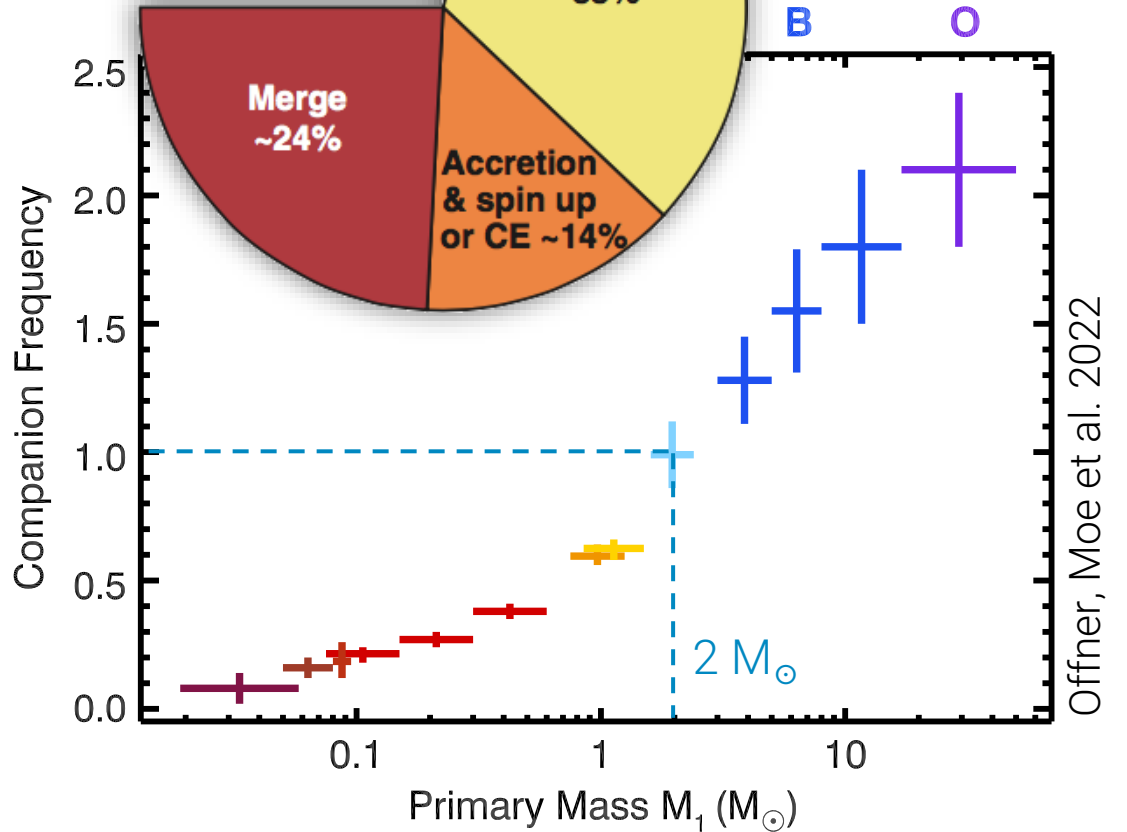
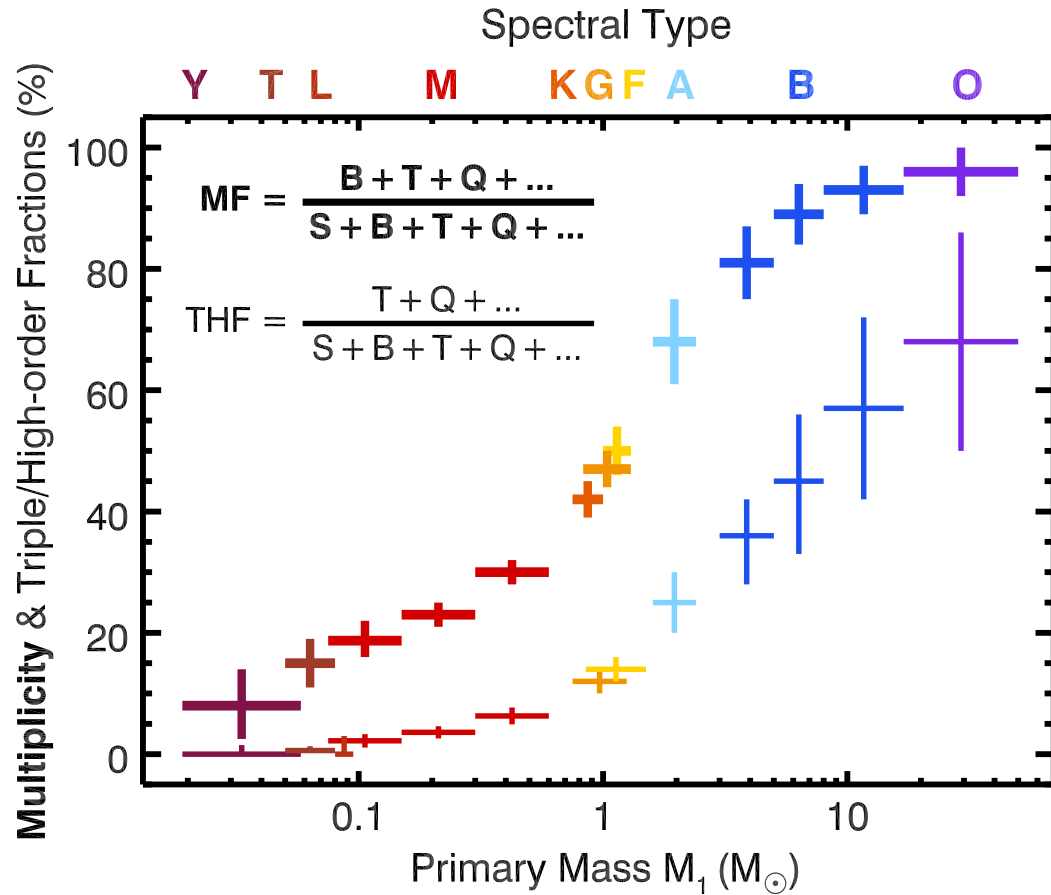


Binary stars

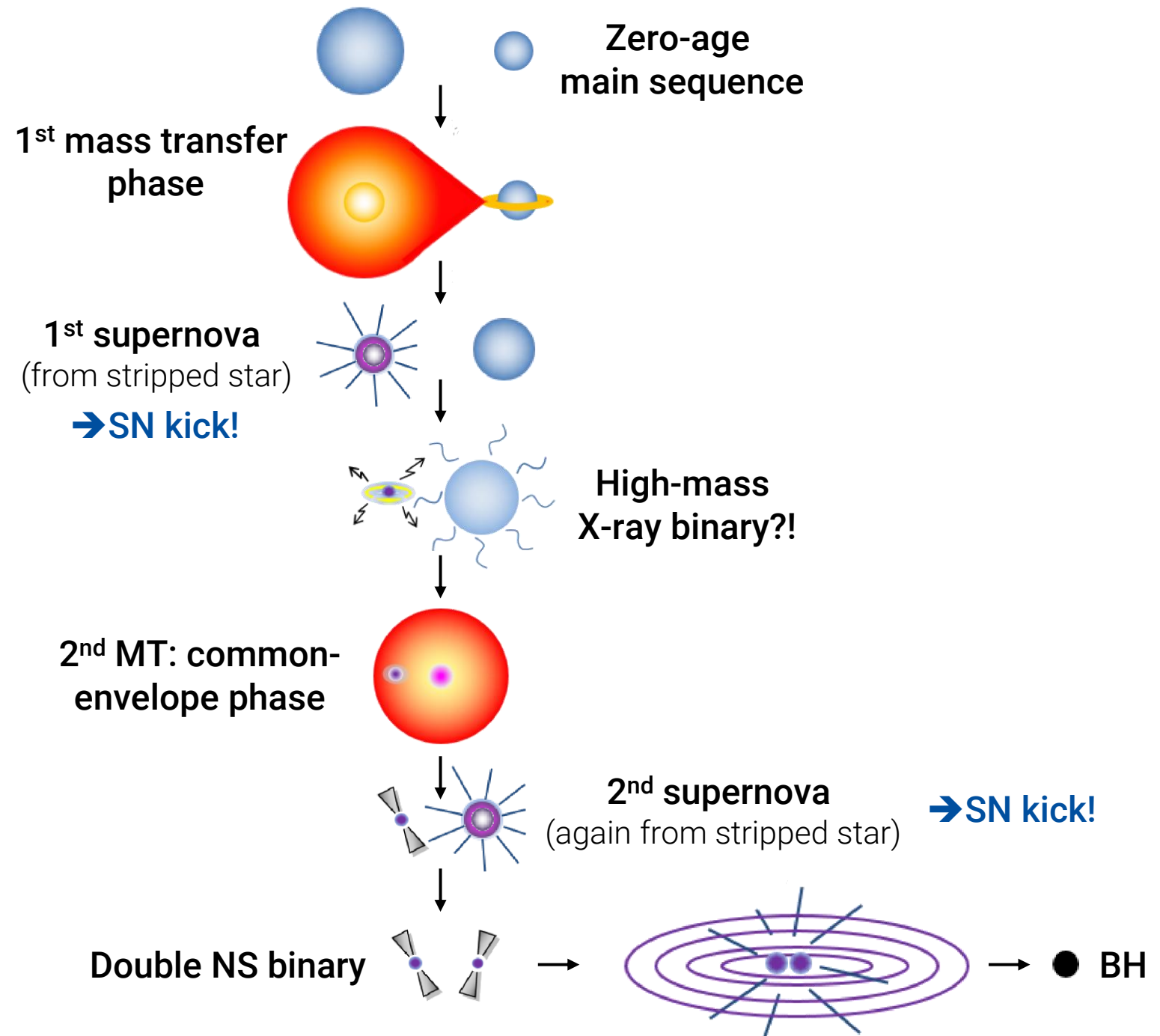
Most massive stars exchange mass with companion during life

Sana et al. 2012

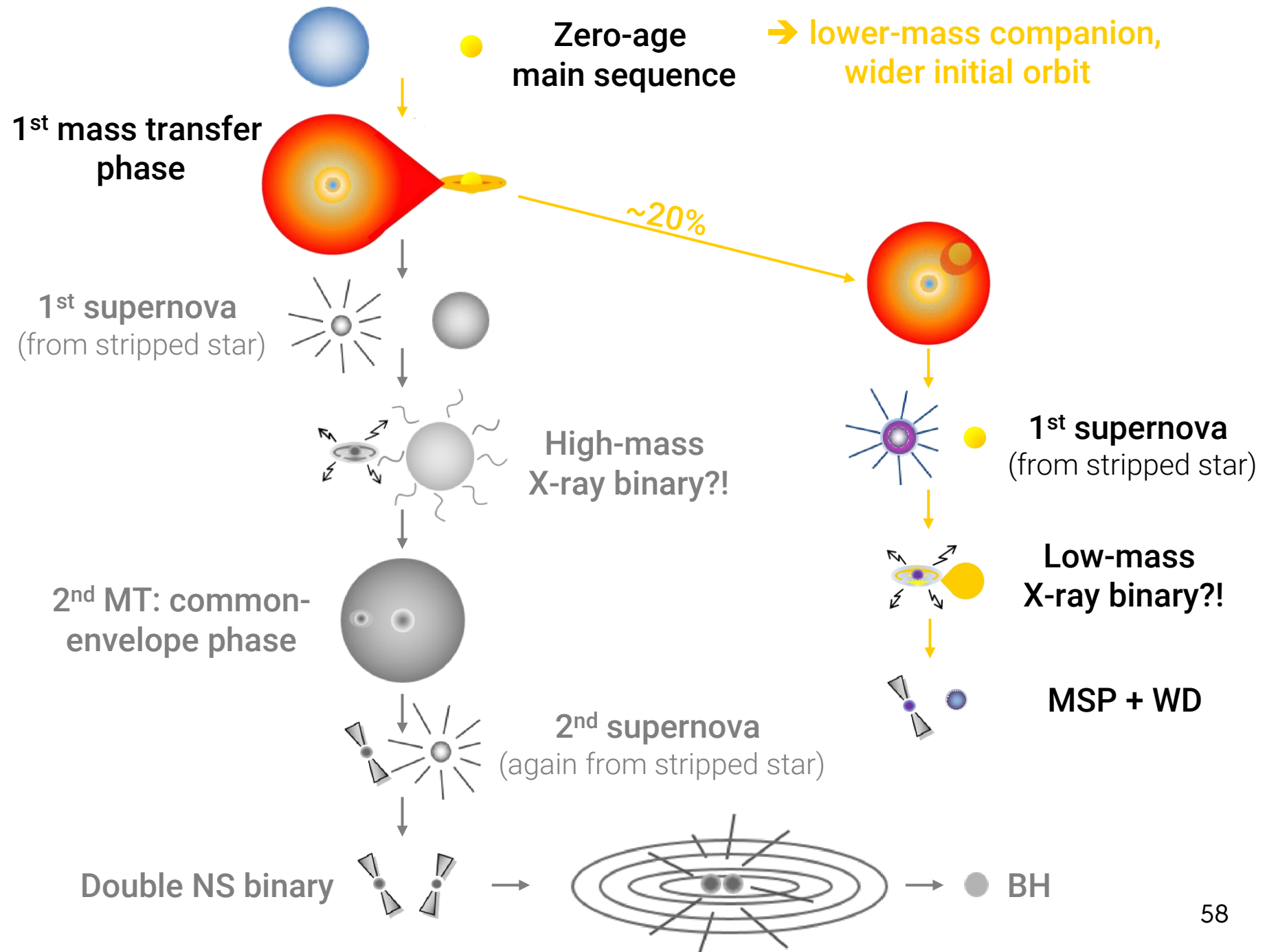
→ Most supernovae and BHs from binary products!



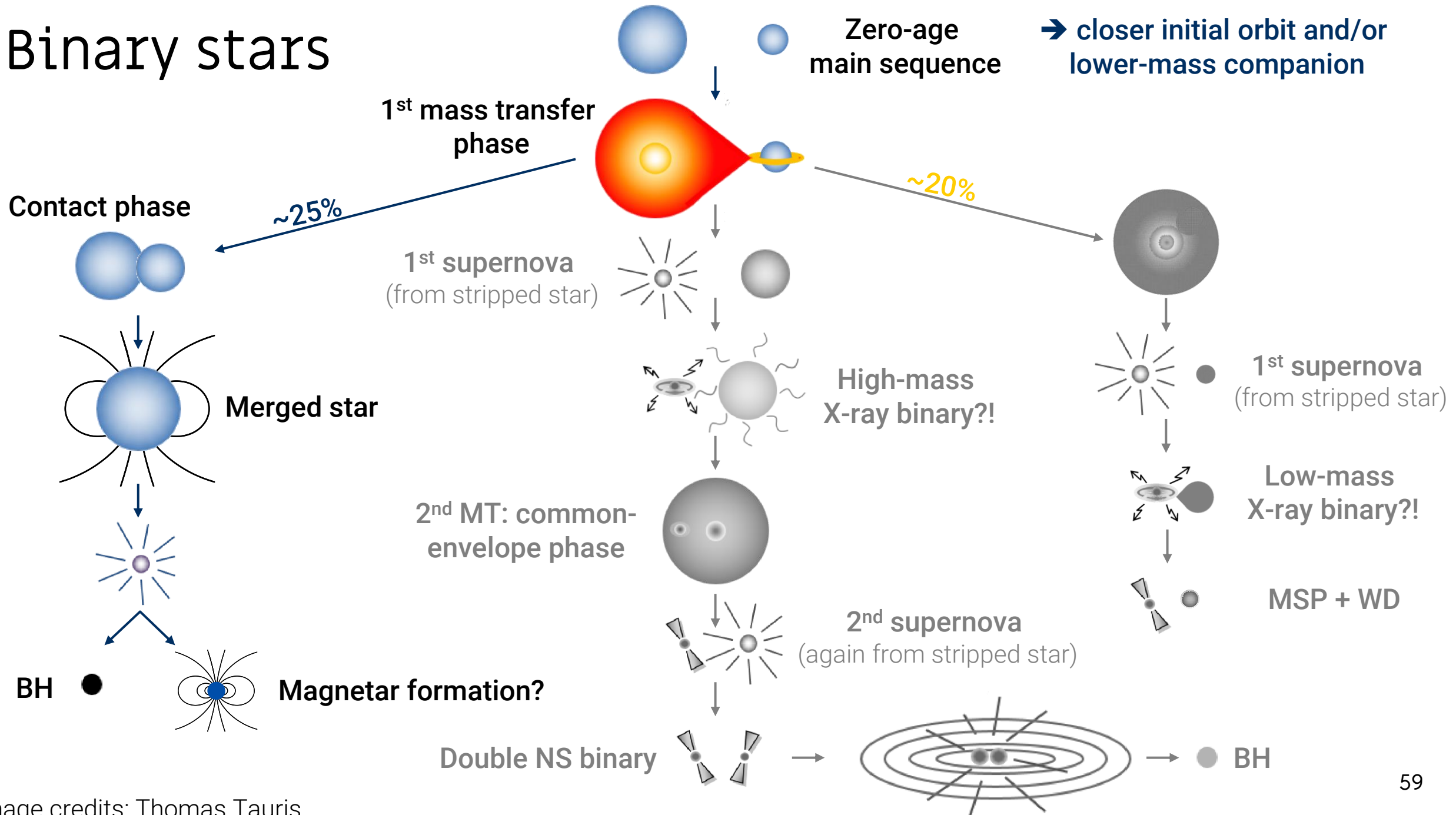
Binary stars



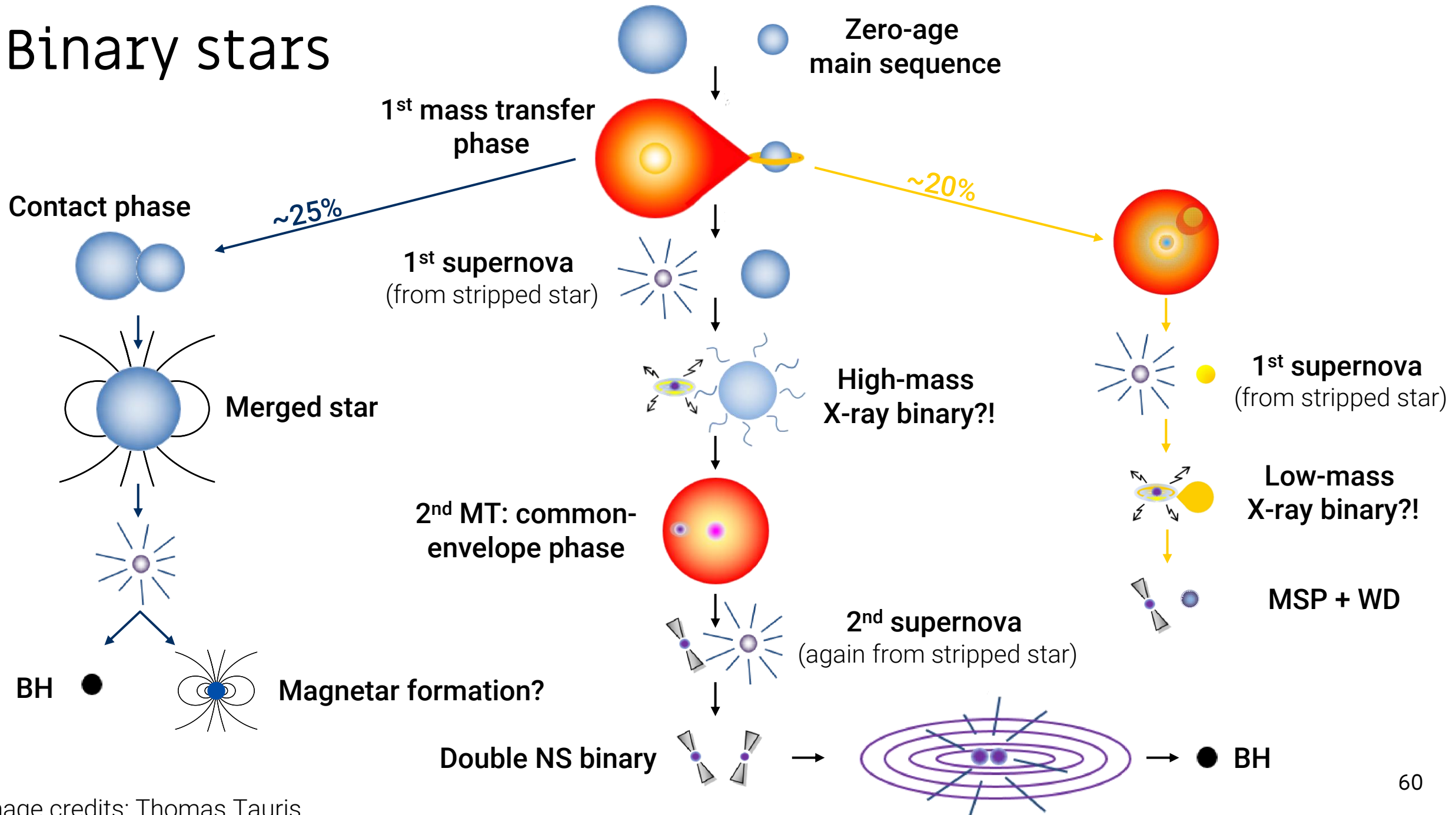
Binary stars



Binary stars



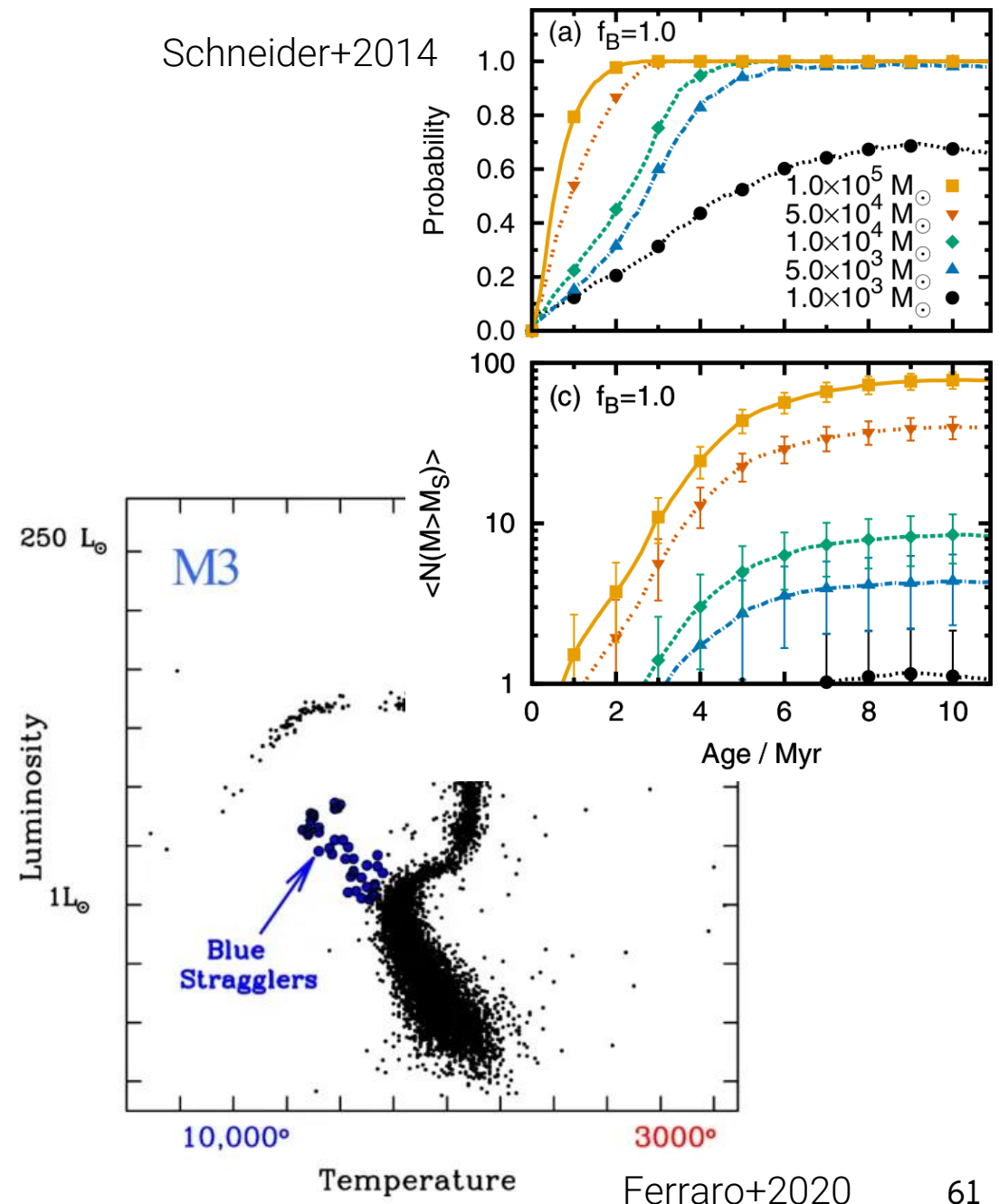
Binary stars



The most massive stars

Blue stragglers and rejuvenation

- **30 Doradus**: observationally, the **most massive ~30** out of ~1200 hot stars (\rightarrow WR, Of, WNh) contribute **40–50% to ionising and wind luminosity** (Doran+2013)
- Good chance that **several are blue stragglers**
- **Blue stragglers**: hot (i.e. blue) and luminous extension of main-sequence in star clusters
- **Origin**: accretors of binary mass transfer and/or stellar mergers (\rightarrow binary evolution & cluster dynamics)
- Blue stragglers are **rejuvenated** and **appear younger**
- Show up within Myr
- \rightarrow delayed feedback from massive stars

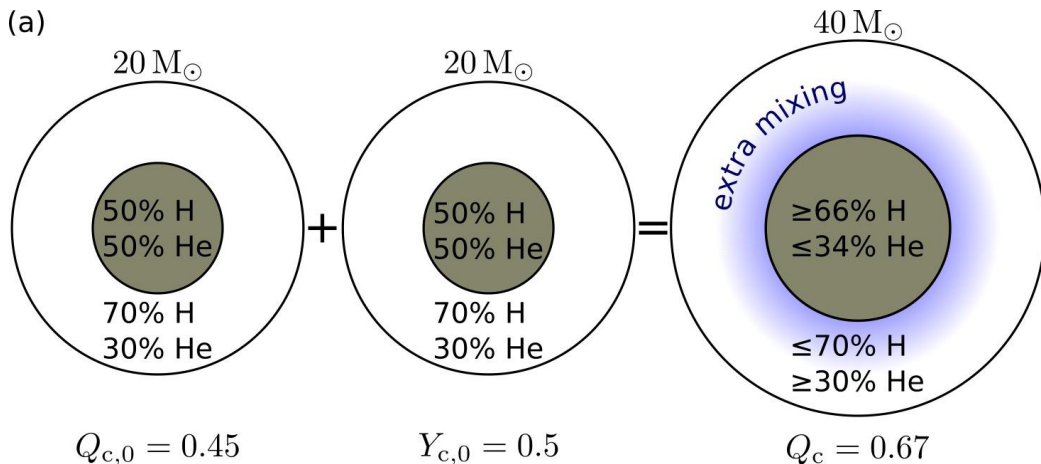


Blue straggler stars: rejuvenation

Rejuvenation

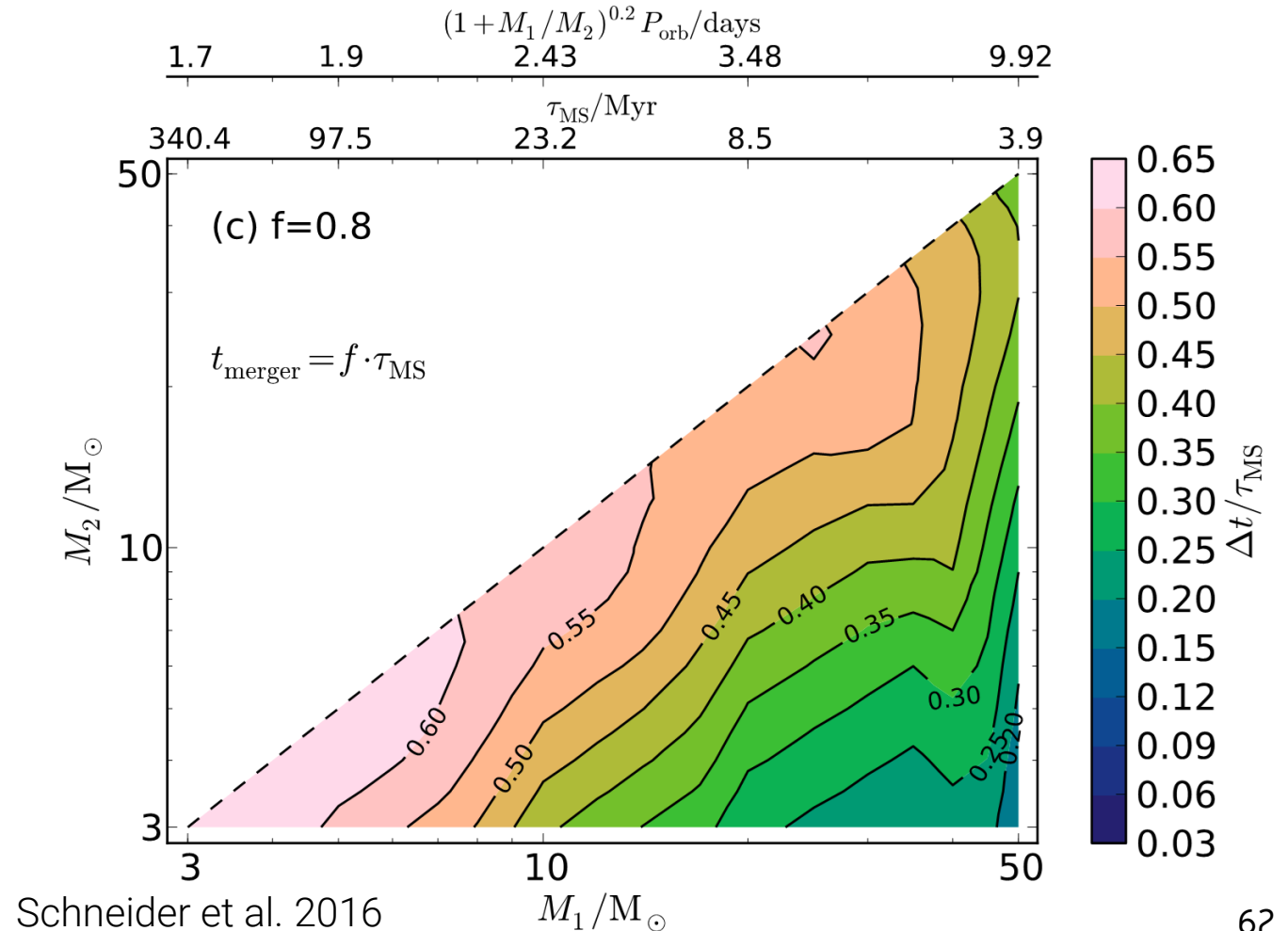
- **Apparent:** shorter lifetimes associated with more massive star
- **True:** mixing of fresh fuel into core (mixing during coalescence & transient conv. core during thermal relaxation)

Warning: cluster age determination



$$\Rightarrow Y_c = \frac{Q_{c,0}}{\alpha Q_c} Y_{c,0} \leq Y_{c,0} \quad \Rightarrow \Delta Y_c = Y_{c,0} - Y_c \propto Y_{c,0}$$

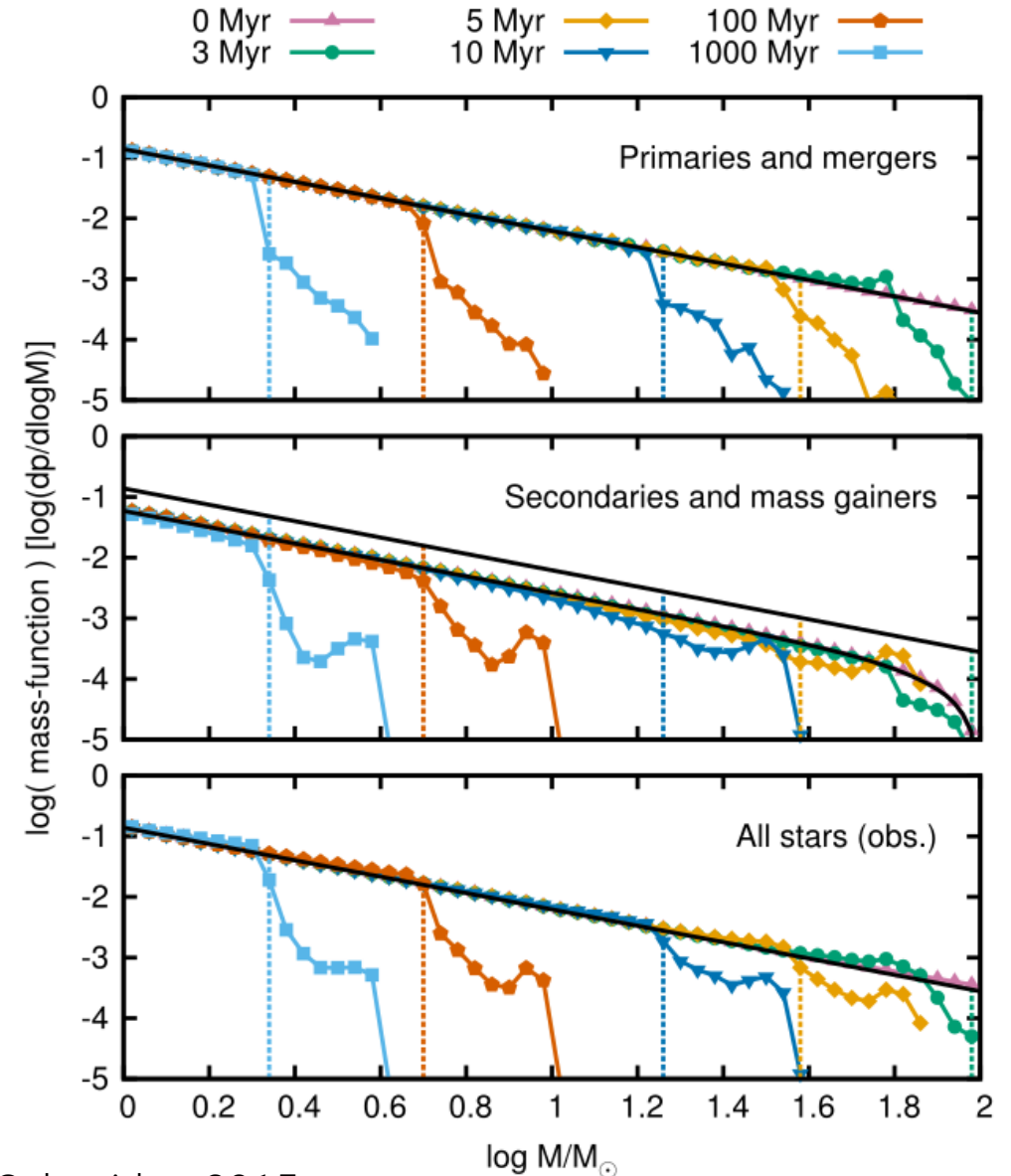
Simple rejuvenation recipes calibrated on mixing in SPH head-on collisions: **Glebbeek & Pols 2008**, Glebbeek et al. 2013, Schneider et al. 2016



The most massive stars

Effective mass function

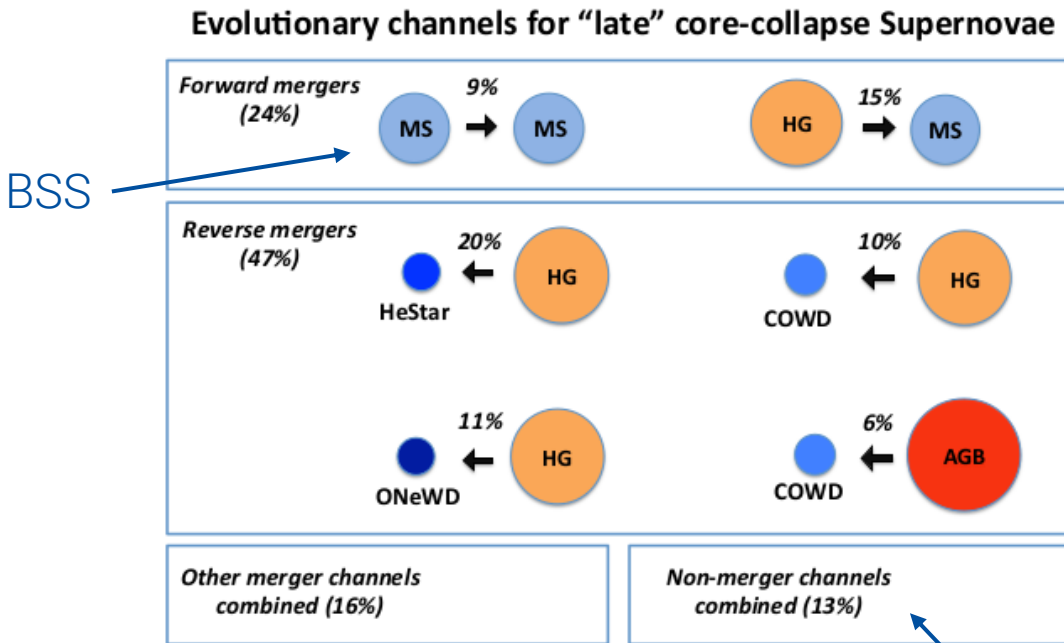
- Blue stragglers up to \sim twice more massive
- For feedback matters, treat as extension of the mass function of a stellar population
- **Example:** massive blue stragglers increase ionizing radiation and lead to delayed feedback because of rejuvenation
→ increase escape fraction of ionizing radiation (e.g. Secunda+2020)



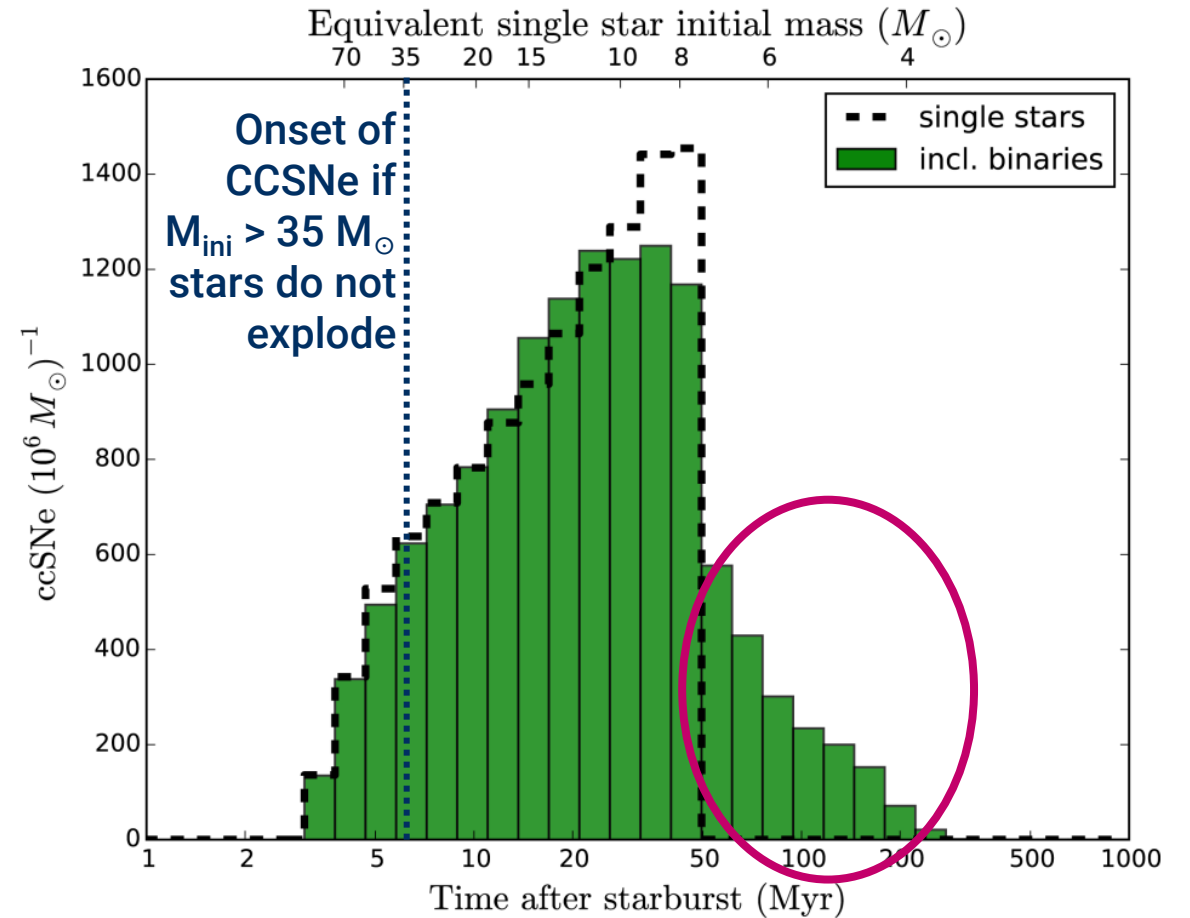
Schneider+2015

Delayed SN feedback from massive binary stars

- Binary mass accretion & mergers lead to late SNe compared to single stars
 - SNe from single stars over after ~ 50 Myr
 - With binaries, **$\sim 15\%$ of SNe at ages > 50 Myr**



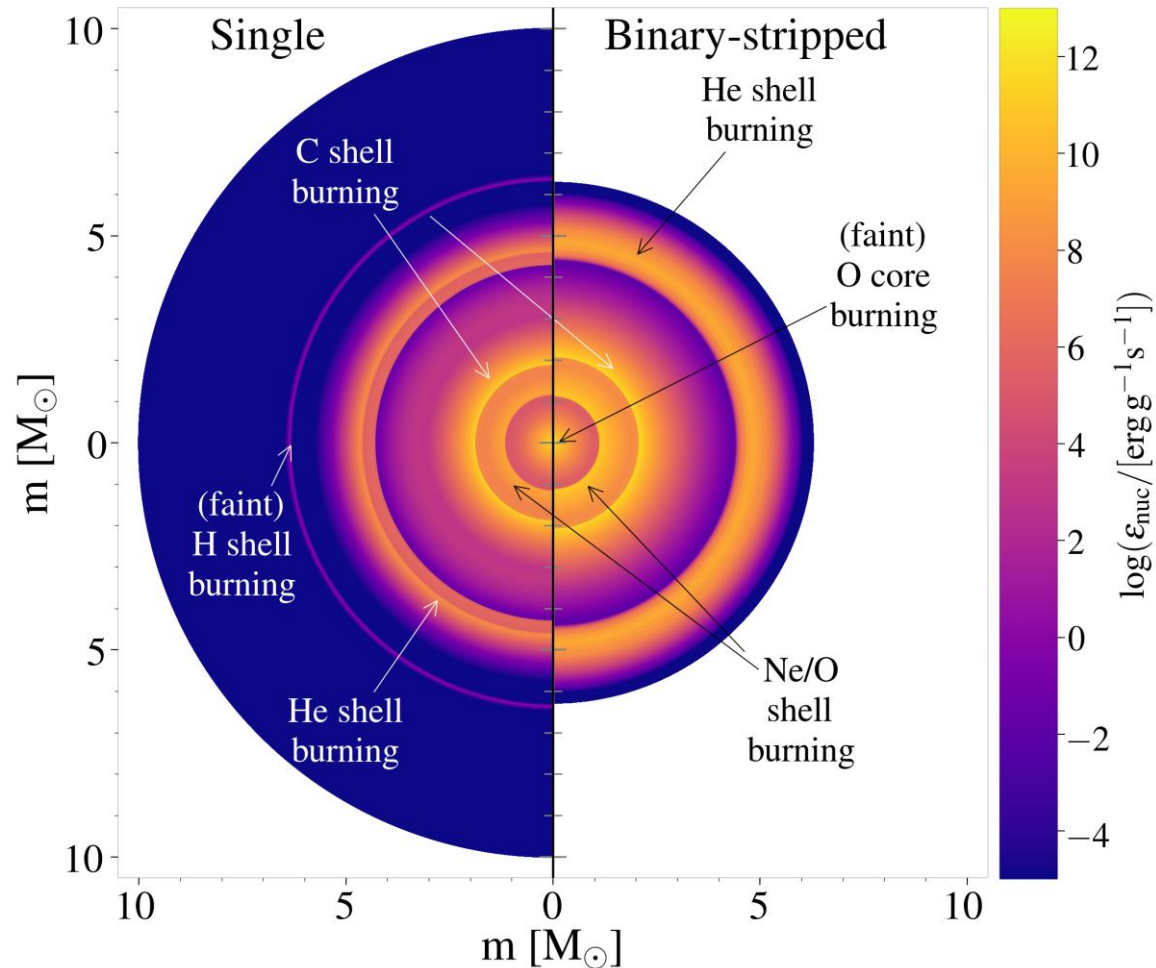
BSS (accreting secondaries)



Zapartas+2017

Stripped binary stars

Explodability and SN energetics



Nuclear burning structure in **single** vs. **binary-stripped star** at end of core oxygen burning
 → same $6.3 M_{\odot}$ core mass (Case B)

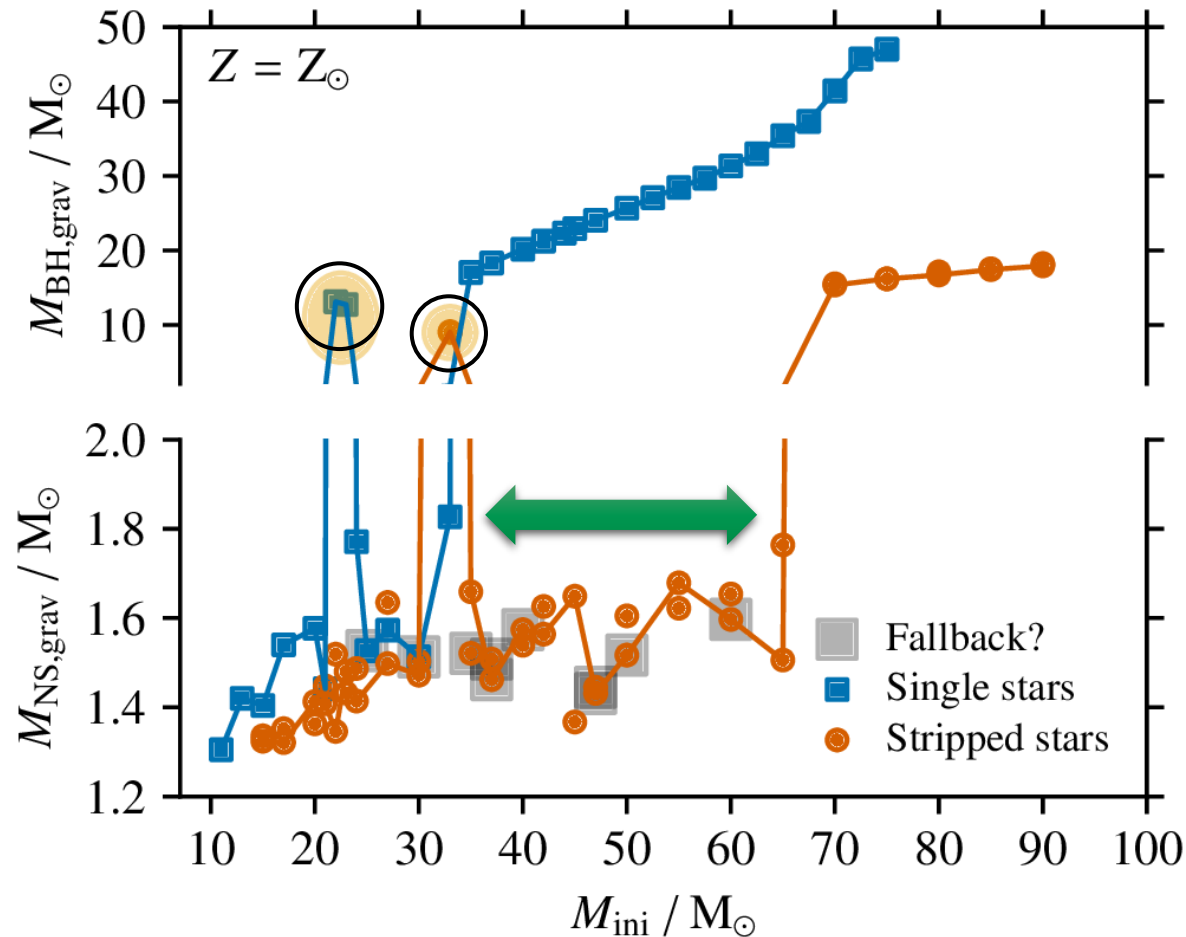
Clearly **different burning properties** mainly caused by absence of hydrogen shell burning

**Stripped stars have
 different pre-SN structure**

(cf. Timmes+1996, Wellstein & Langer 1999, Brown+2001, Podsiadlowski+2004, Woosley 2019, Schneider+2021, ...)

Stripped binary stars

Explodability and SN energetics

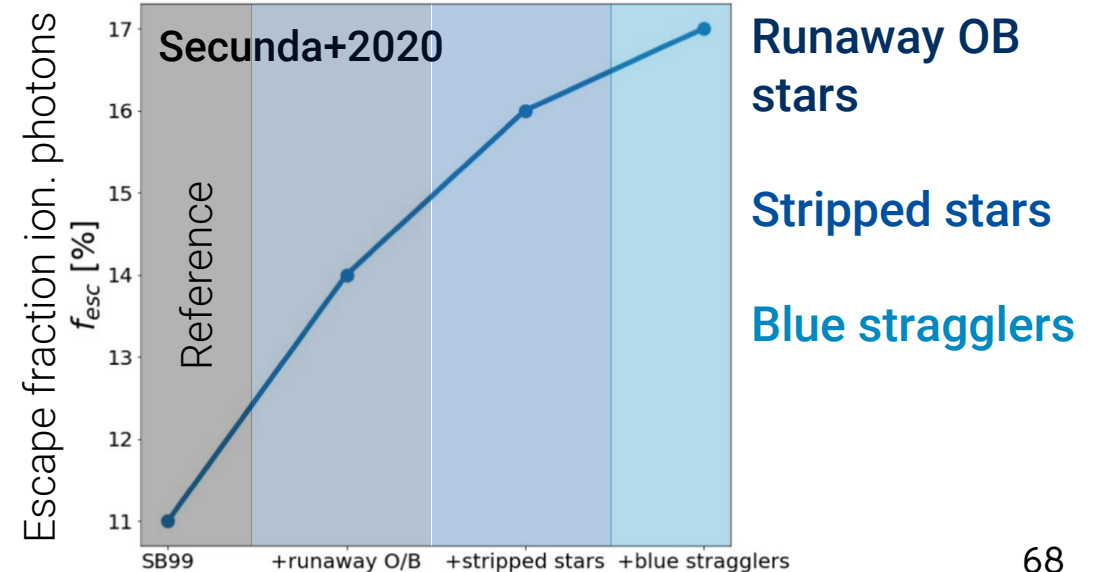
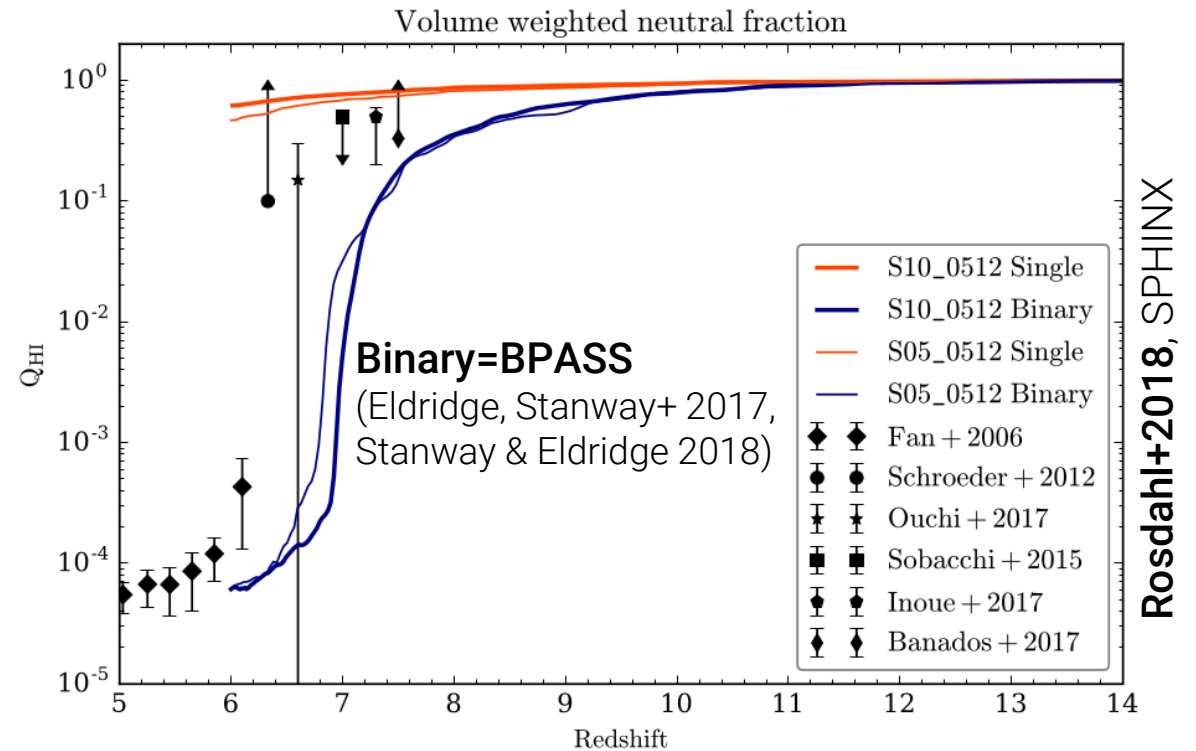
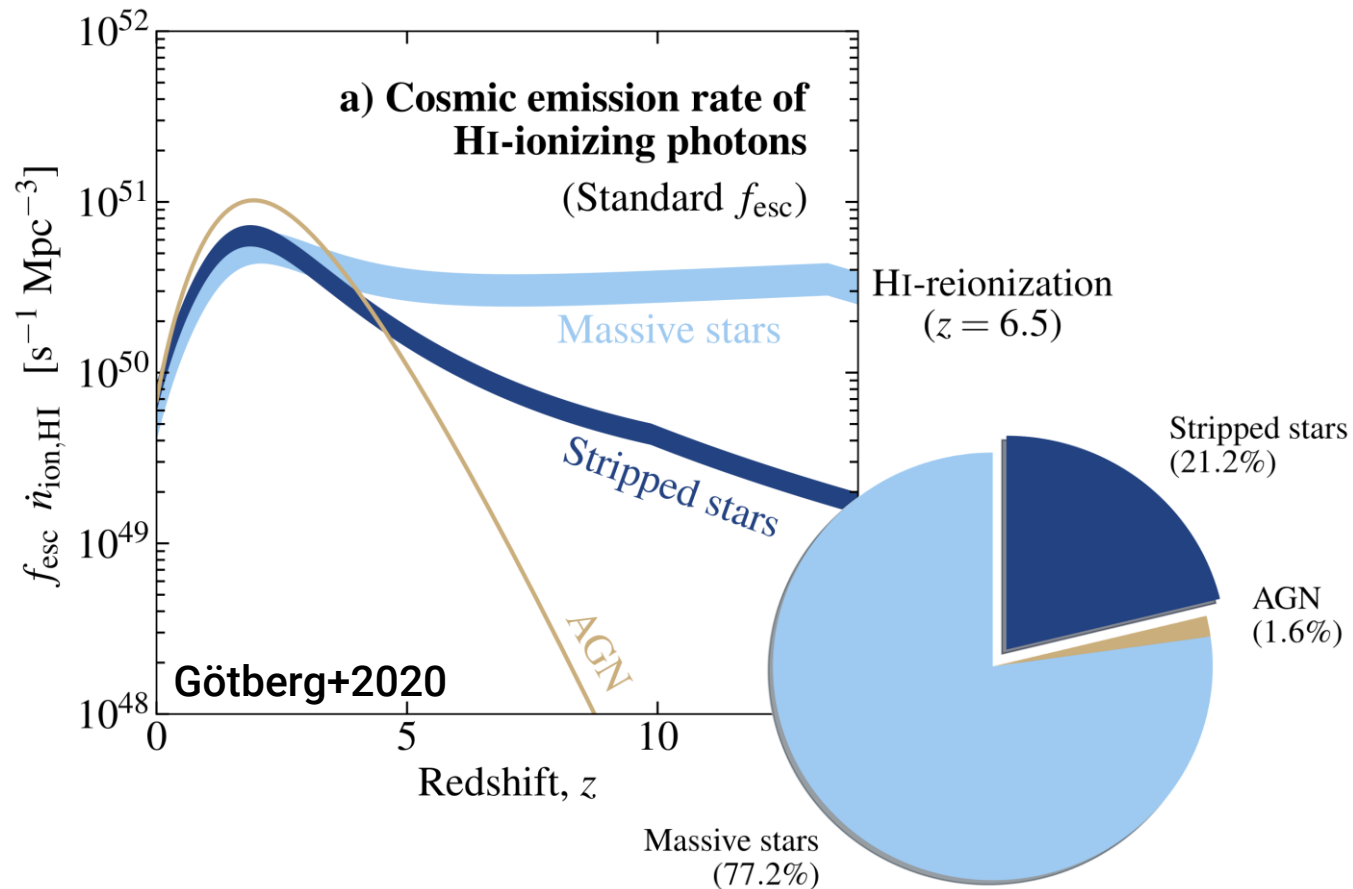


- **Island of BH formation** (compactness-peak BHs)
- **Neutron star formation/SN explosions:**
 - **Single stars:** $M_{\text{ini}} < 35 M_{\odot}$
 - **Binary-stripped stars:** $M_{\text{ini}} < 70 M_{\odot}$
- Stripped stars have **higher explosion energies** ($\Delta E_{\text{expl}} \sim 0.2 B$)
- **Implications for feedback**
 - First SNe from stripped stars (~ 4 Myr) & on average more energetic
 - SNe from genuine single stars after ~ 6 Myr
 - Metallicity dependence: **higher/lower $Z \rightarrow$ earlier/later SNe**
 - Number of SNe: 35 vs $40 M_{\odot} \rightarrow 3\%$ diff.; island of BH formation $\rightarrow 5\%$ diff.

Stripped binary stars

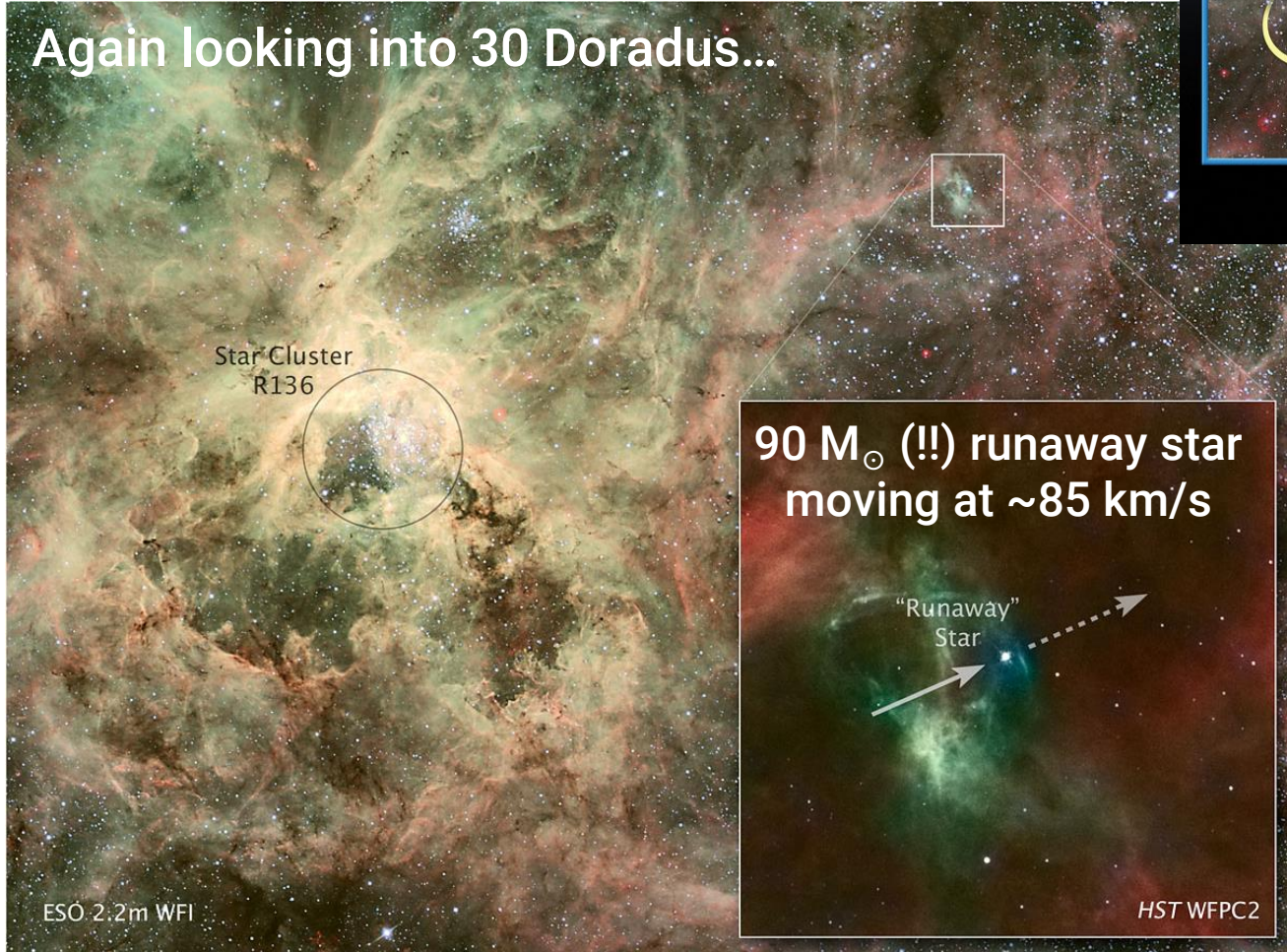
Cosmic reionisation

- **Massive stars:** hot, luminosities of up to $\sim 10^6$ solar
- **Stripped stars:** high $T_{\text{eff}} \rightarrow$ ionising radiation

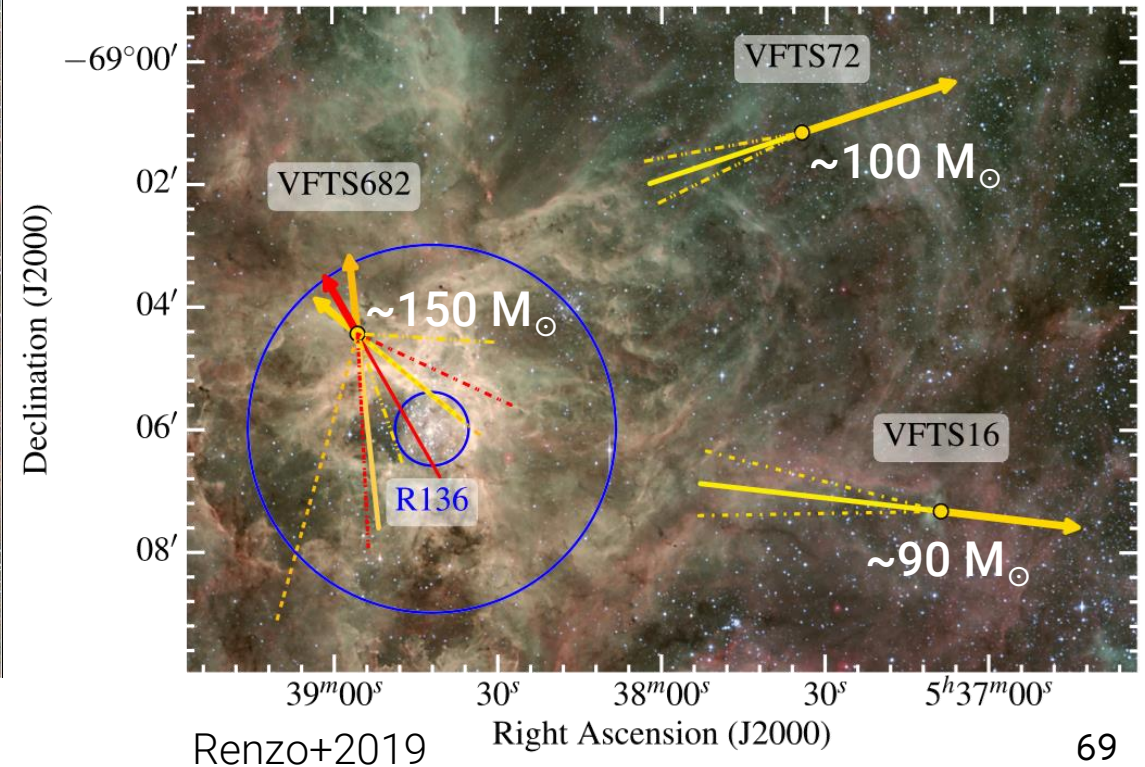
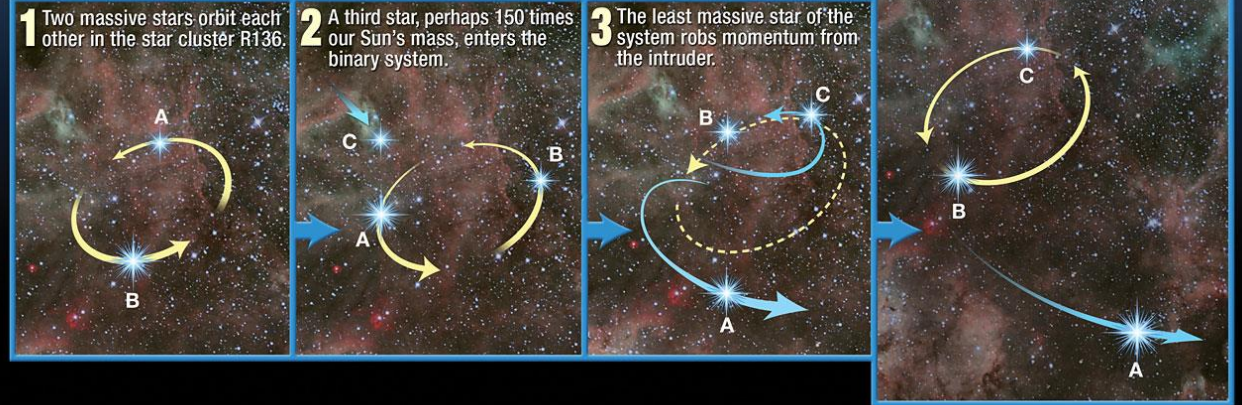


Runaway stars

Evans+2010



Massive star is ejected from a young star cluster



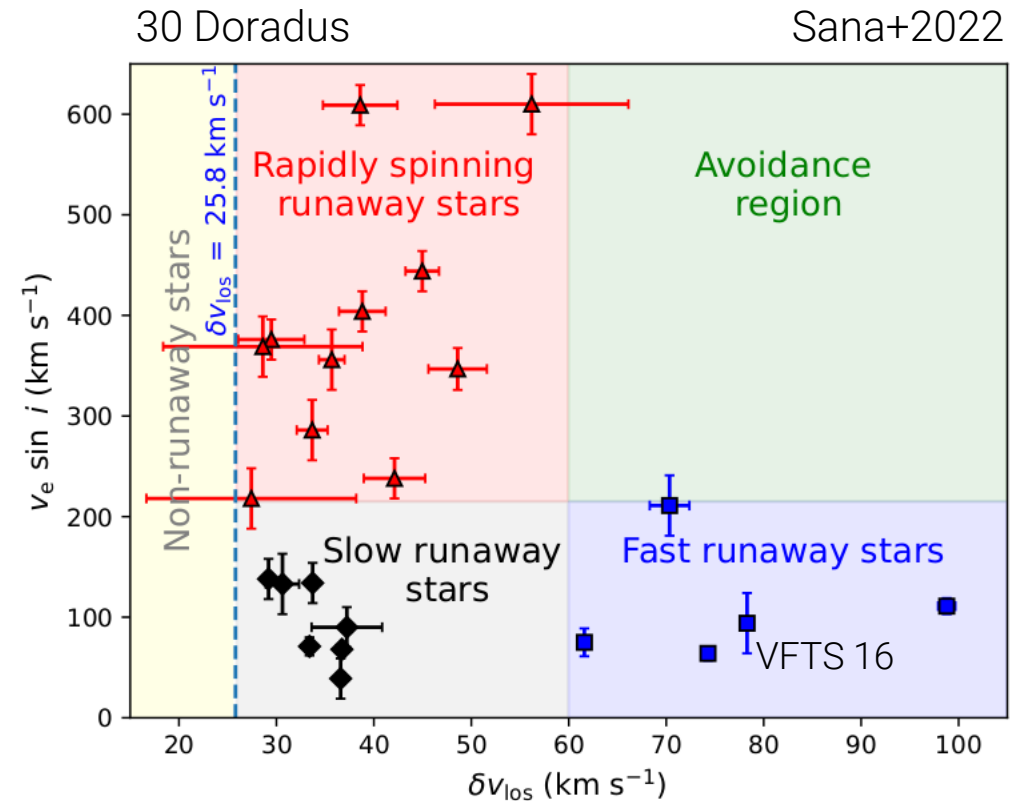
Credit: NASA/ESA/STScI/ESO

Runaway stars

- Up to 30% of OB stars move at $v > 30$ km/s
→ runaways
- Formation channels:
 - Dynamical ejection from dense stellar region
 - SN explosion in binary star system
 - Disruption if SN ejecta mass $>$ half of binary mass

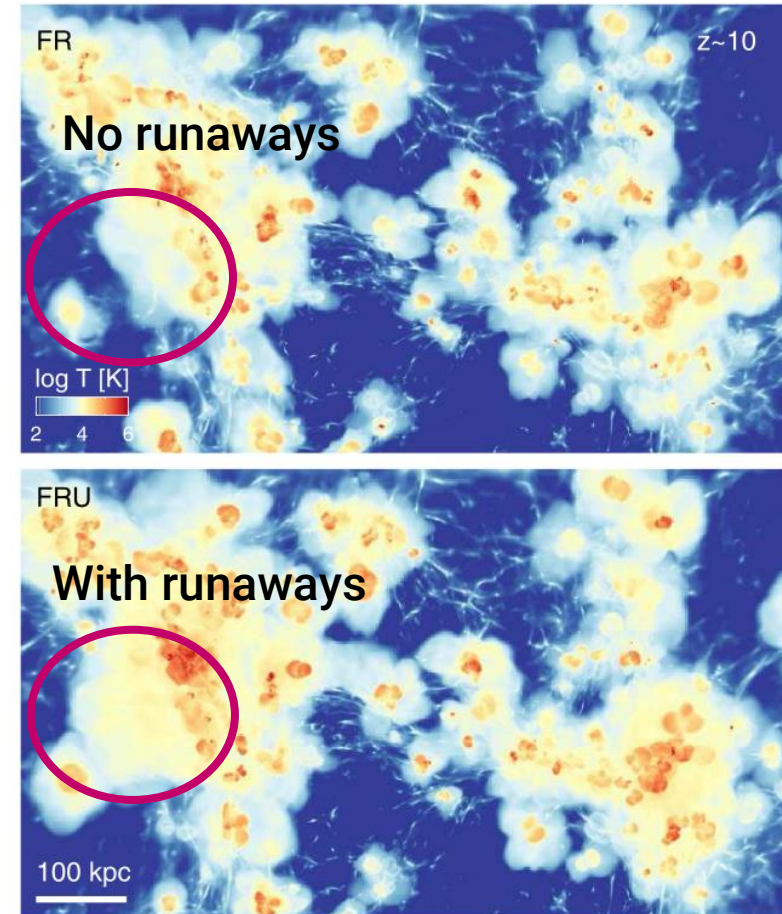
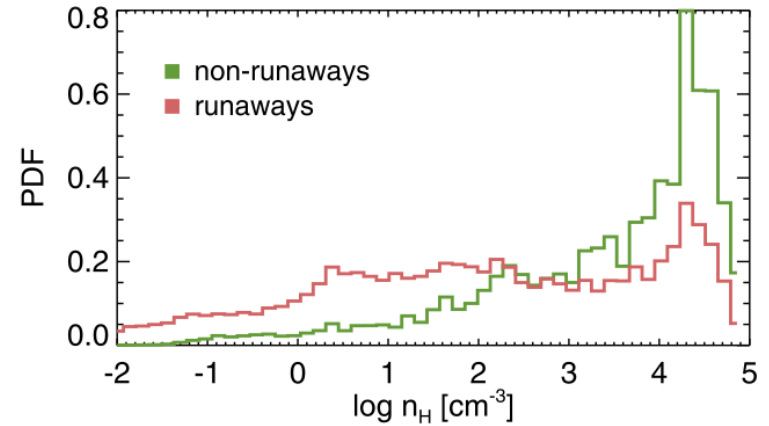
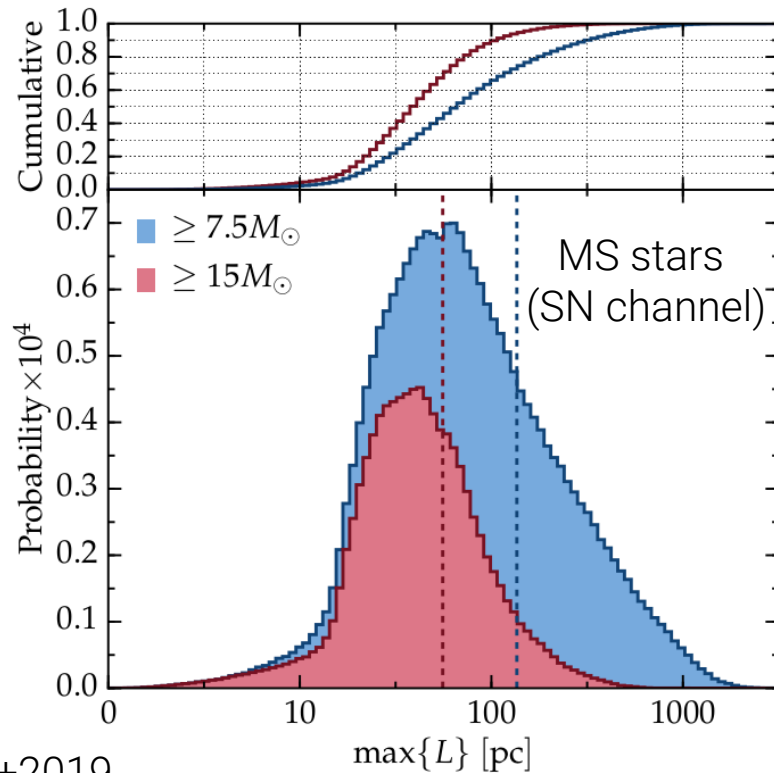
$$\Delta M > \frac{M_1 + M_2}{2}$$

- star ejected with orbital velocity
- Kick of NS/BH from SN explosion disrupts binary
→ star ejected with orbital velocity
→ **dominant contribution**
- Ejected MS stars: slow velocities → walkaways
(<30 - 40 km/s, Renzo+2019)



Runaway stars

- Runaways can travel large distances and deposit their feedback (radiation, SNe, winds, ...) far from SF sites
 → e.g. increases escape fraction of ionizing radiation by factors 1.1–8 (Conroy & Kratter 2012, Kimm & Cen 2014, Ma+2016, Secunda+2020)



X-ray binaries

Accreting NSs and BHs

Casares 2017

Three categories defined by mass of donor star

- **Low-mass X-ray binary (LMXB)**

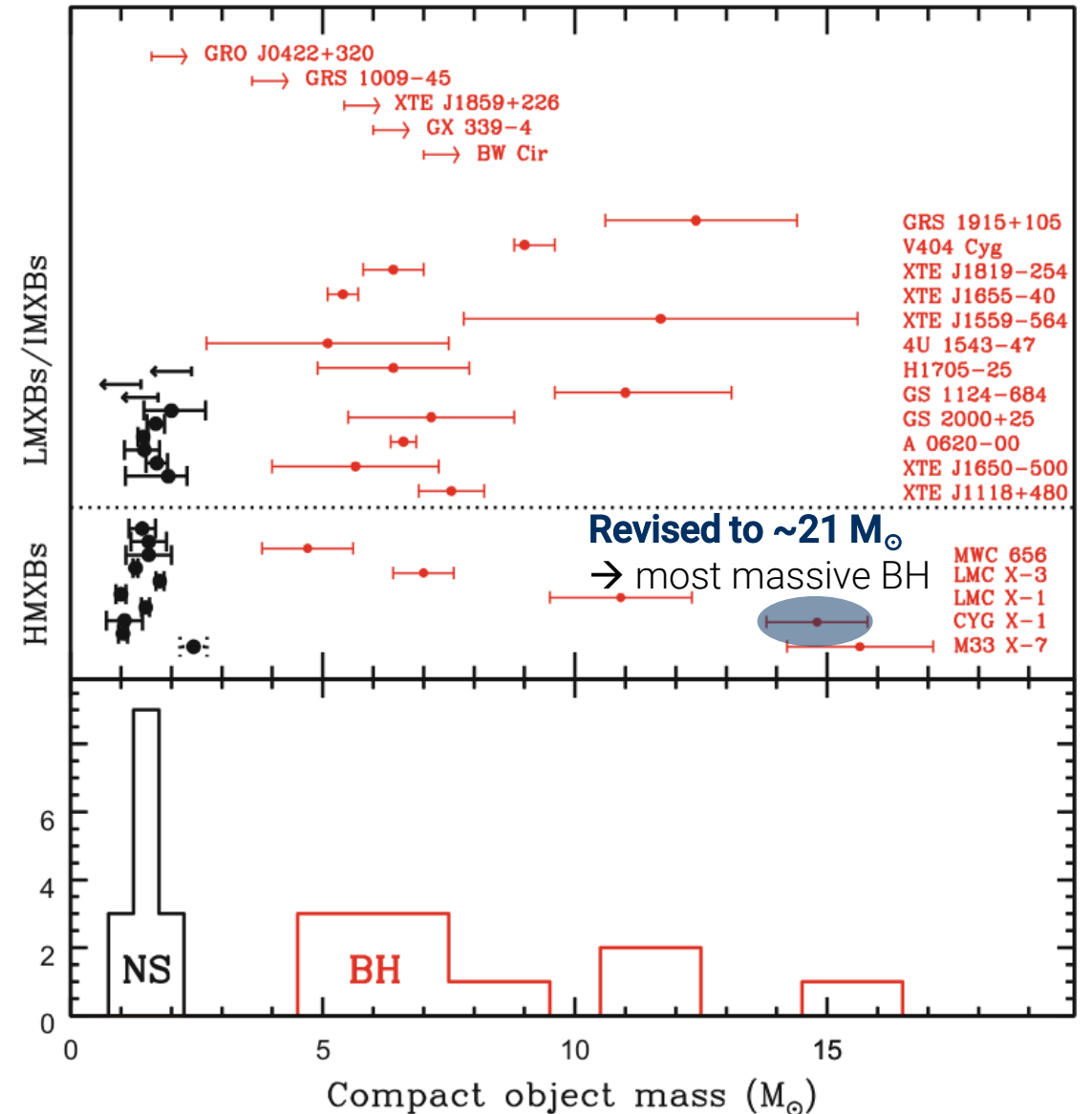
→ $M_{\text{donor}} < 1.5 M_{\odot}$ → **RLOF**, globular clusters, young & old stellar pops.

- **Intermediate-mass X-ray binary (IMXB)**

→ $1.5 M_{\odot} < M_{\text{donor}} < 5 M_{\odot}$

- **High-mass X-ray binary (HMXB)**

→ $M_{\text{donor}} > 10 M_{\odot}$ → **wind-fed**, young stellar pops.



X-ray binaries

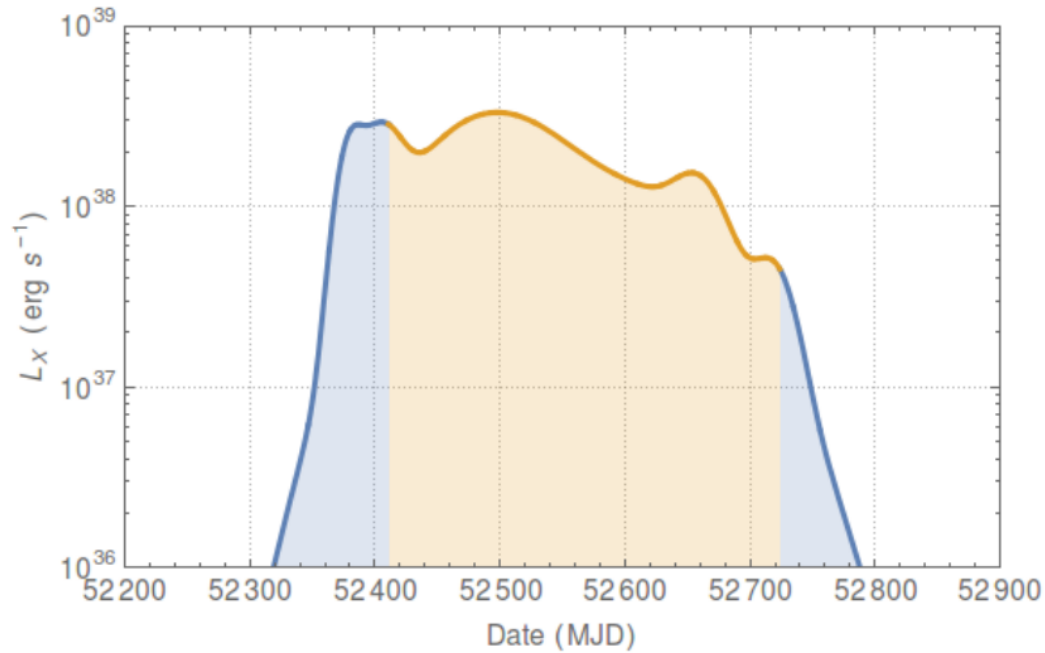
- **Persistent** and **transient sources**

- Transients → related to disk instability → AGNs

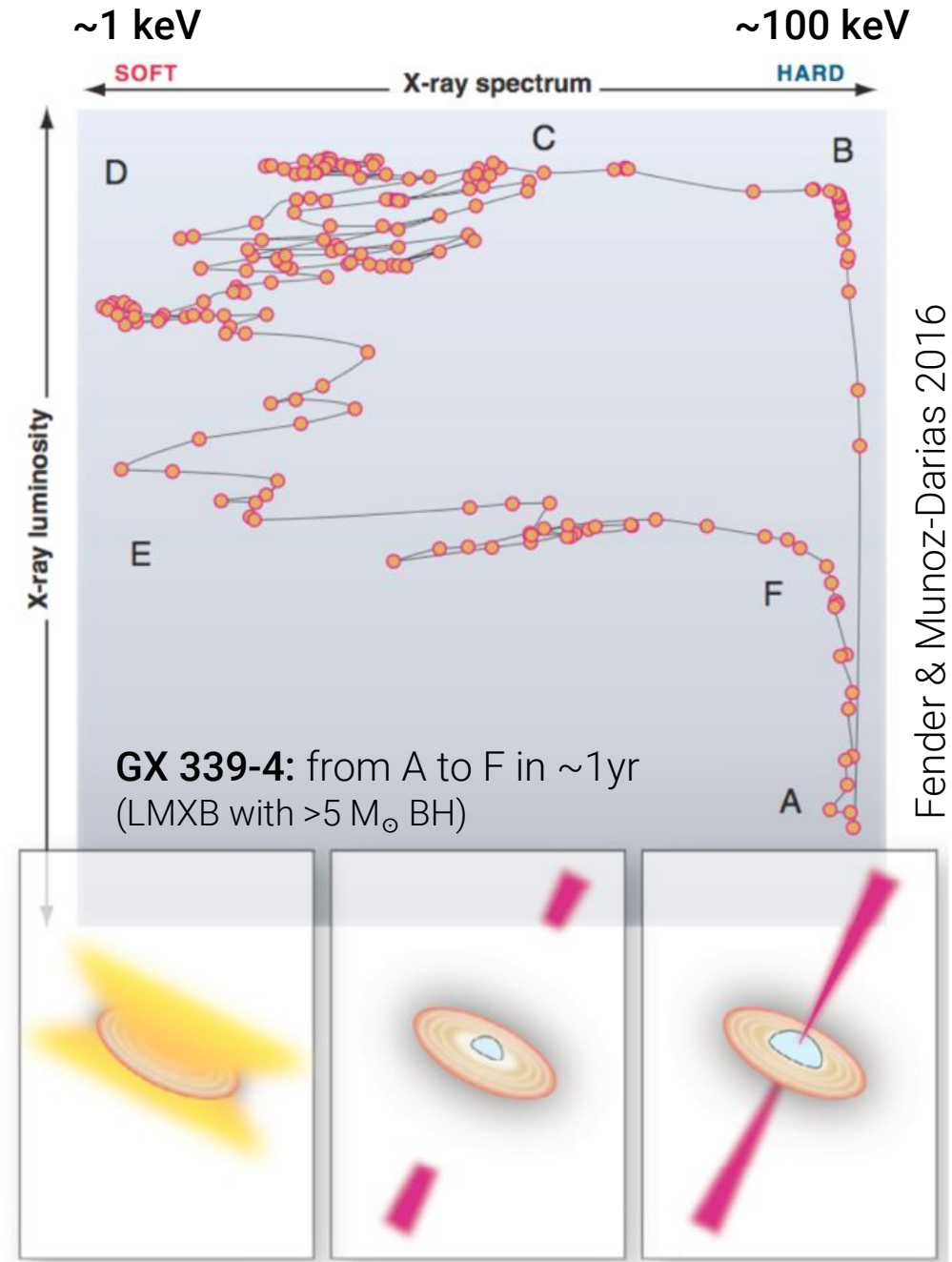
- Hard state: steady jet

- Soft state: disk wind, no jet

- Available accretion power: $P_{\text{acc}} = \frac{GM\dot{m}}{r} = \eta\dot{m}c^2$



Fender & Munoz-Darias 2016



Fender & Munoz-Darias 2016

X-ray binaries

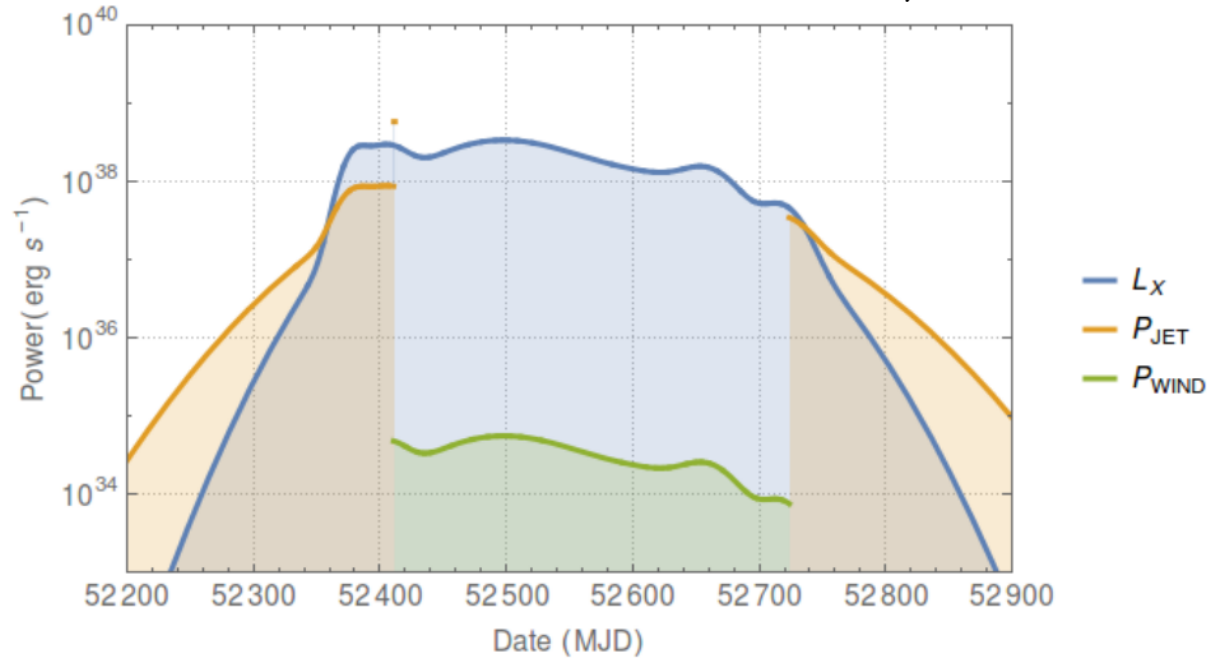
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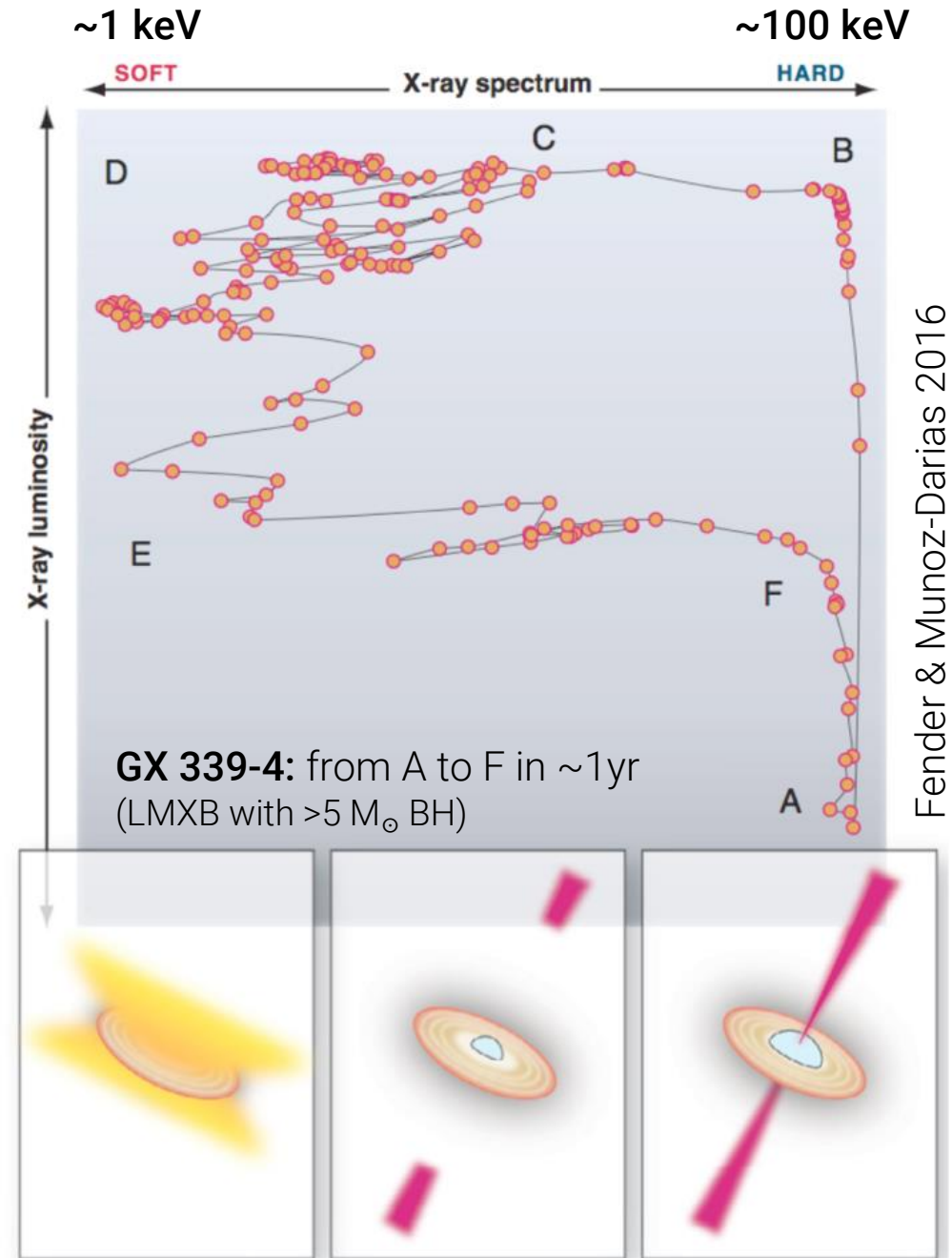
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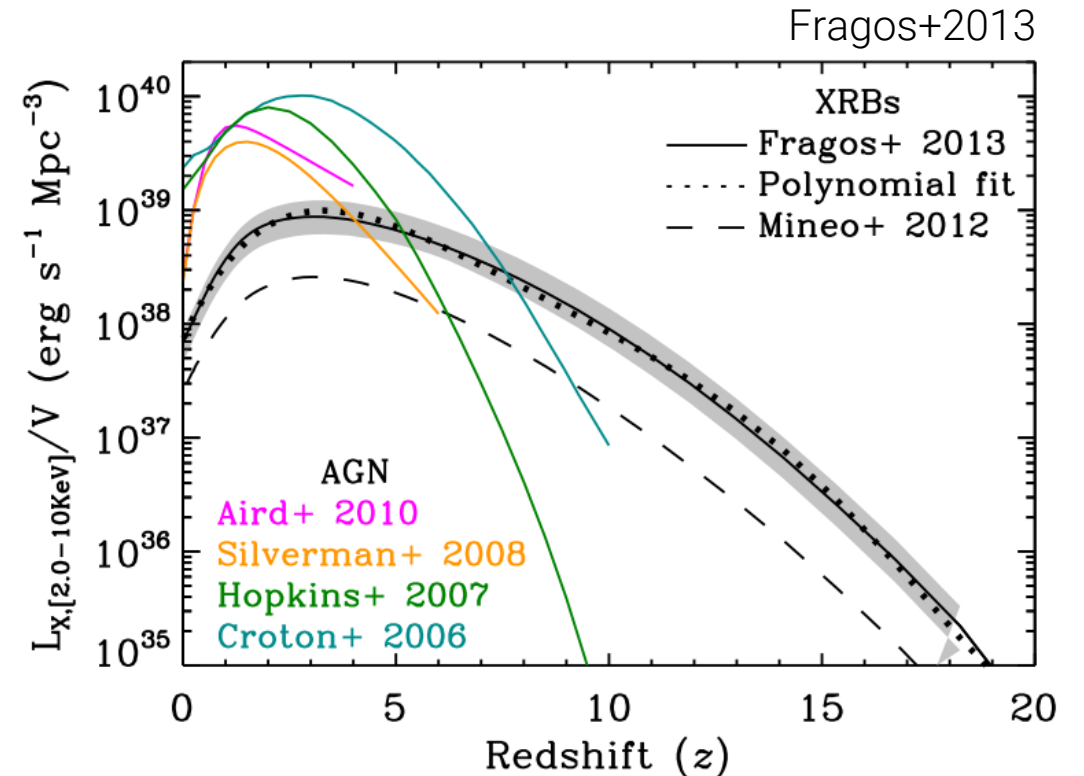


Fender & Munoz-Darias 2016



X-ray binaries

- X-ray sources in universe: AGNs and X-ray binaries (+ colliding wind binaries, massive stars, ...)
- High redshift: X-ray luminosity from X-ray binaries larger than that of AGNs?!
→ reionization!
- Models for X-ray binary feedback: Fragos+2013
- Further reading: Justham & Schawinski 2012

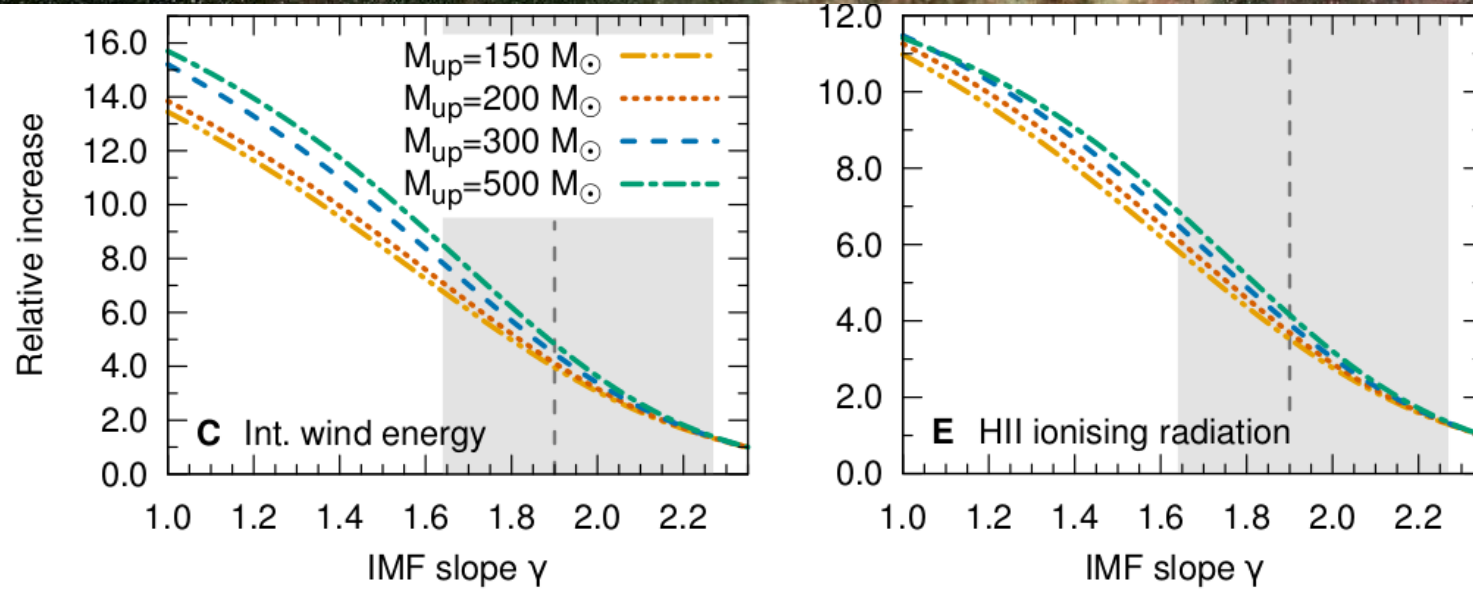
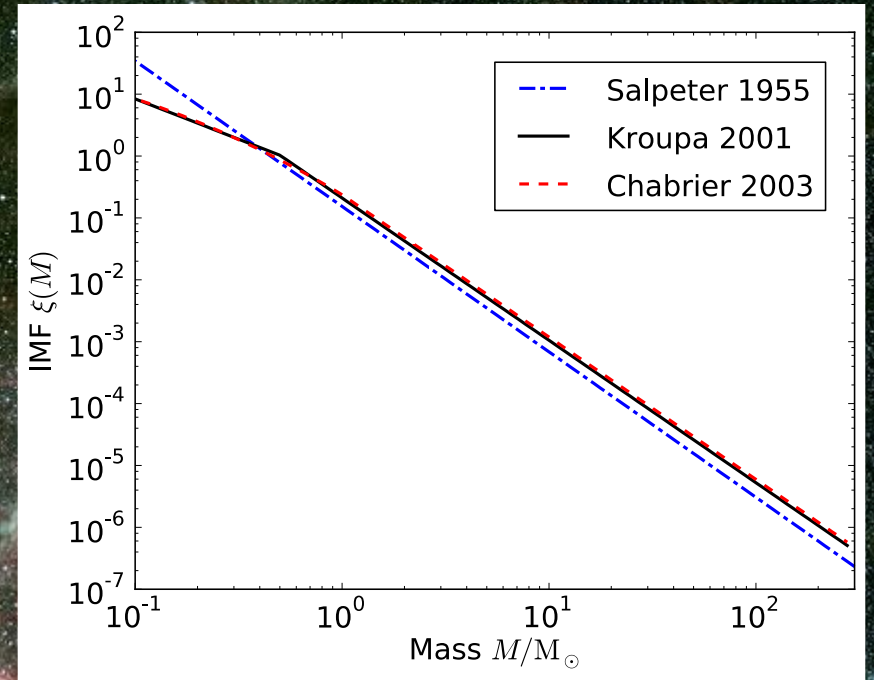


IMF and integrated stellar feedback

Stellar initial mass function (IMF)

- **0.5% of stars $>10 M_{\odot}$ contain 20% of mass** (1 CCSN/ $10^2 M_{\odot}$ and 1 PISN/ $10^4 M_{\odot}$)
- $\gamma < 2.00$: most mass in high mass stars & need M_{\max}

30 Doradus: top-heavy IMF ($\gamma \sim 1.90$; Schneider+2018)
 → enhanced feedback!



$$\xi(M) = \frac{dN}{dM} \propto M^{-\gamma}$$

$$(\Gamma = \gamma - 1)$$

Salpeter:
 $\gamma=2.35$
 $(\Gamma=1.35)$

Summary stellar feedback

Stellar feedback = injection of **matter**, **momentum** and **energy** by stars into the circumstellar medium
→ **effective feedback depends on environment**

Feedback processes

- Photons & ionising radiation
- Winds
- Accretion, jets & outbursts (proto-/pre-MS stars)
- Supernovae
- Cosmic rays
- Binary stars
- X-ray binaries
- Runaway stars
- ...

Stellar feedback...

- ...enriches Universe with chemical elements
- ...helped to reionise Universe after dark ages
- ...dictates appearance & evolution of galaxies
- ...regulates star formation (inhibits & triggers SF)
→ (in-)efficiency of star formation
- ...drives galactic winds and affects intergalactic medium (chemical enrichment, temperature, ...)
- ...sets life-cycle of giant molecular clouds
- ...is necessary for forming realistic galaxies in simulations (e.g. properties at $z=0$)
→ also limits formation of dwarf galaxies