

# Dust dynamics and **dust evolution** during star formation

~~(no) more equations~~

**Less equations**

# Short summary of yesterday

- Dust dynamics is important and can be very different from the gas during star formation and disk evolution.
- The dynamics of grains is controlled by the Stokes number which depend on the grain and gas properties.
- Dust grains settle fast in protoplanetary disks and drift radially.
- The radial drift is fast compared to the disk lifetime. To understand for planet formation it is key to find mechanisms to overcome the radial drift barrier

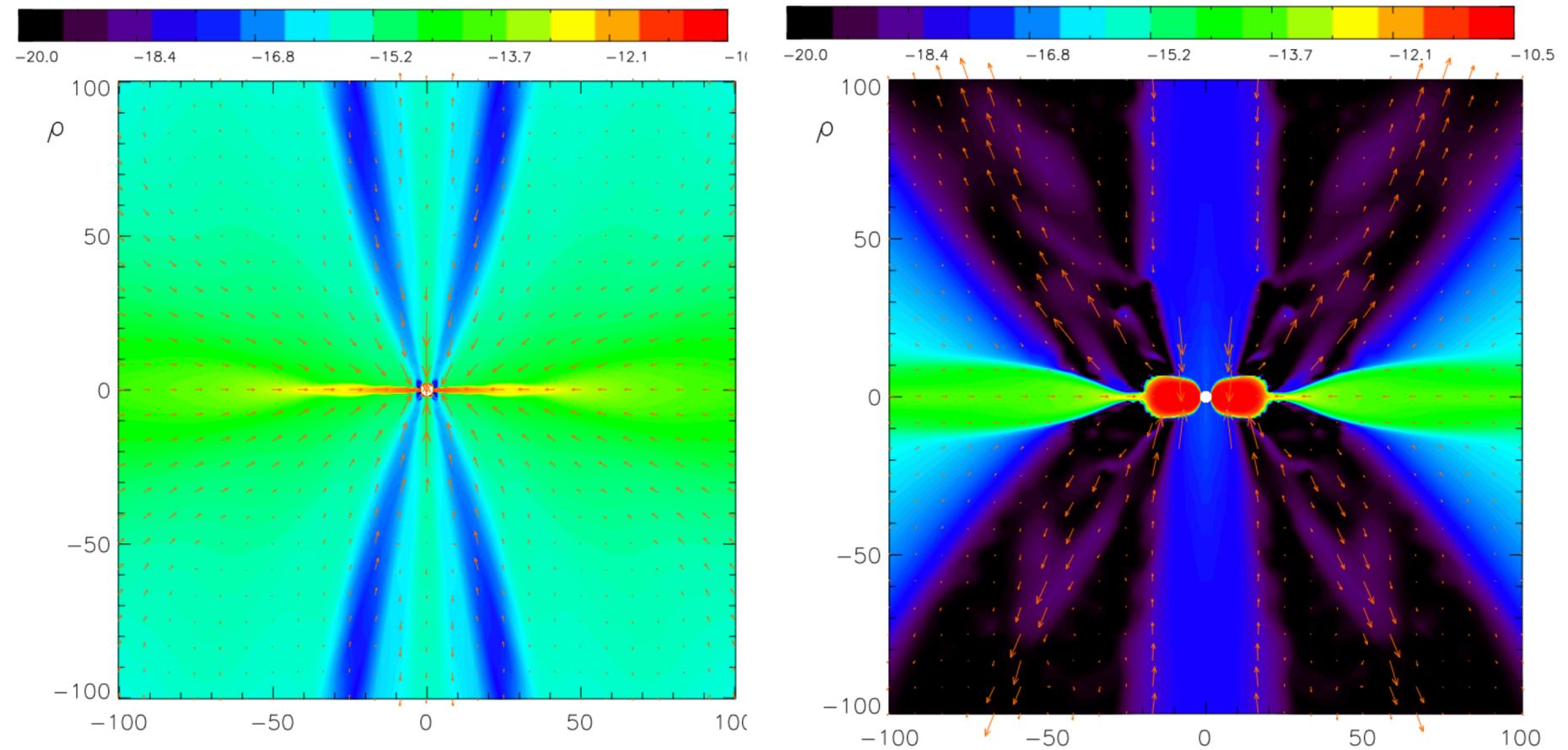
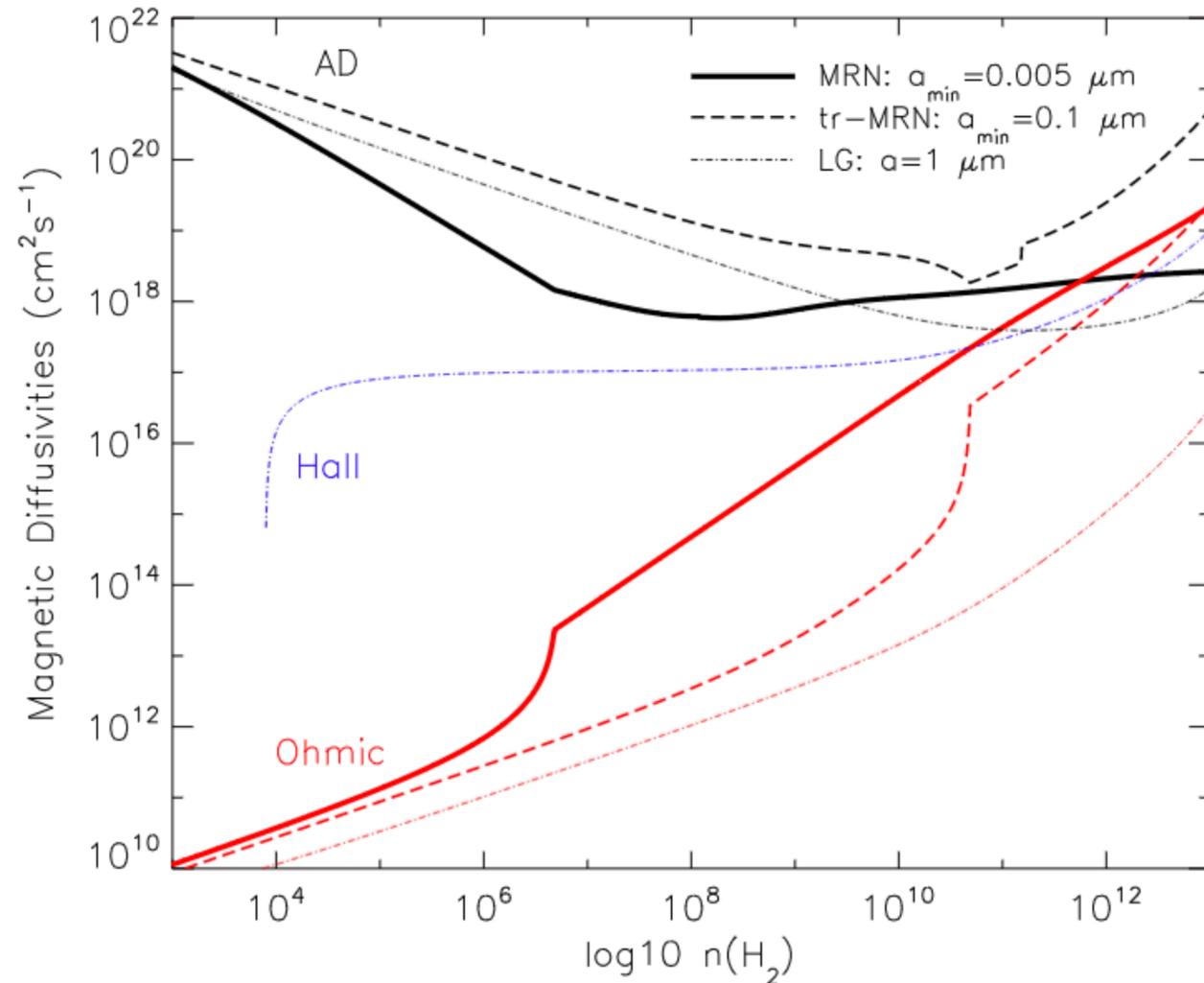
# Short summary of yesterday

$$\frac{\partial \vec{B}}{\partial t} + \underbrace{\nabla \times (\vec{v} \times \vec{B})}_{\text{Advection}} = \underbrace{\nabla \times [\eta_o \nabla \times \vec{B}]}_{\text{Ohmic dissipation}} - \underbrace{\nabla \times \left[ \frac{\eta_A}{|\vec{B}|^2} \left( (\nabla \times \vec{B}) \times \vec{B} \right) \times \vec{B} \right]}_{\substack{\text{Ambipolar "diffusion"} \\ \text{Or ion-neutral drift}}} + \underbrace{\nabla \times \left[ \frac{\eta_H}{|\vec{B}|^2} (\nabla \times \vec{B}) \times \vec{B} \right]}_{\text{Hall effect}}$$

The resistivities depends on the **abundance** of charged particles, their **charge**, their **Hall factor**, the magnetic field strength

# Role of grains in the MHD equations

Changing the dust size distribution (by for example removing the very small grains) can have a dramatic influence on the resistivities and therefore on the disk formation



*Zhao, Caselli et al. (2016) : See also Marchand et al. (2020)*

$$r_{d,AD} \simeq 18 \text{ au}$$

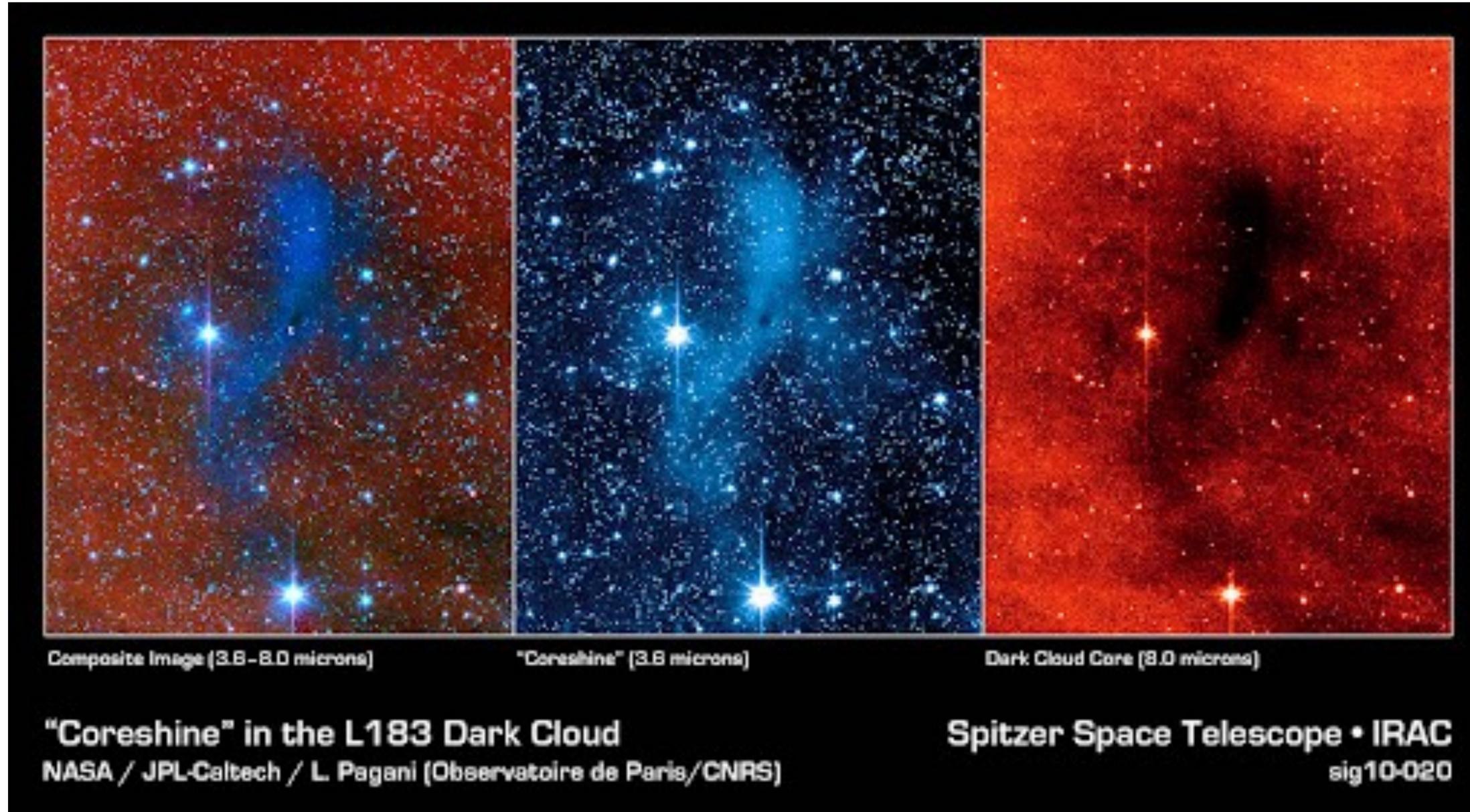
$$\times \delta^{2/9} \left( \frac{\eta_{AD}}{0.1 \text{ s}} \right)^{2/9} \left( \frac{B_z}{0.1 \text{ G}} \right)^{-4/9} \left( \frac{M_d + M_*}{0.1 M_\odot} \right)^{1/3} \text{ Hennebelle et al. (2016)}$$

# Dust evolution

# Observational evidence of dust evolution

## In prestellar cores :

IR excess (coreshine) could be explained by micron sized grains. (Pagani 2011)

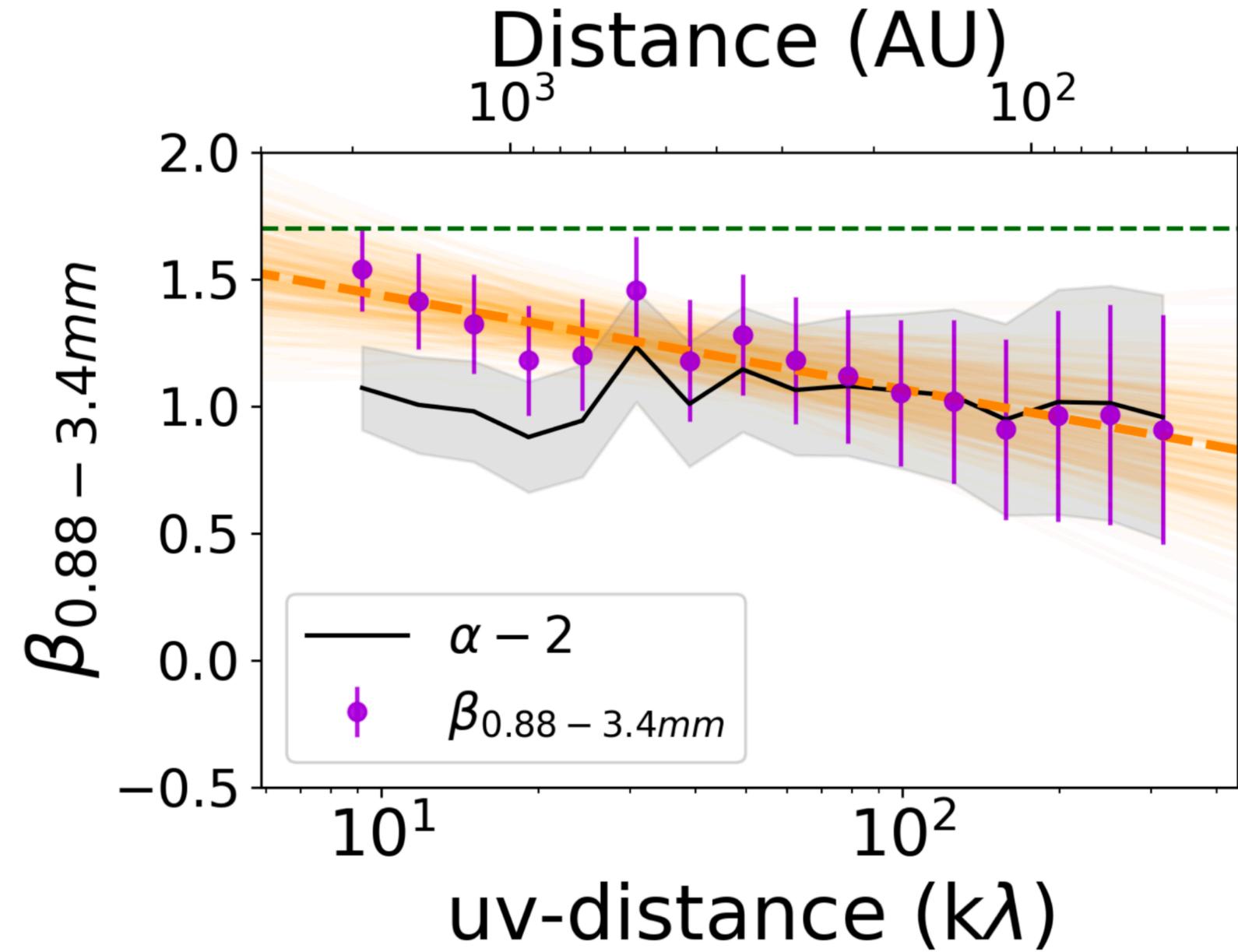


# Observational evidence of dust evolution

## In protostars envelopes :

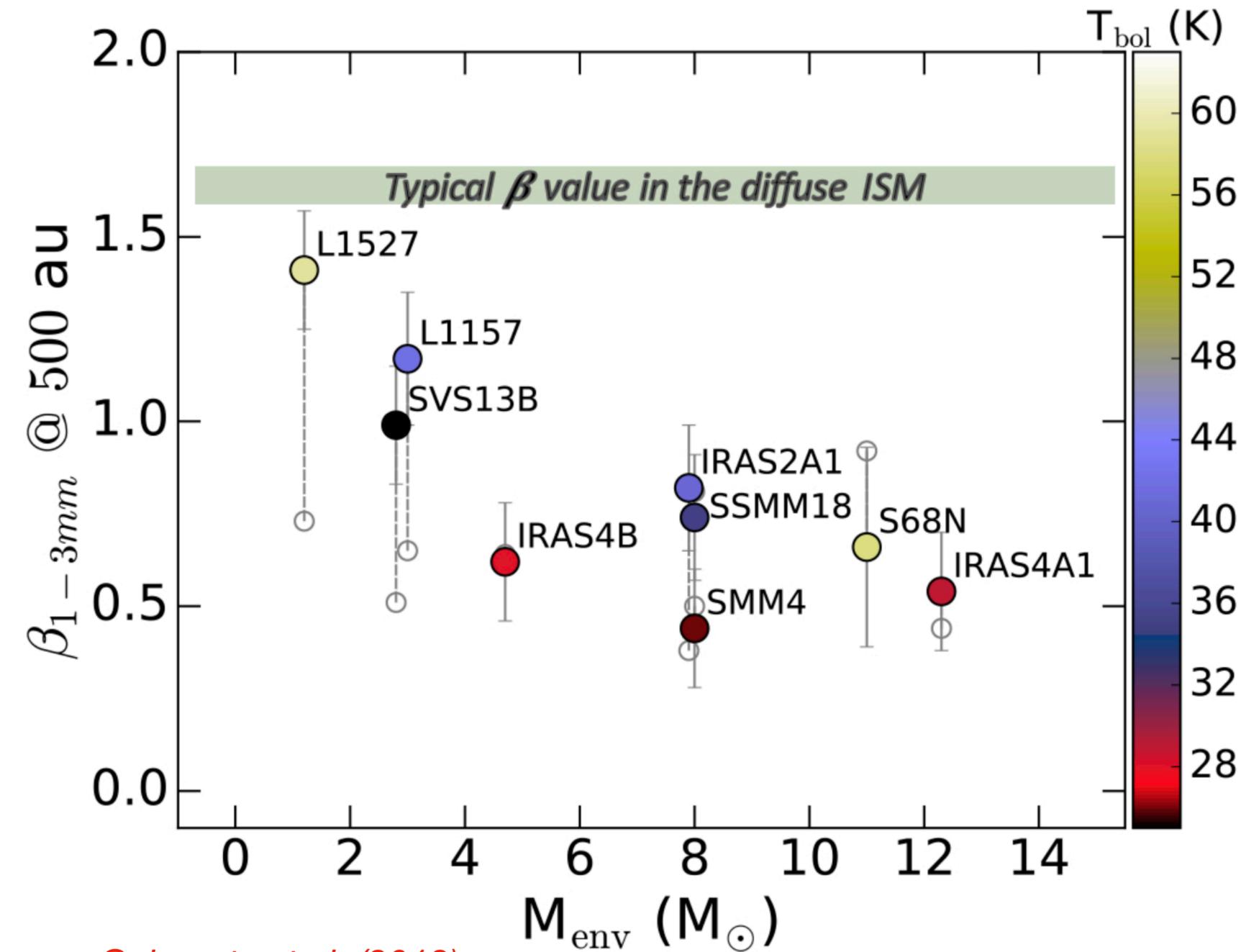
- Slope of the SED might be explained by grown  $\sim 100$  micron grains. (e.g. [Galamez et al. 2019](#); [Cacciapuoti et al. 2023,2024a](#))
- Same goes with polarisation fraction ([Valdivia et al. 2019](#))

**Note :** this is tentative because large grain optical properties are not very well constrained (yet).

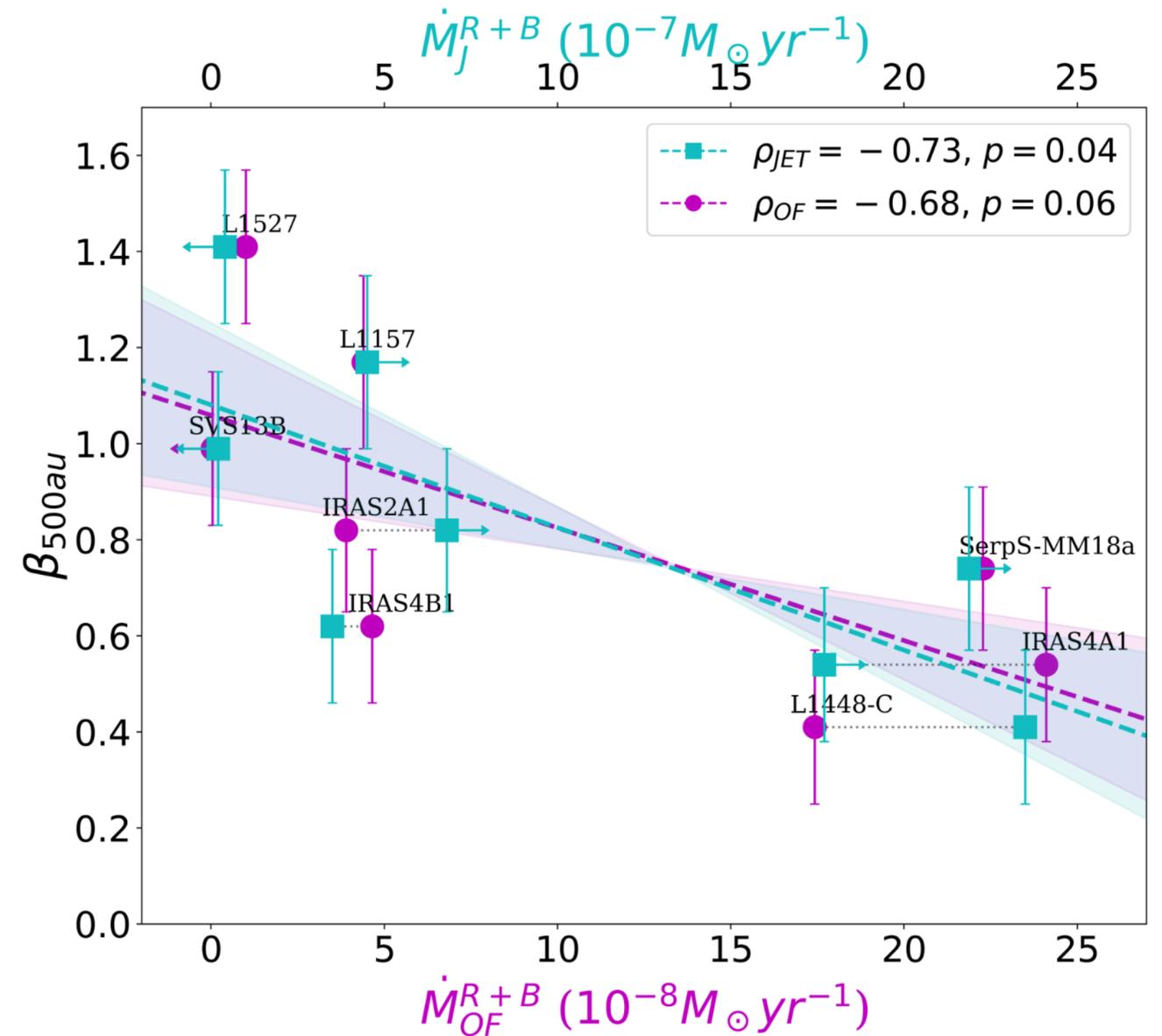


[Cacciapuoti et al. 2023](#)

# Observational evidence of dust evolution



*Galametz et al. (2019)*



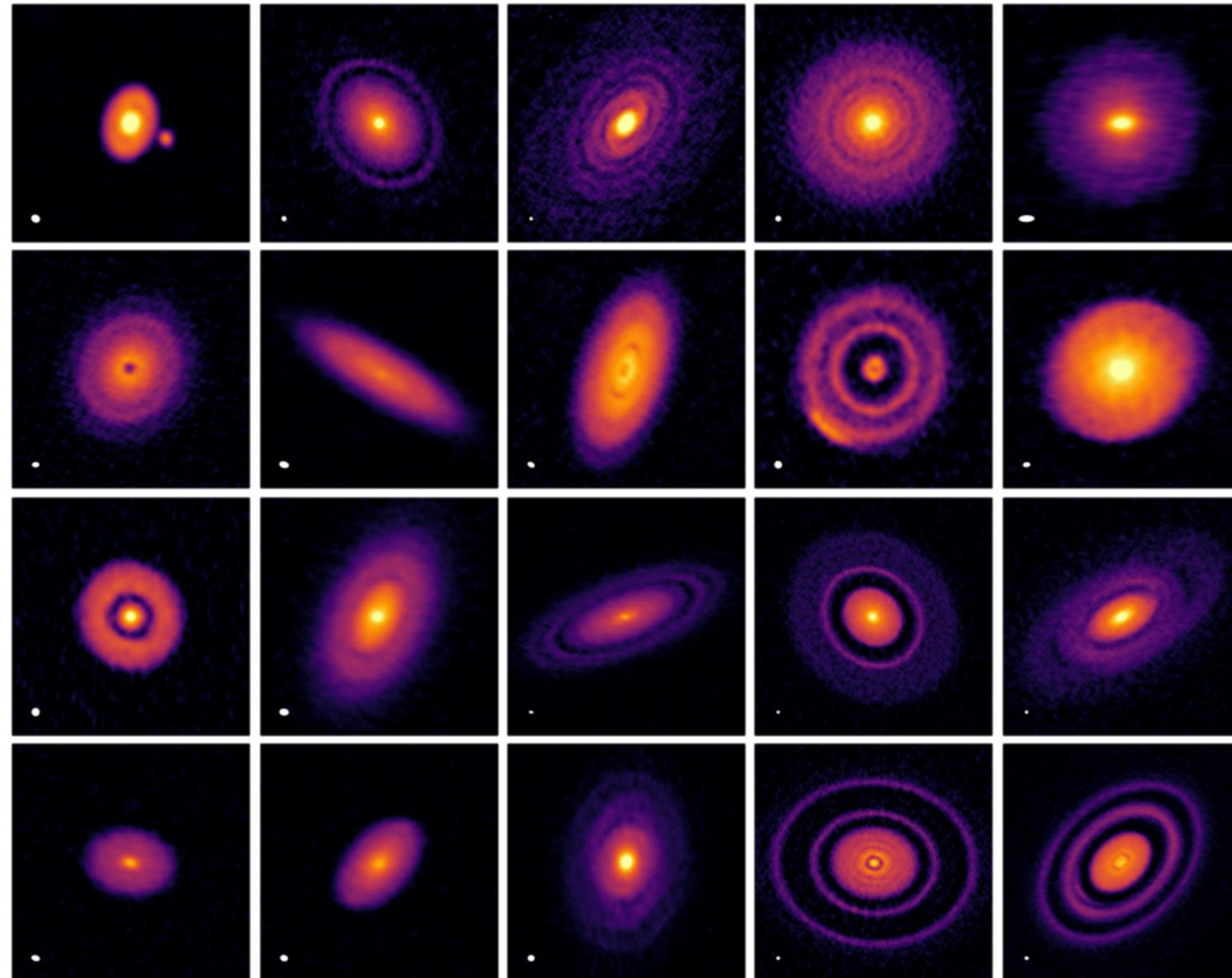
*Cacciapuoti et al. 2024a*

# Observational evidence of dust evolution

In protoplanetary disks :

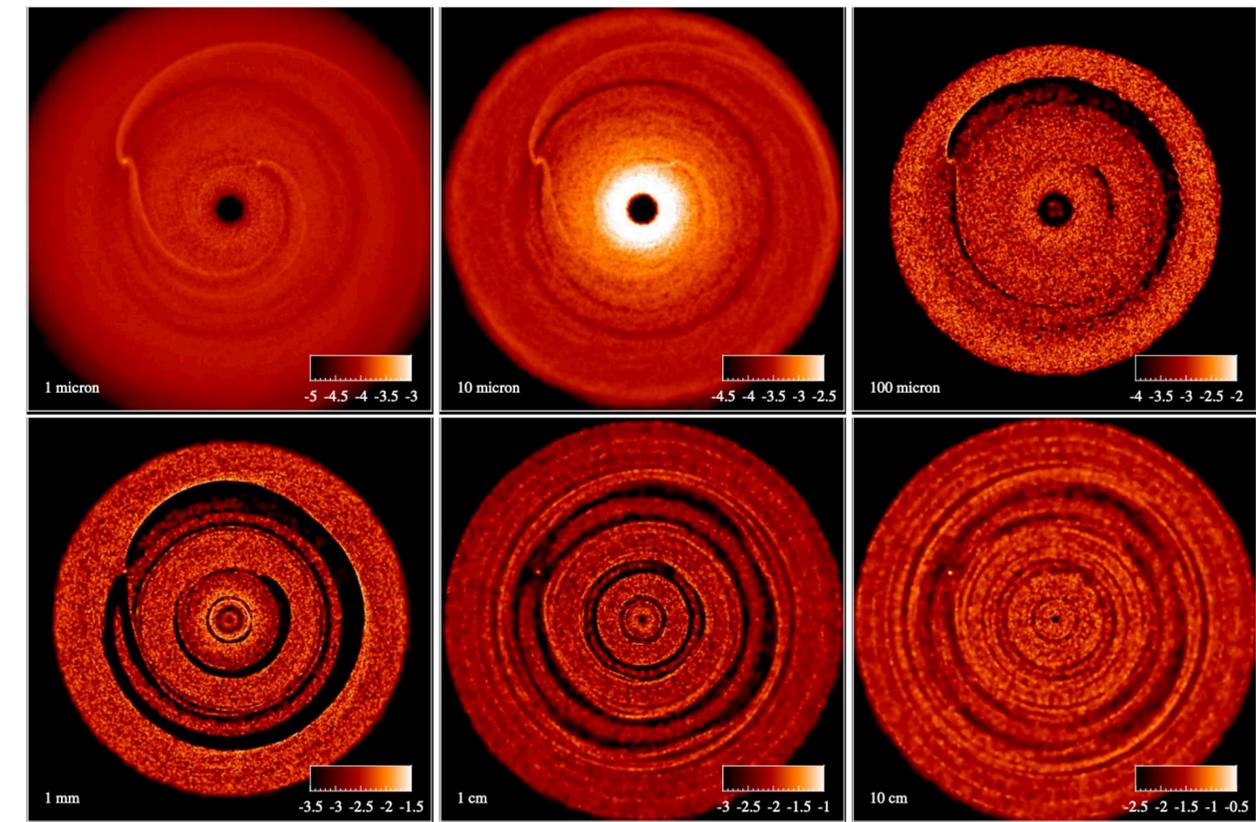
Dust emission in the mm in disks indicate grains of similar sizes.

—> Also, sub-structures and vertical settling could be indication that large  $St \sim 1$  are reached



ALMA observations of disks (DSHARP).

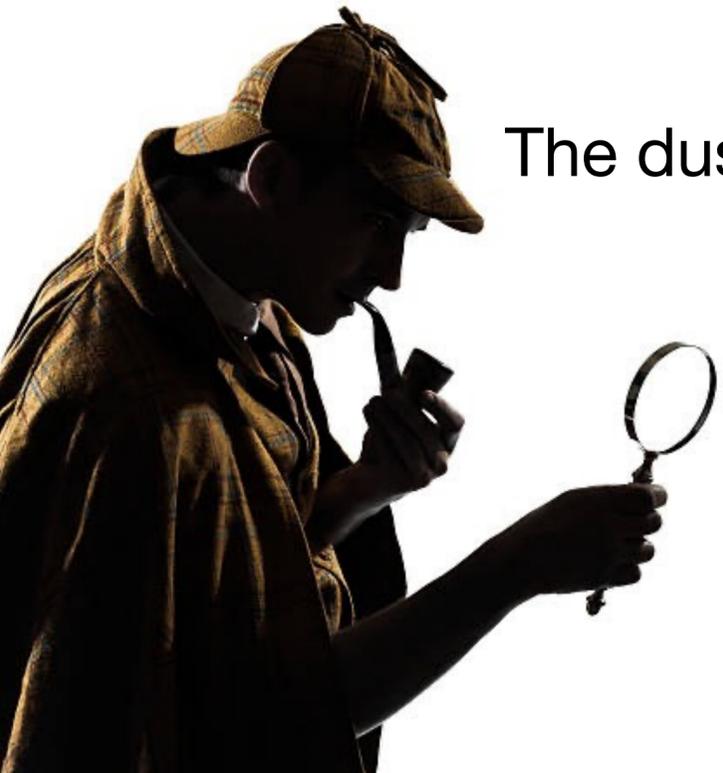
*Andrews et al. (2018)*



*Dipierro et al. (2015)*

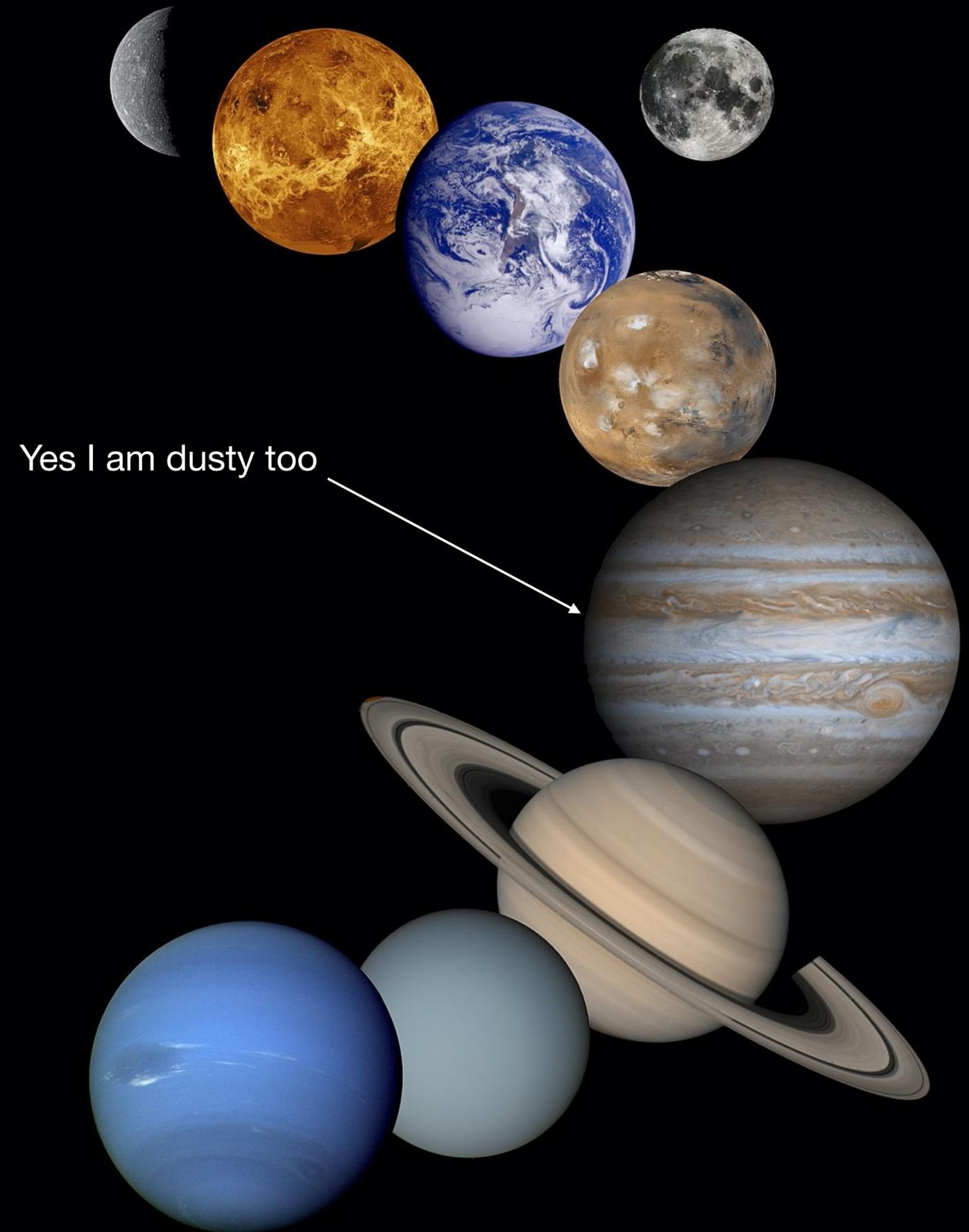
# Last evidence

I wonder what could that be.



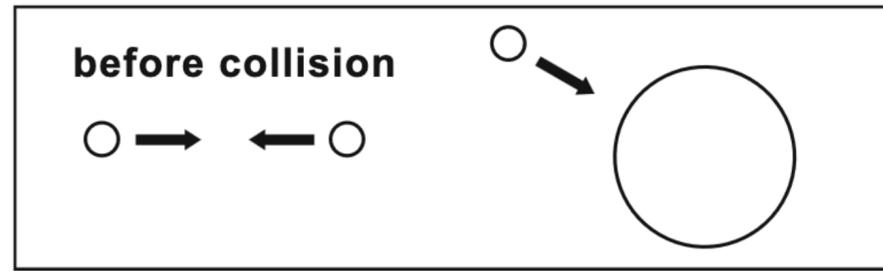
The dust under my bed ?

Thank you Sherlock



Yes I am dusty too

# The many outcomes of grain-grain collisions



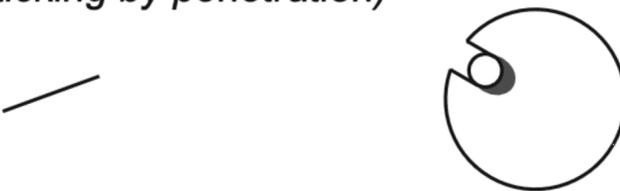
**S 1** (*hit & stick*)



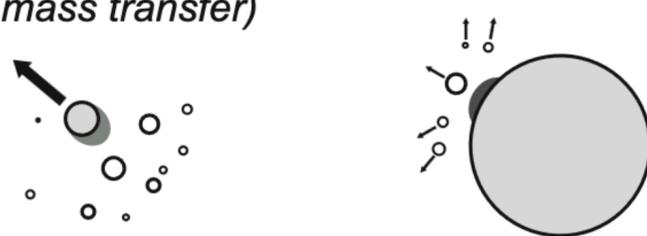
**S 2** (*sticking through surface effects*)



**S 3** (*sticking by penetration*)



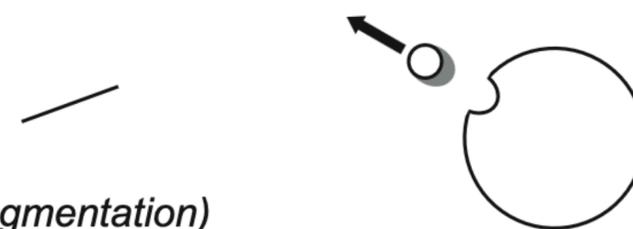
**S 4** (*mass transfer*)



**B 1** (*bouncing with compaction*)



**B 2** (*bouncing with mass transfer*)



**F 1** (*fragmentation*)



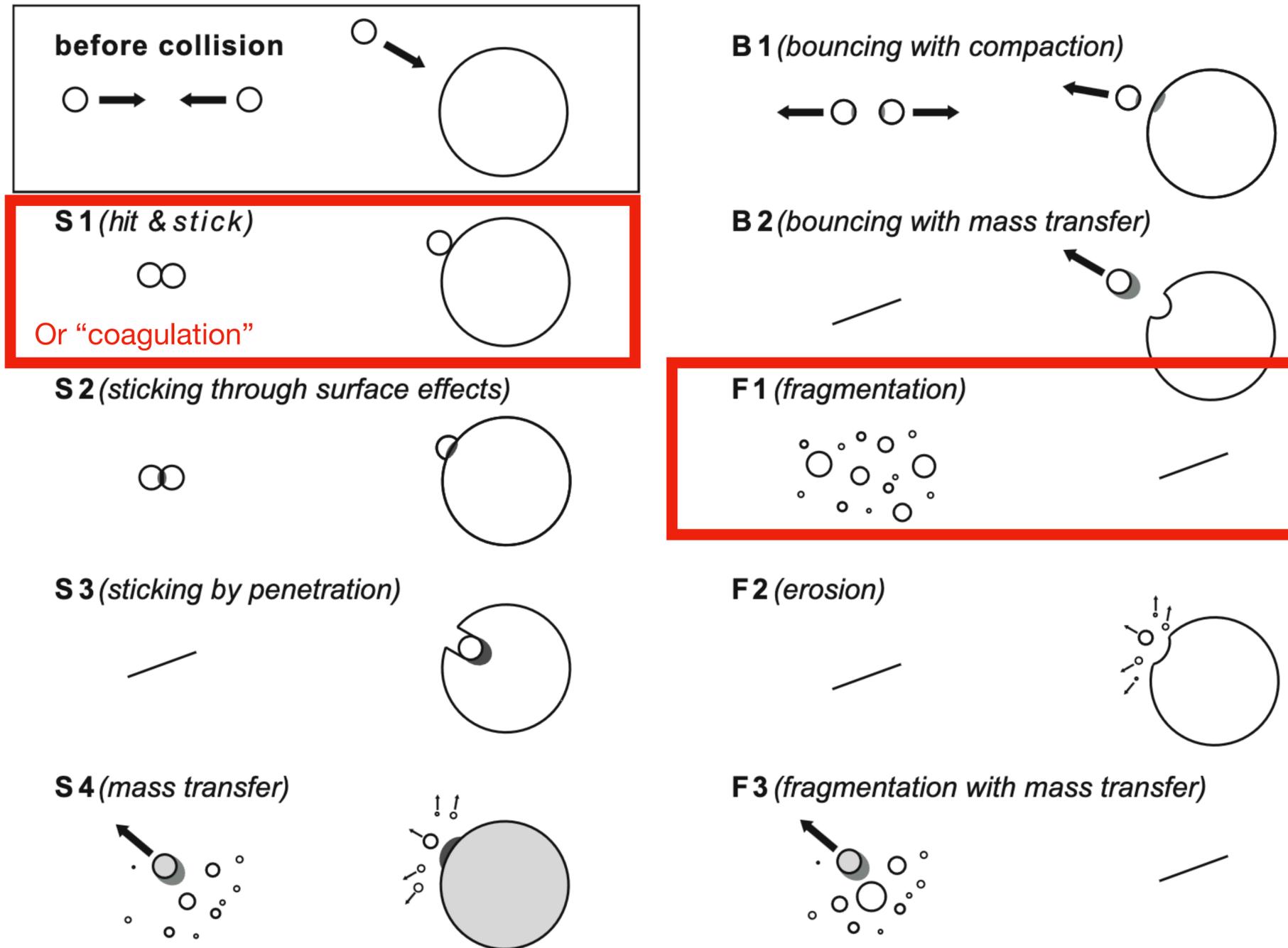
**F 2** (*erosion*)



**F 3** (*fragmentation with mass transfer*)



# The many outcomes of grain-grain collisions



# Dust evolution

## 1. How to model the dust coagulation ?

# Smoluchowski equation(s)

**Discrete (original) form (Smoluchowski 1916) :**

$$\frac{dn_k}{dt} = \underbrace{\frac{1}{2} \sum_{j=1}^{k-1} K_{j,k-j} n_j n_{k-j}}_{\text{Gain : Grain k formed from the collisions between smaller grains}} - \underbrace{n_k \sum_{j=1}^{\infty} K_{k,j} n_j}_{\text{Loss : k sticking with the other grain sizes (including k)}}$$

In the discrete form you solve  $\mathcal{N}$  (the larger it is the better) coupled equations.

**Coagulation Kernel** (where you put the physics it's a volume per unit time) :  $K_{i,j} = \pi(a_i + a_j)^2 \Delta v_{i,j}$

**Continuous form (Müller 1928) :**

$$\frac{\partial \tilde{n}}{\partial t} = \frac{1}{2} \int_0^m K(m', m - m') \tilde{n}(m') \tilde{n}(m - m') dm' - \int_0^\infty K(m', m') \tilde{n}(m') \tilde{n}(m') dm', \quad \tilde{n} = \frac{dn}{dm}$$

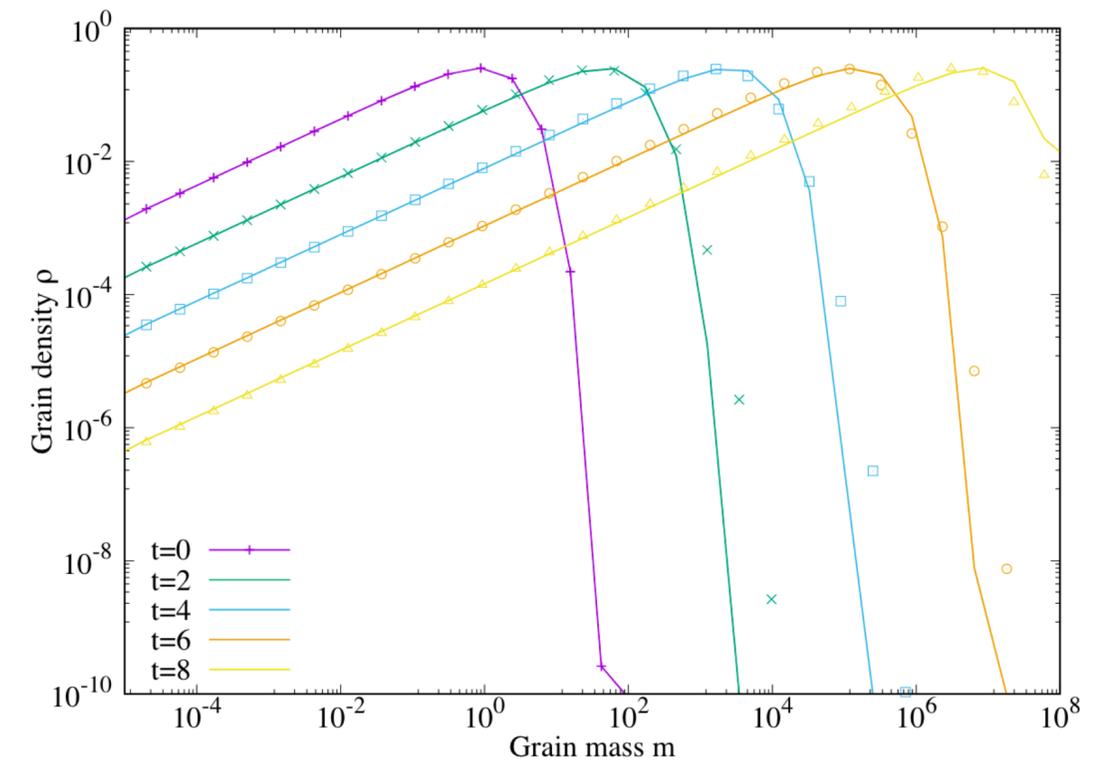
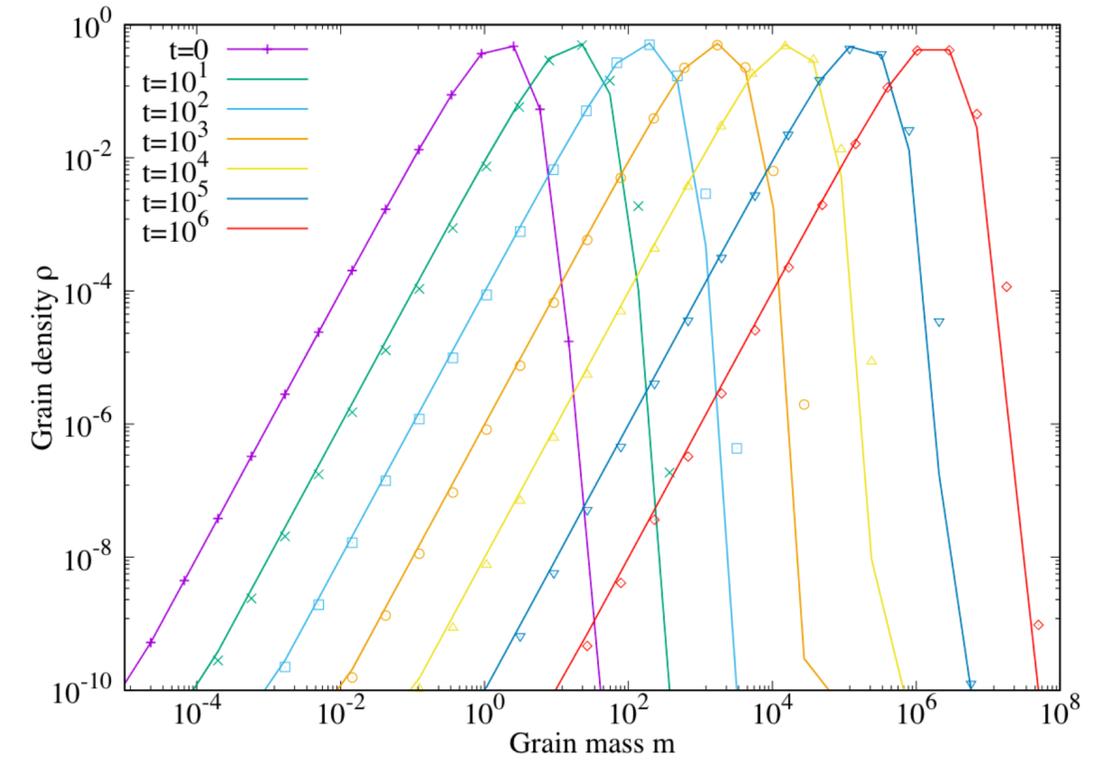
# Smoluchowski equation

No analytical solutions exist for :

- Constant kernels  $K_{i,j} = C$
- Additive kernel  $K_{i,j} = m_i + m_j$
- Multiplicative kernels  $K_{i,j} = m_i m_j$

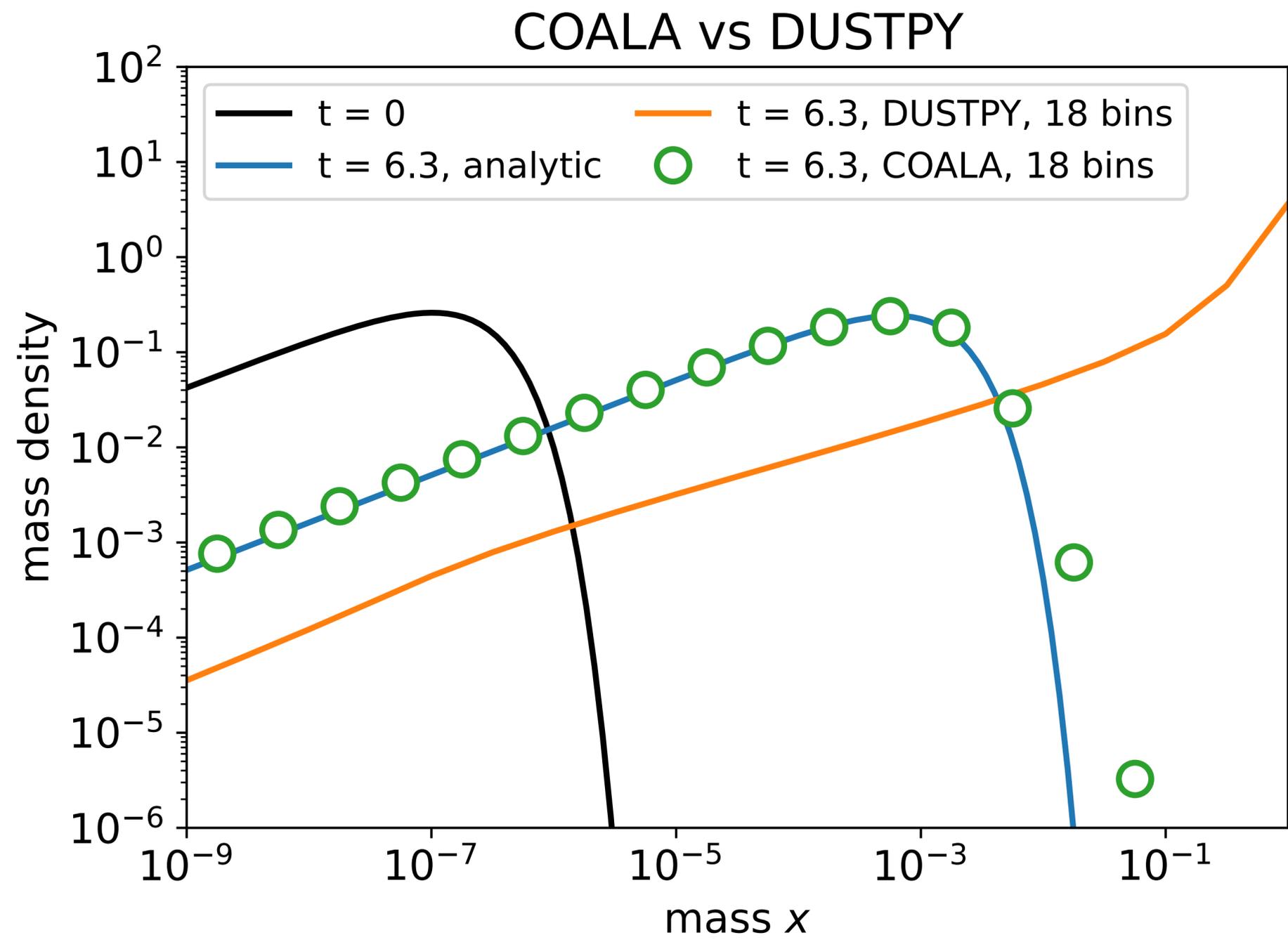
Self-similar solutions derived in *Menon & Pego (2004)*

==> *Can be used to test dust coagulation algorithms*



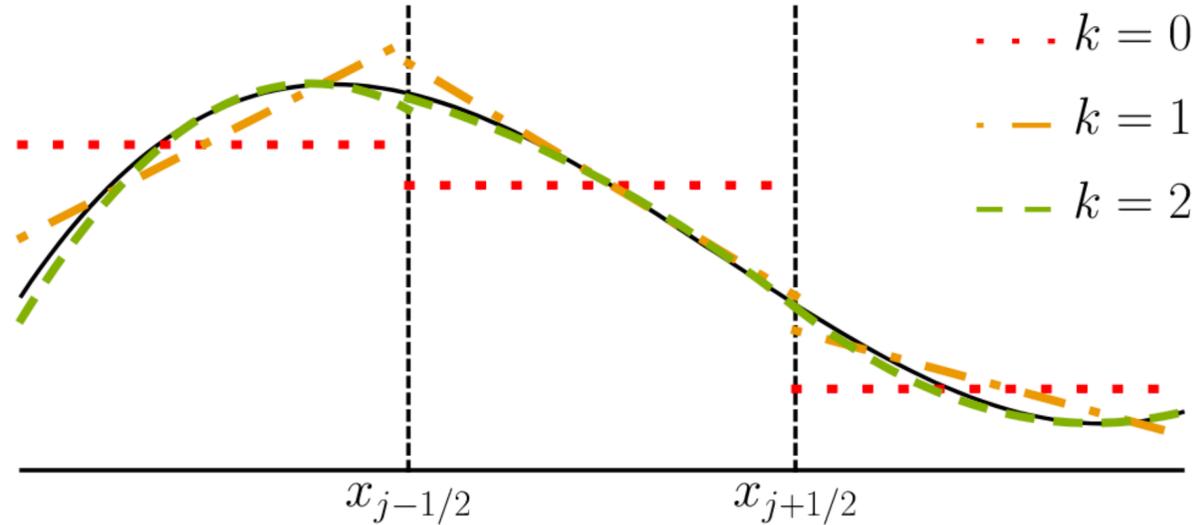
e.g. *Marchand et al. (2021)*

# Problem with solving the Smoluchowski equation numerically



Coala : [Lombart & Laibe 2021](#)

Dustpy : [Stammler et al. 2023](#)



Idea : polynomial approx of the distribution in a bin

**Most coagulation numerical methods need a lot of dust bins to converge** : this is a problem for 3D simulations

New algorithm by [Lombart & Laibe \(2021\)](#) uses a Galerkin discontinuous methods and can get a good agreement with a low number of bins

# An interesting way to overcome the cost of the Smoluchowski equation

If you assume a coagulation kernel that has the form :  $K(m, m') = f_{\text{gas}}(n_H, T, \text{etc.}) h(m, m')$

You can remove the gas properties from the integral (X -> dust to gas ratio) :

$$\frac{dX}{dt} = f_{\text{gas}} n_H I(X, a, t)$$

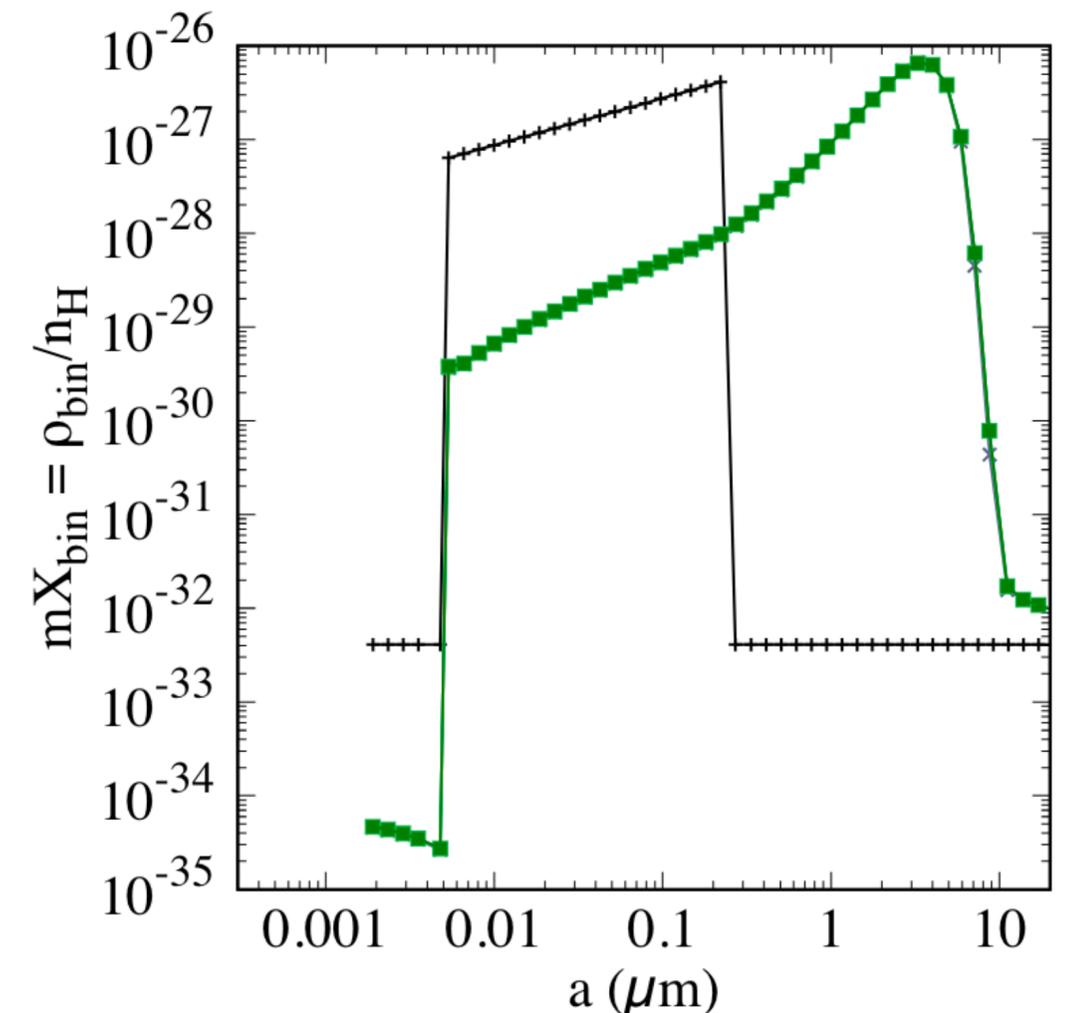
$$I(X, a, t) = \frac{1}{2} \int_0^m h(m - m', m') X(m - m', t) X(m', t) dm' - \int_0^\infty h(m, m') X(m, t) X(m', t) dm'$$

Coagulation becomes a 1D process :

$$d\chi = f_{\text{gas}} n_H dt \quad \longrightarrow \quad \frac{dX}{d\chi} = I(X, a, t)$$

**Caveats :**

- Only works for certain kernels (e.g turbulence)
- Incompatible with fragmentation



# Stepinski (or monodisperse) approach

$$\frac{da_{\text{grain}}}{dt} = \frac{a_{\text{grain}}}{3t_{\text{coag}}} \quad (\text{Stepinski \& Valageas 1996})$$

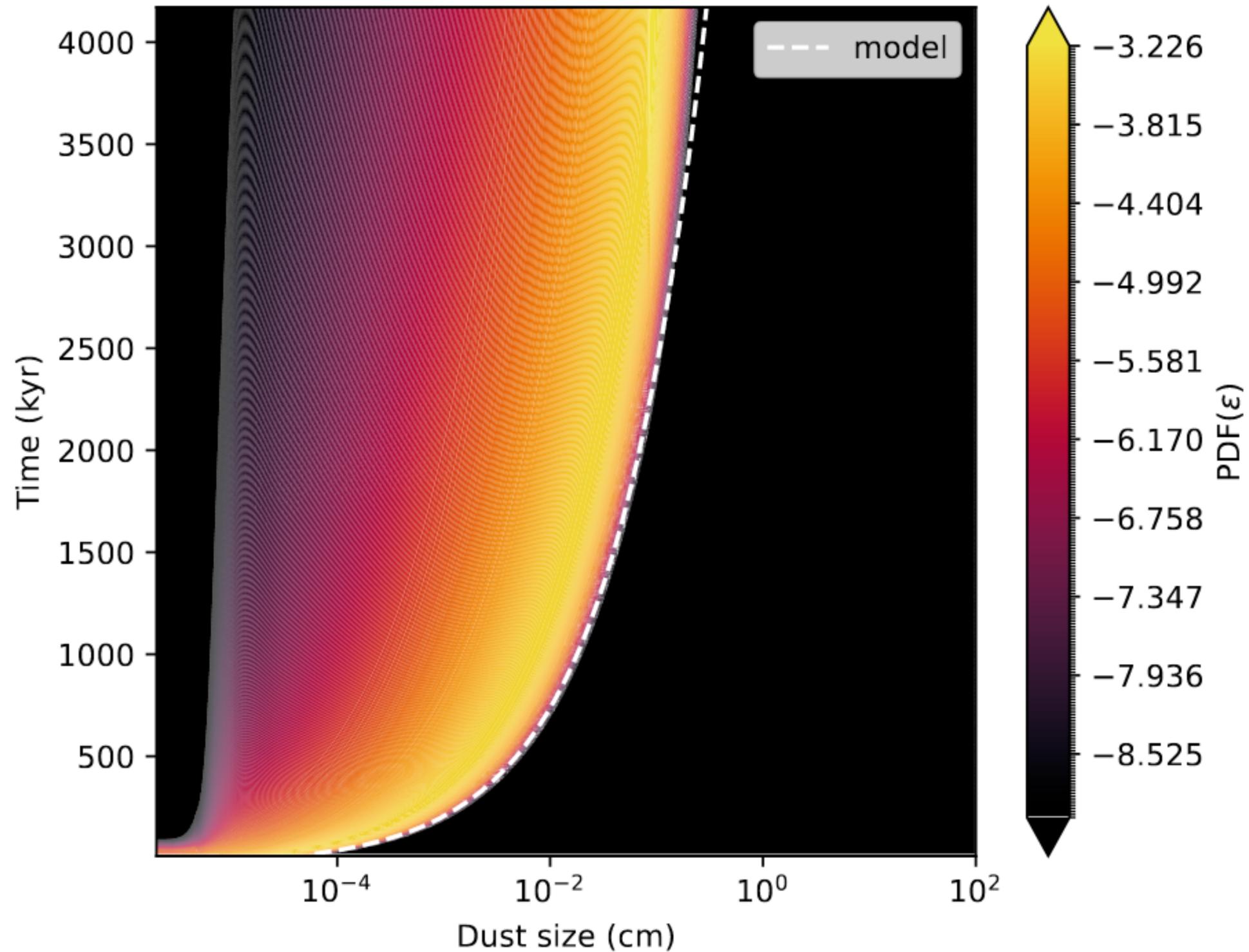
Timescale for coagulating from a grain :  $t_{\text{coag}} = \frac{1}{\pi n_{\text{d}} a_{\text{grain}}^2 \Delta v}$  (*Draine 1985*)

## Hypothesis :

- > Grain collisions happen only for grain of similar sizes (which is wrong of course)
- > Only growth by turbulence or brownian motions (which are the only velocity source able to coagulate grain of similar sizes) — No systematic drift included (hydro drift, ambipolar drift)

**Simplified but powerful to get the peak of the distribution**

# Stepinski approach



Vallucci-Goy et al. (in prep)

# Dust evolution

## 2. Sources of grain-grain collisions

# Sources of grain-grain collisions (velocities)

“Hydrodynamical” drift :

$$\Delta \vec{v} = \vec{v}_g - \vec{v}_d$$

this one comes from solving the fluid equations. (remember the 1st part of the course)

“Ambipolar” drift or **magnetic drift** : Same, but for charged grains  $\Delta \vec{v} = \frac{c}{|\vec{B}|} \left( \frac{\Gamma_d^2}{1 + \Gamma_d^2} \vec{E}_b \times \vec{b} + \frac{\Gamma_d}{1 + \Gamma_d^2} \vec{E}_{b,\perp} + \Gamma_d \vec{E}_{b,\parallel} \right)$

$$c \vec{E}_b = \eta_o \nabla \times \vec{B} - \frac{\eta_A}{|\vec{B}|^2} \left( (\nabla \times \vec{B}) \times \vec{B} \right) \times \vec{B} + \frac{\eta_H}{|\vec{B}|^2} (\nabla \times \vec{B}) \times \vec{B}$$

--> Grains of different sizes have a different coupling with the magnetic field.

--> Small grains stick to the field and ‘big’ grains follow(-ish) the gas.

In principle, these two velocities are parts of the same systematic drift between the gas and the dust that should be obtained by solving the full **multifluid-MHD (RT, gravo, etc.) gas and dust (and electrons and ions.. CRs?) equations**

So far we sum the two contributions (with a simplified formula for the ambipolar drift)

# Sources of grain-grain collisions (velocities)

**Brownian motions :**  $\Delta v_{B,i,j} = \sqrt{\frac{8k_B T}{\pi}} \sqrt{\frac{m_i + m_j}{m_i m_j}}$

If we fix  $m_j$ ,  $\Delta v_{B,i,j}$  diverges with a decreasing  $m_i$ . Brownian motions are efficient at sticking small grains to big grains. (So not really to make large grains)

**Turbulent motions** (*Völk et al. 1980; Ormel & Cuzzi 2007*) more general derivation in *Gong et al. (2020;2021)* :

Grains of different sizes couple with eddies of sizes comparable to their stopping length. So the dynamics of the grains depends on how the stopping times compare with the injection timescale of the turbulence and its dissipation timescales.

Expression complicated (you can find in *Ormel & Cuzzi 2007*).

**Important points :**

- Turbulent can make grains of similar sizes collide so it is efficient at making large grains
- Expression of *Ormel & Cuzzi (2007)* widely used but remember that these are analytical expressions of something very complicated and neglect some key effects : the dust back-reaction (through collisions or magnetic fields) and changes of dust concentration (but we know dust can concentrate in the turbulence)

# Dust evolution

**Any questions ?**

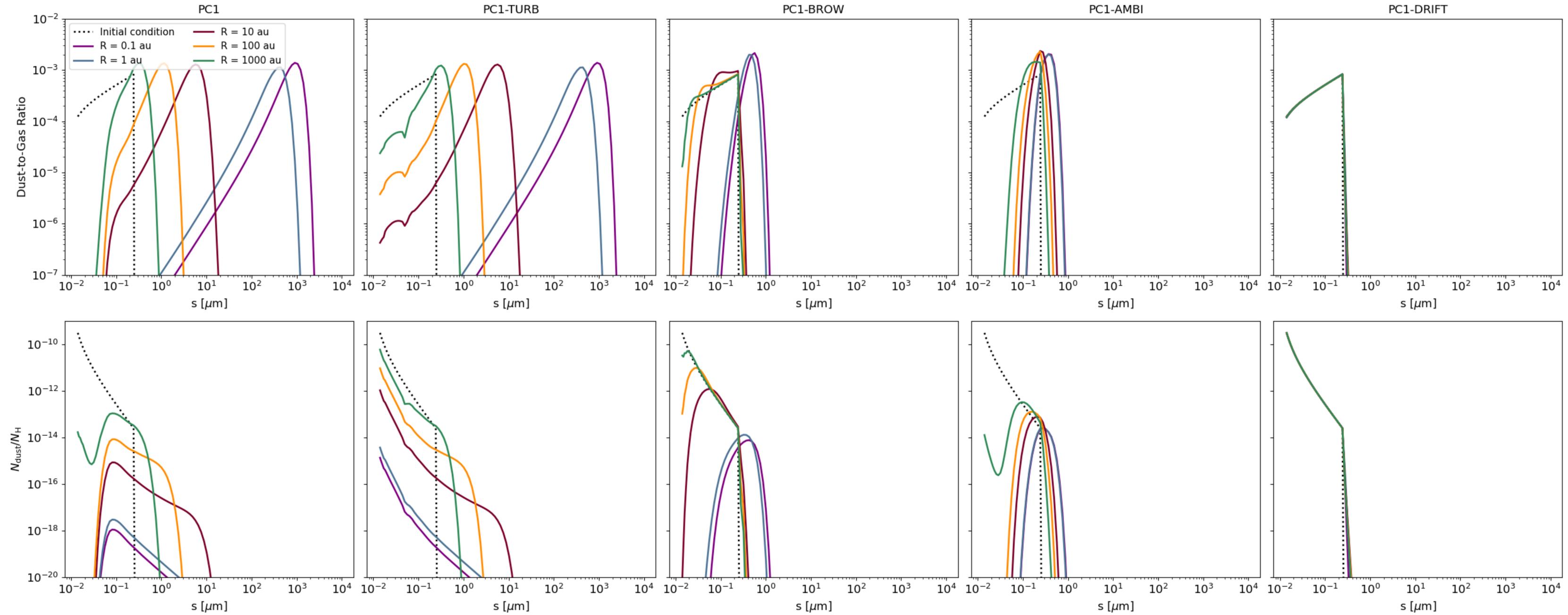
# Dust evolution

## 3. Coagulation/fragmentation in astrophysical environments

# Dust evolution

## Protostellar collapse

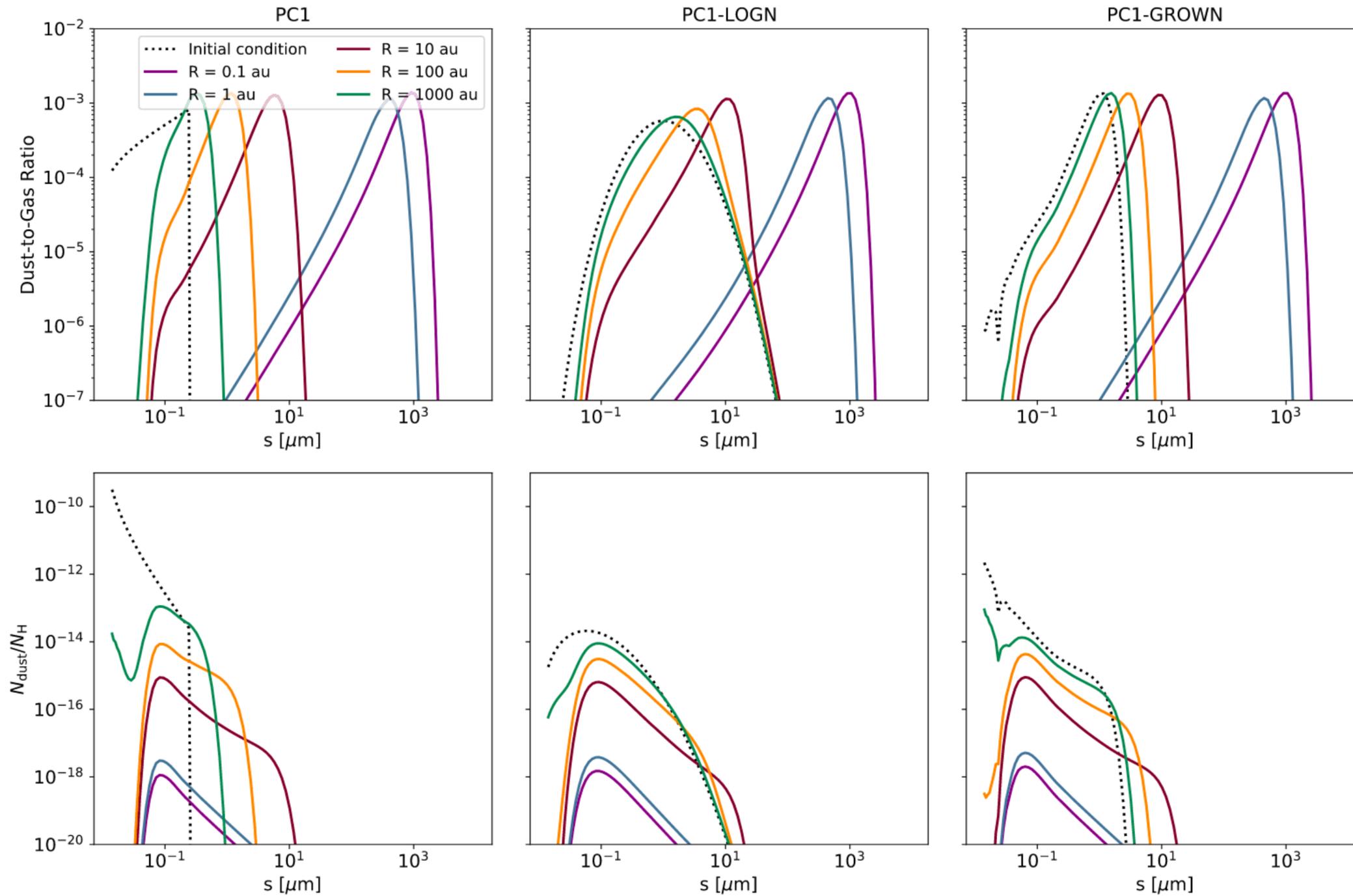
# Sources of grain-grain collisions (velocities)



*Lebreuilly et al. 2023a* see also :

*Guillet et al. 2020; Silsee et al. 2020, Vallucci-Goy et al. sub*

# Memory loss of coagulation

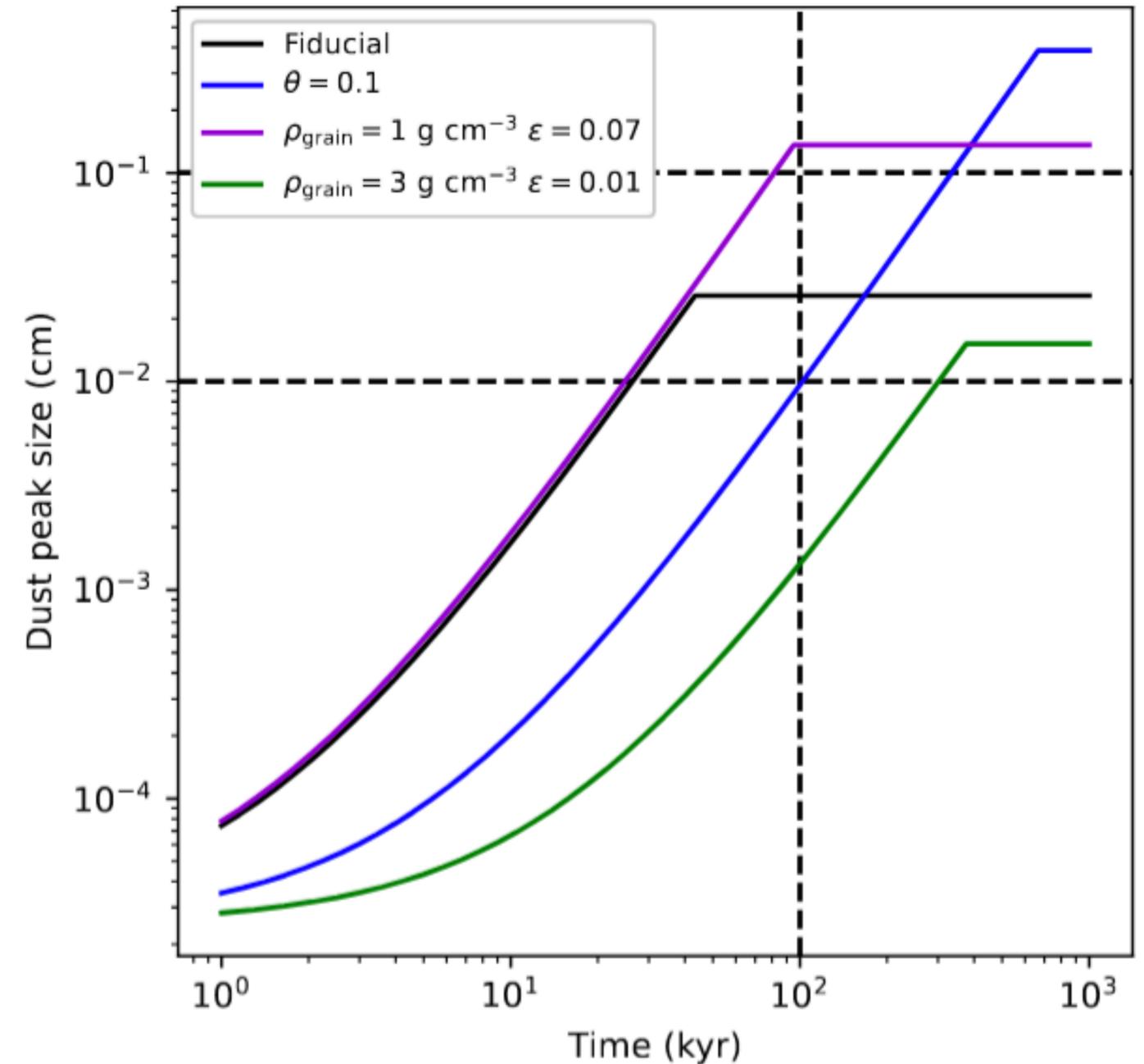
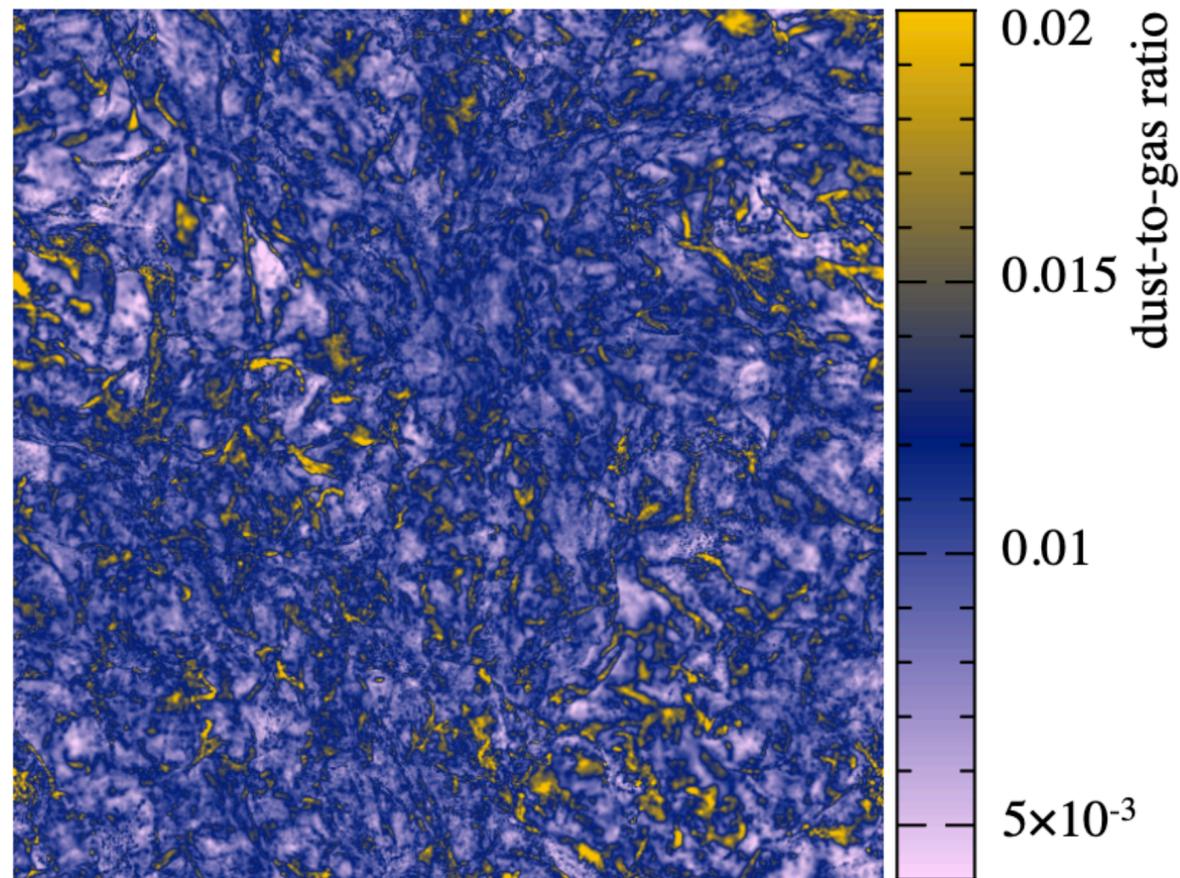


# Is there a timescale issue for coagulation ?

Growing millimetre grains in protostellar envelopes is hard : but not impossible

## 3 ingredients are needed :

- Low velocities (or high fragmentation threshold)
- Porous grains (which accelerates the coagulation)
- Some dust clumping (small scale dust clumping is fine)

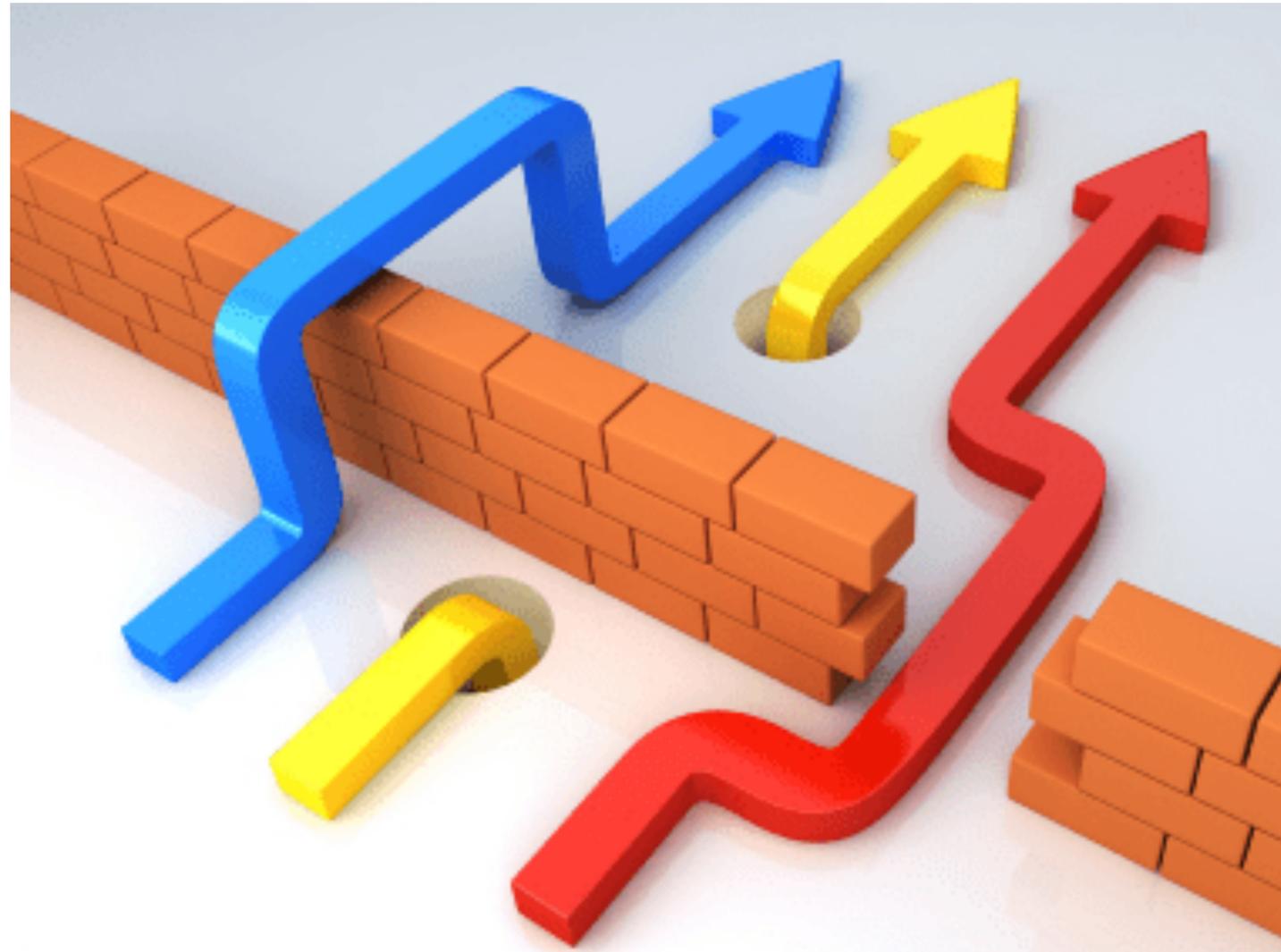


Tricco et al. 2016; Hopkins & Lee (2016), Lee, Hopkins & Squire (2017) ; Commerçon et al. 2023

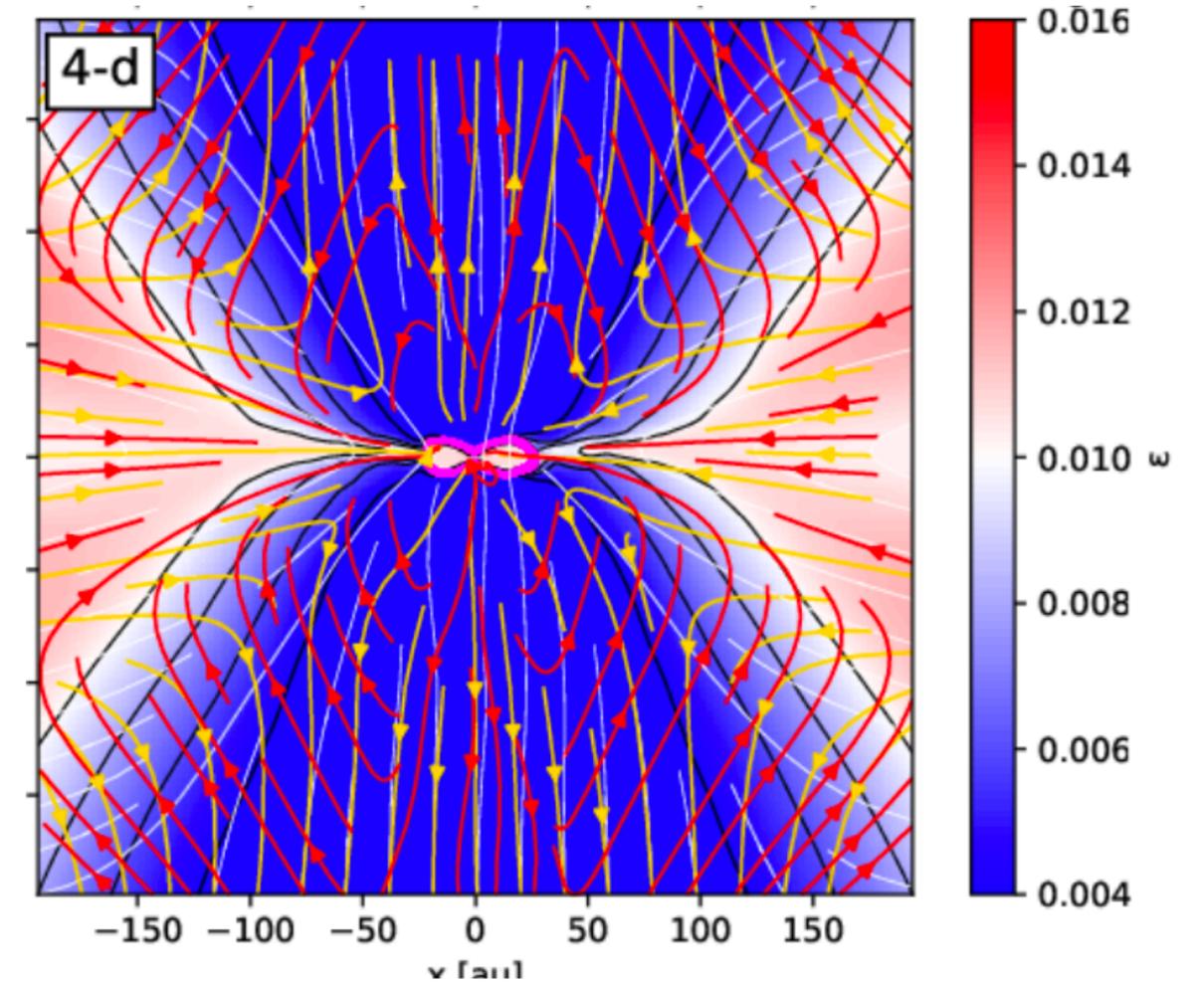
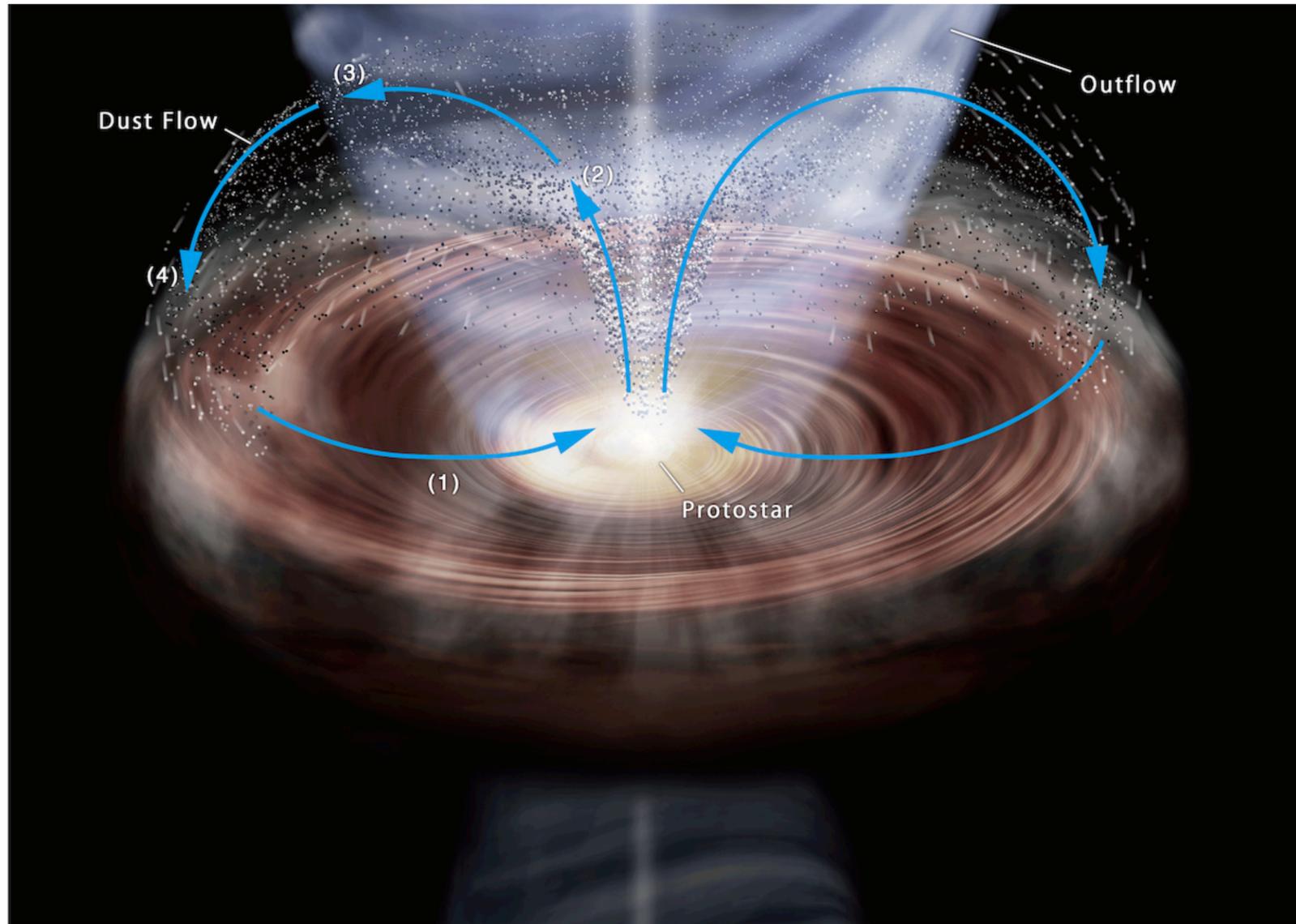
Vallucci-Goy et al. (in prep)

# Protoplanetary disks

## Overcoming the radial drift barrier



# “Ash-fall” phenomenon

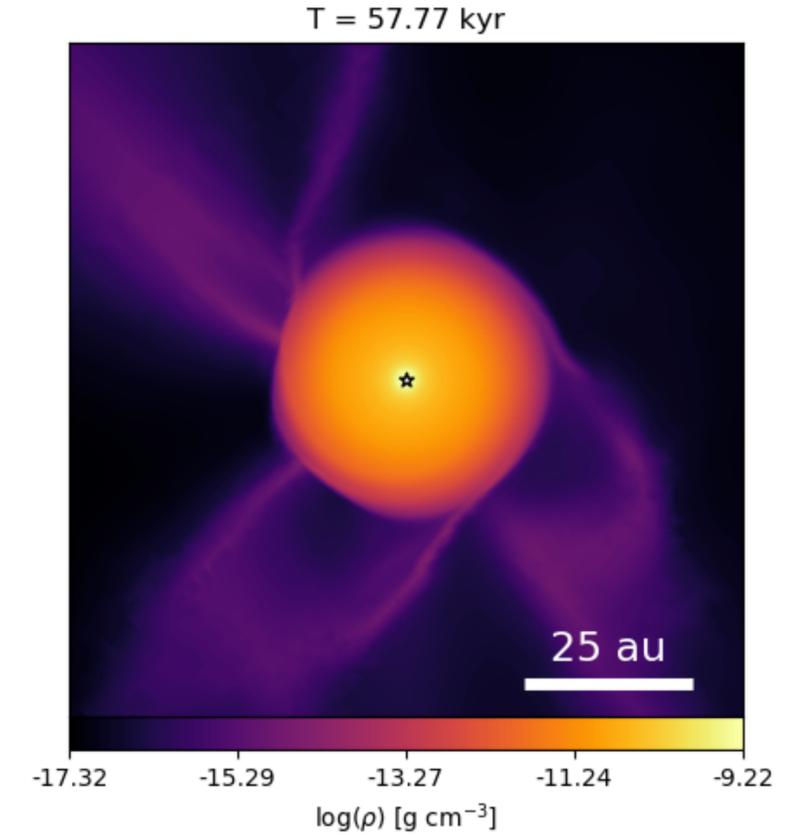
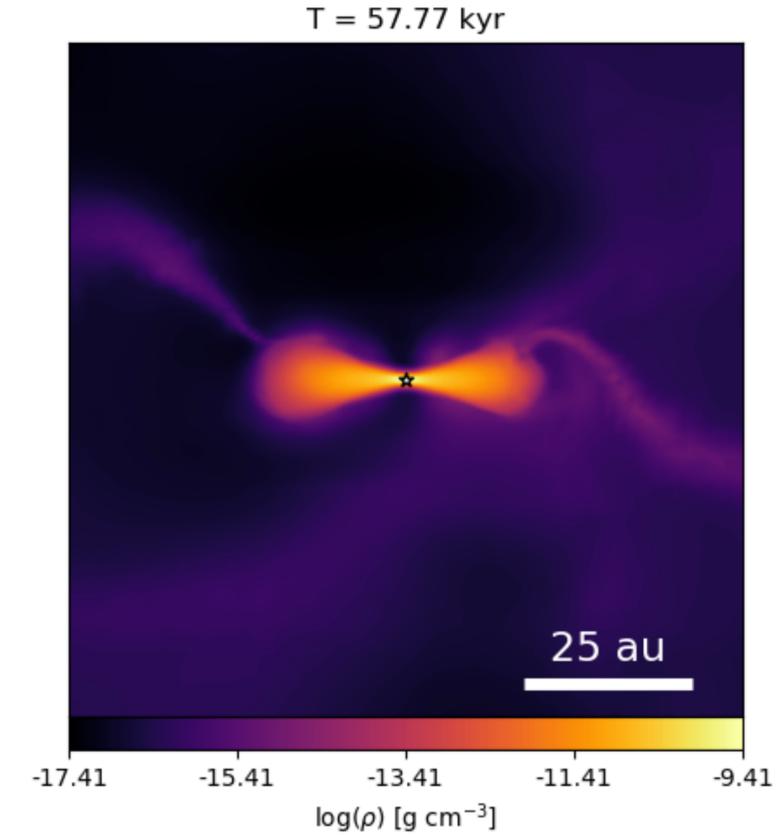
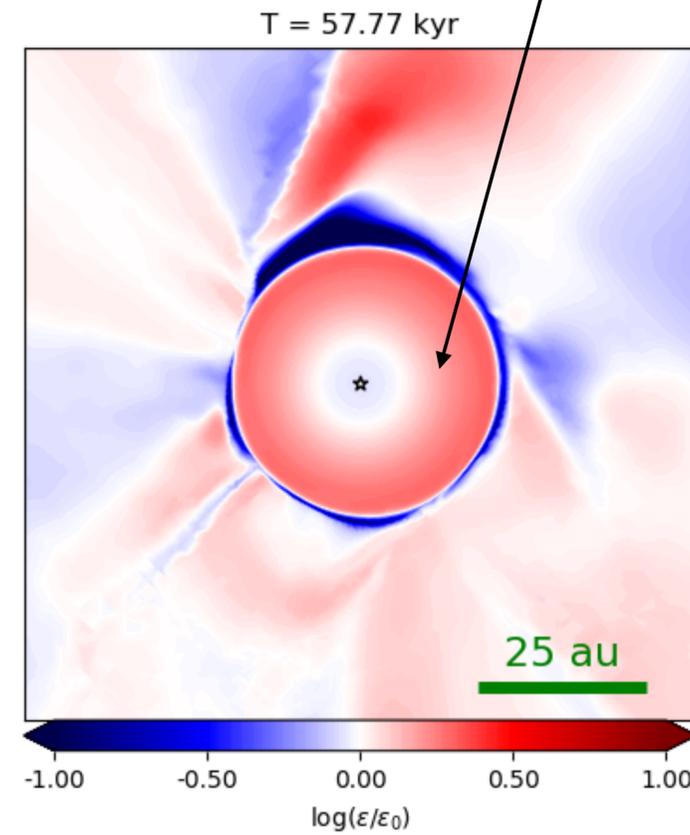
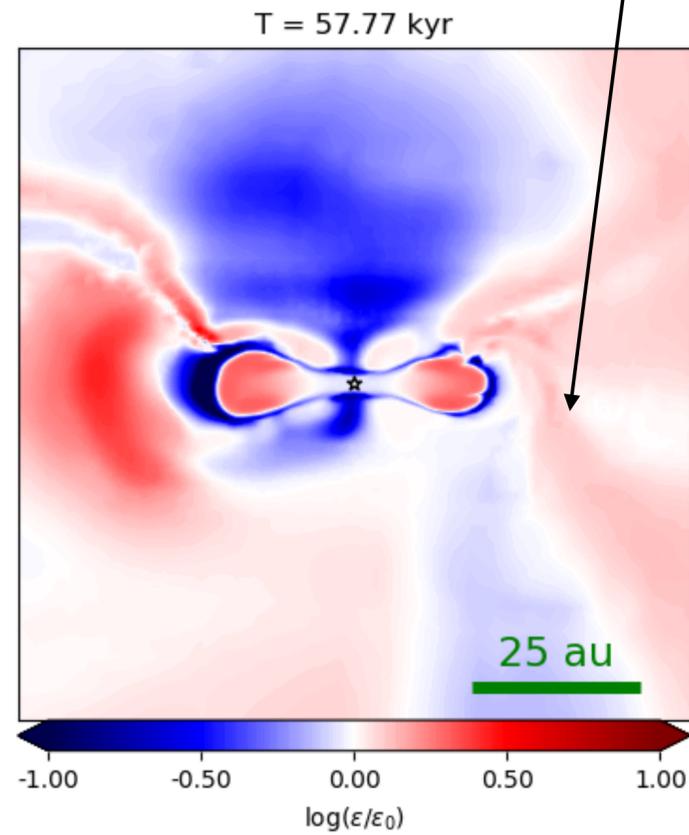


$$\frac{\partial \rho_d a_{\text{grain}}}{\partial t} + \nabla \cdot \left( \rho_d a_{\text{grain}} \vec{v}_{\text{grain}} \right) = \frac{\rho_d a_{\text{grain}}}{t_{\text{coag}}}$$

# Dust replenishment from the outer edge of the disk

Smaller grains but lower density

Larger densities but also larger grains



# Coagulation instability : in disks (but not only?)

## Mechanism of Coagulation Instability

**Coagulation**

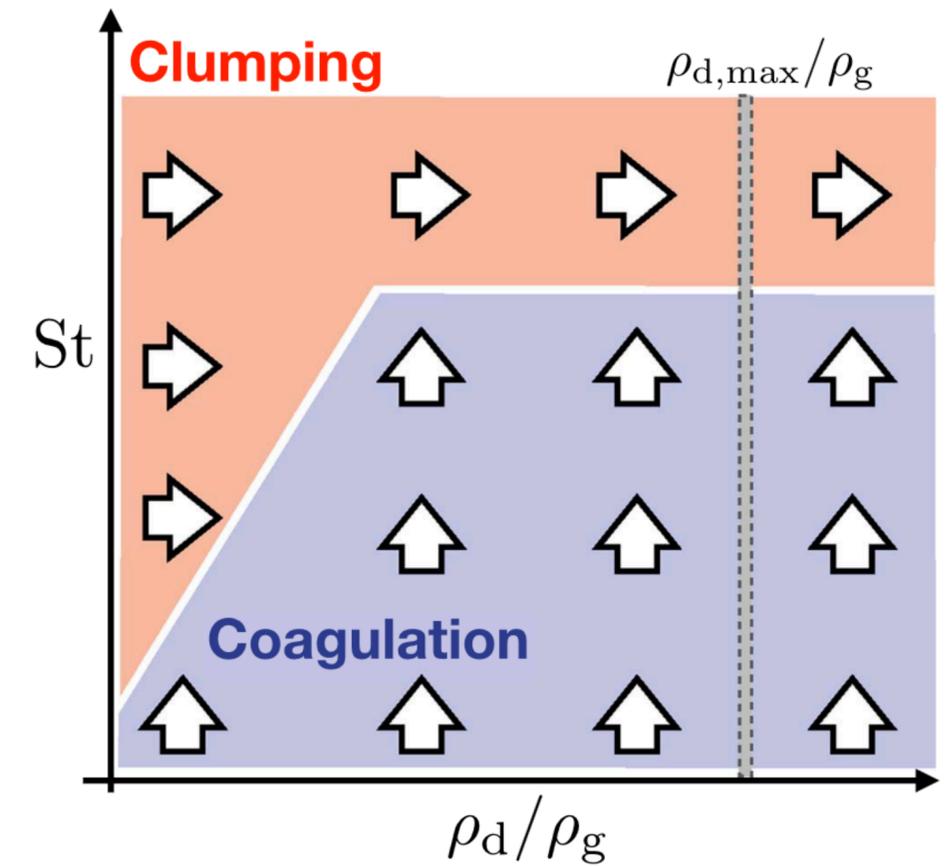
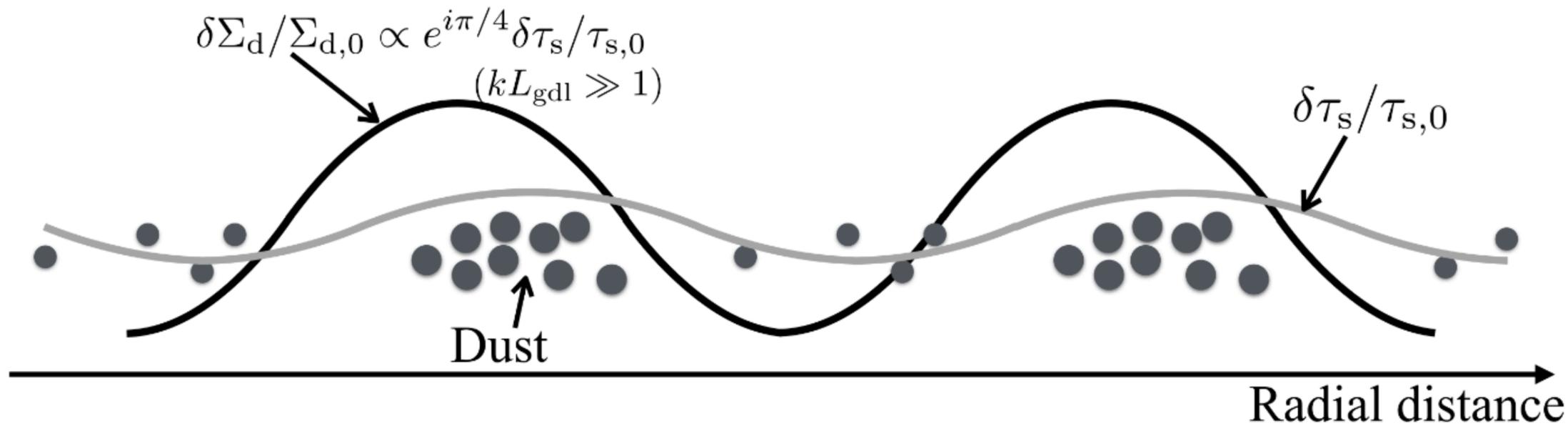
$$\delta\Sigma_d \nearrow \Rightarrow \delta\tau_s \nearrow$$

**Traffic jam**

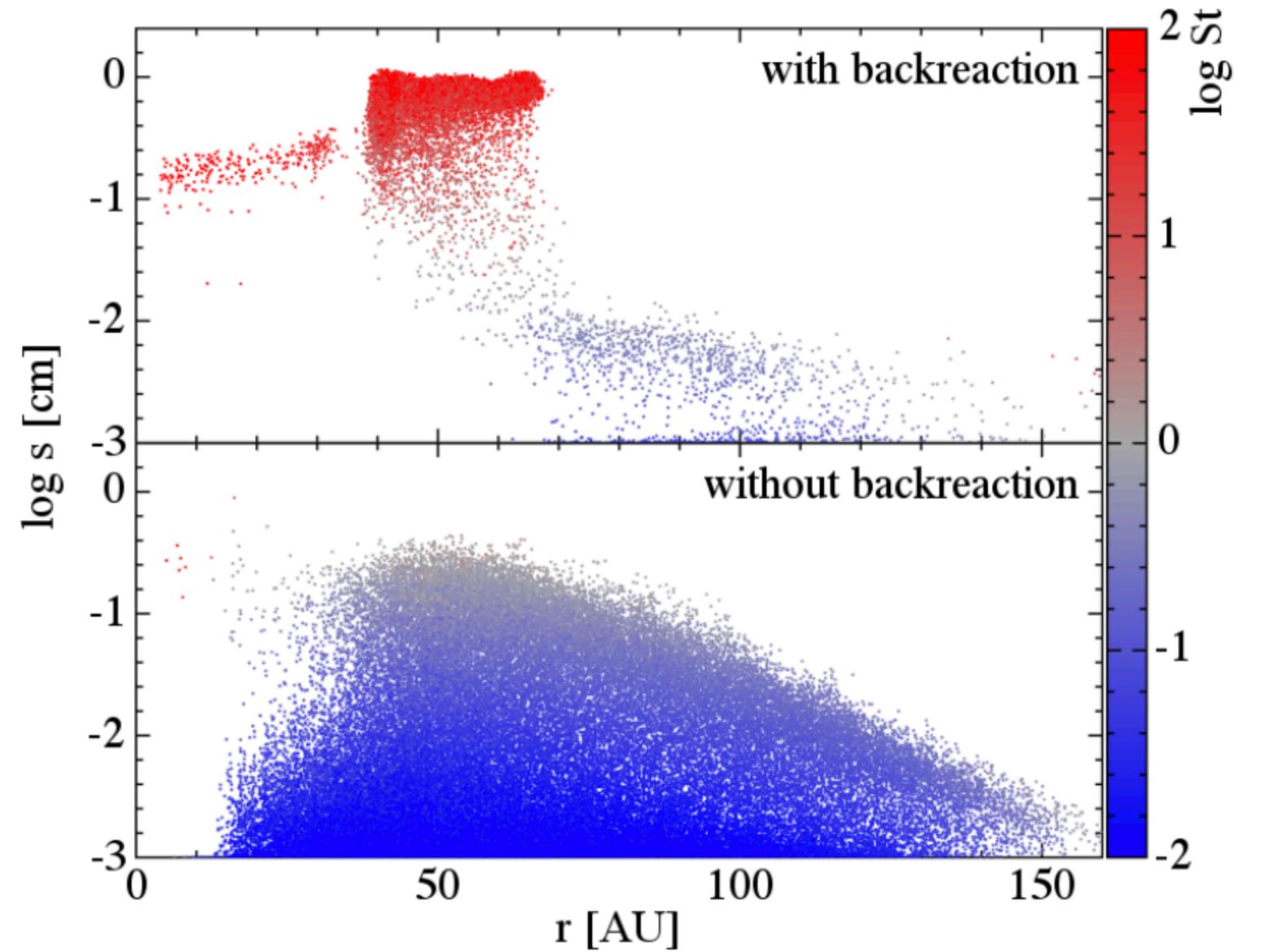
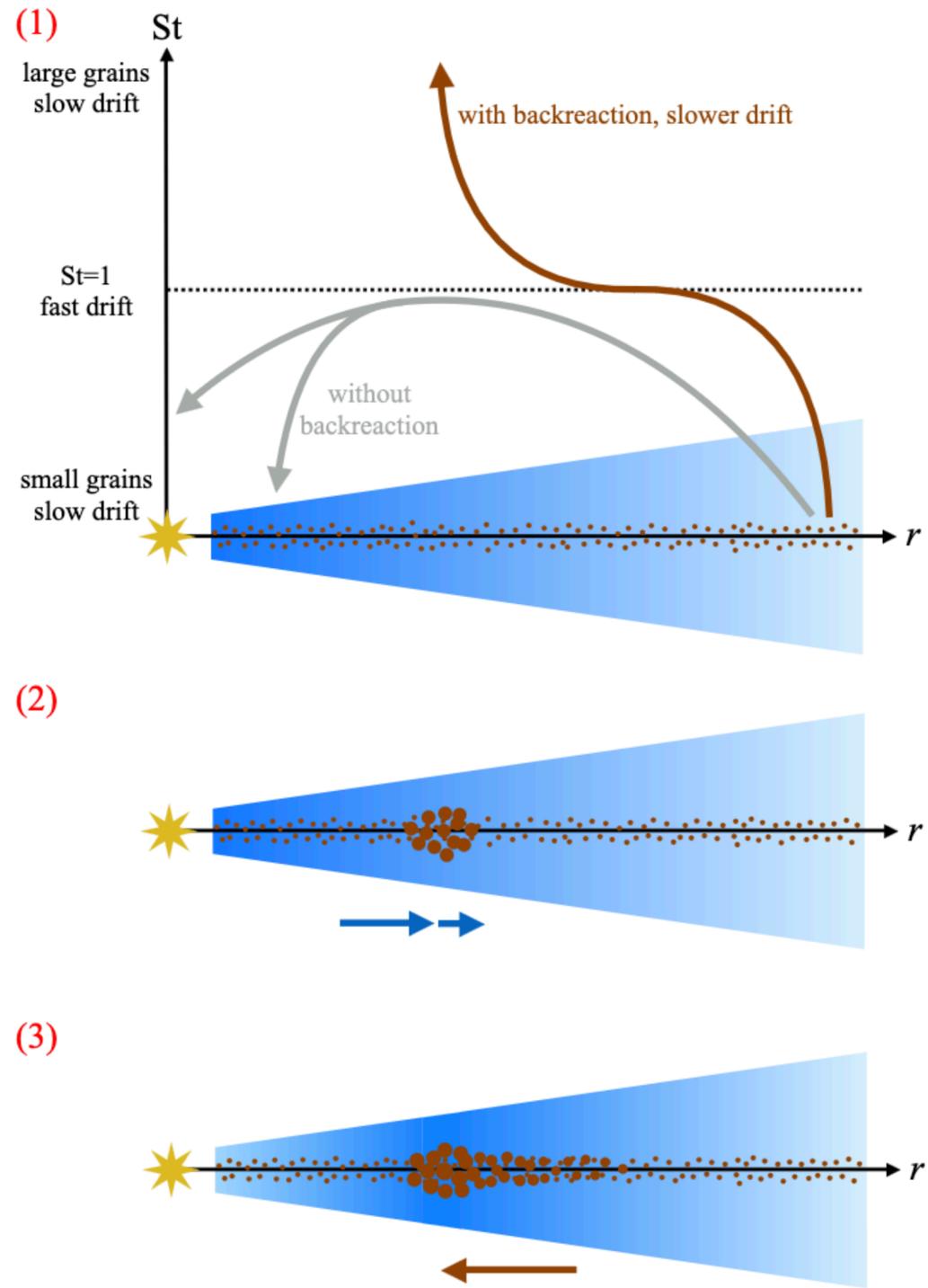
$$\delta\tau_s \nearrow \Rightarrow \delta\Sigma_d \nearrow$$

×

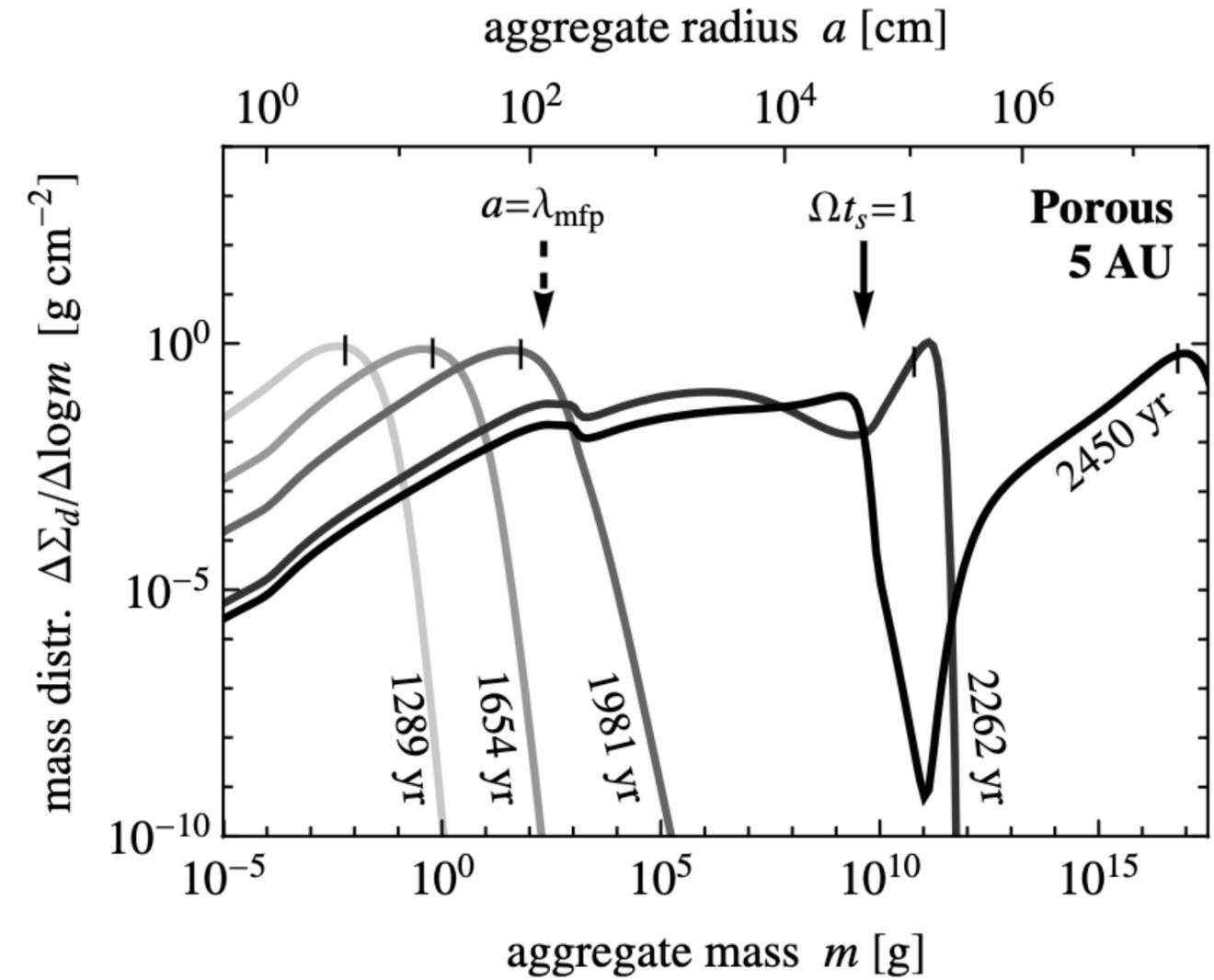
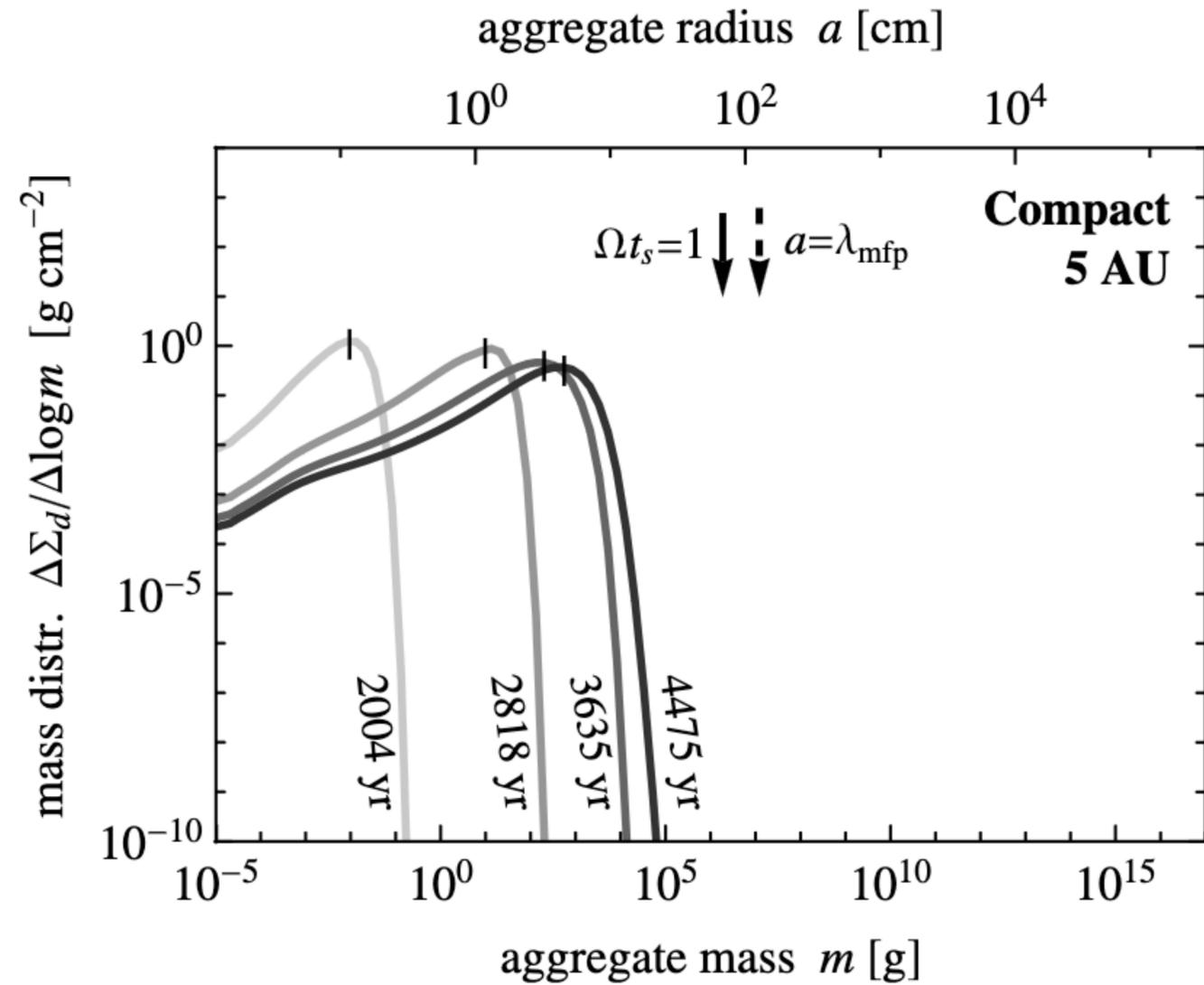
$$\delta\Sigma_d / \Sigma_{d,0} \propto e^{i\pi/4} \delta\tau_s / \tau_{s,0} \quad (kL_{gdl} \gg 1)$$



# Self-induced dust traps



# Porous grains could overcome the radial drift barrier



Porous grain might grow fast enough to break the drift barrier (but they must not break) :  $t_{\text{coag}} \propto \sqrt{\rho_{\text{grain}}} \sqrt{a_{\text{grain}}}$  in the turbulent case

# More barriers...



**Dammit...**



# Fragmentation

If the grain-grain collision becomes too energetic, then the bonds between the ‘monomers’ constituting the grains can break.

—————▶ **This is fragmentation**

Experimentally (*Dominik & Tielens 1997*, see also *Ormel et al. 2009*) a collision between a grain  $i$  and  $j$  will lead to completed destruction of both grains when :

$$E_{\text{kin},i,j} > E_{\text{frag},i,j} = 5N_{\text{mono}} E_{\text{br}}$$

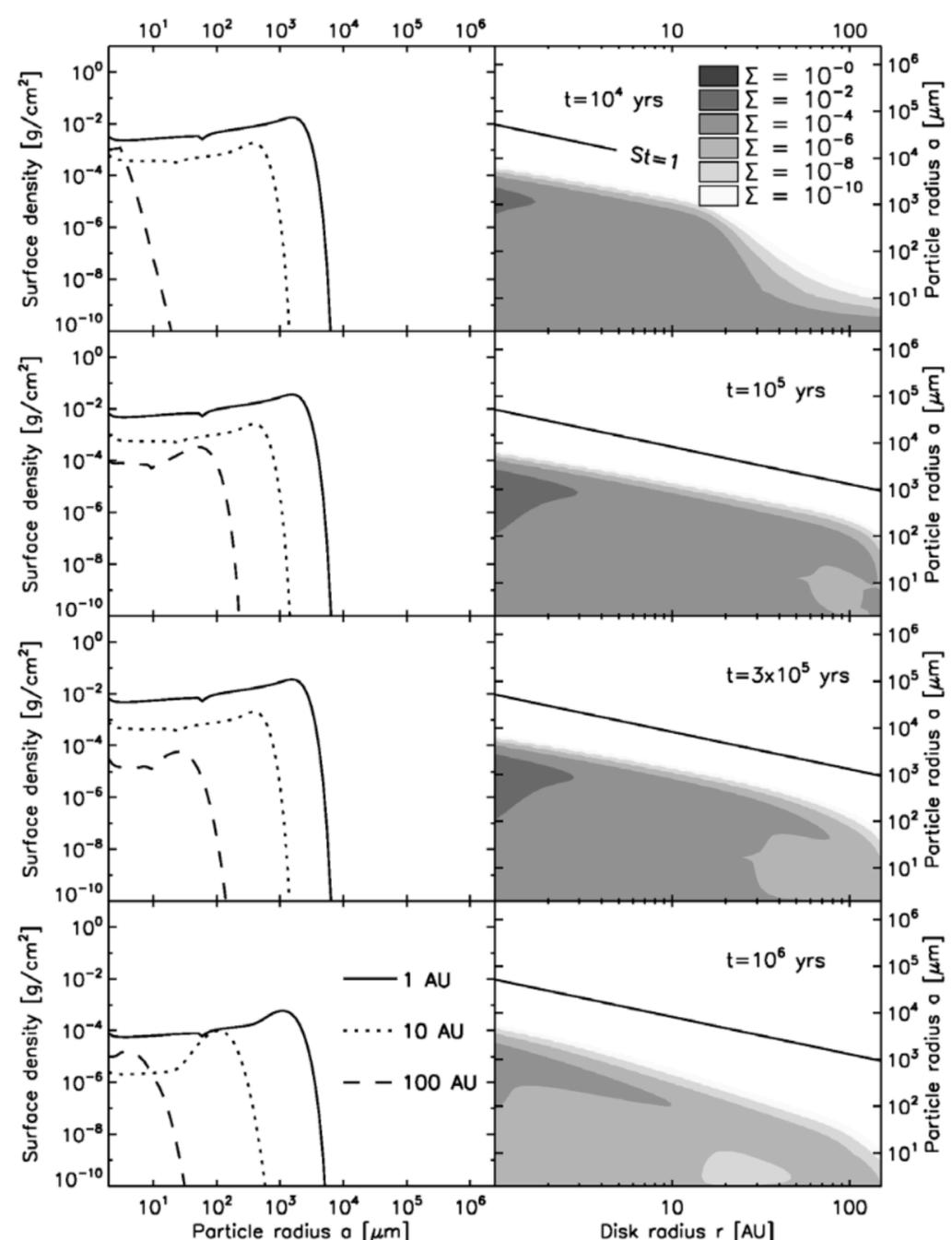
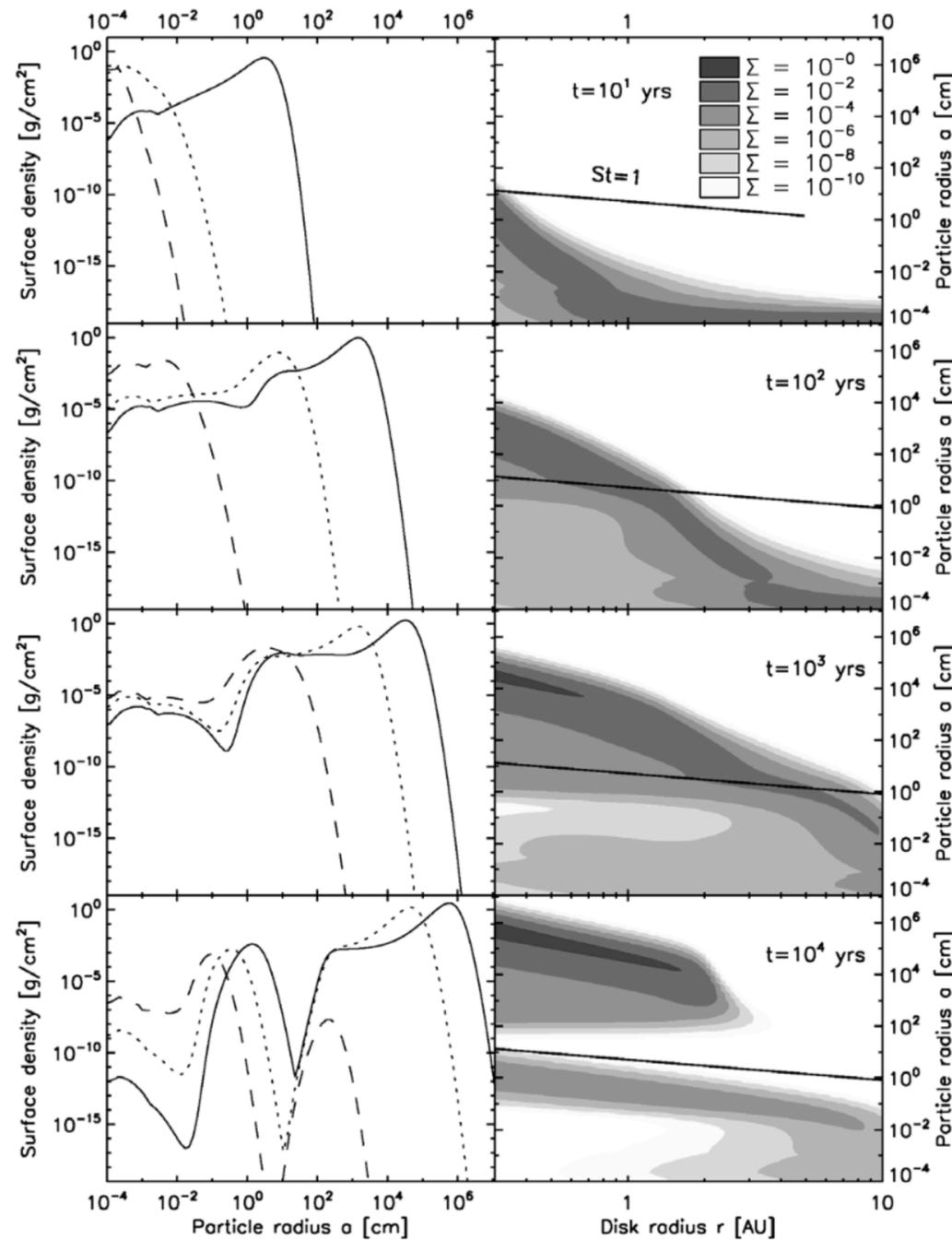
Or

$$\frac{1}{2} \frac{m_i m_j}{m_i + m_j} \Delta v_{\text{frag},i,j}^2 = 5 \left( \frac{m_i + m_j}{m_{\text{mono}}} \right) E_{\text{br}}$$

If we assume that the dust grains is composed of monomers that all have the same size (or a distribution see Kawasaki et al. 2022) we have

$$\Delta v_{i,j,\text{frag}} = \left( 1 + m_j/m_i \right) \sqrt{\frac{10m_i E_{\text{br}}}{m_j m_{\text{mono}}}} \longrightarrow \text{the fragmentation velocity is not constant but depend on the mass ratio between the grains}$$

# Fragmentation: yet another barrier



If grains grow fast (e.g. high dust-to-gas ratio) —> they can overcome the radial drift barrier.

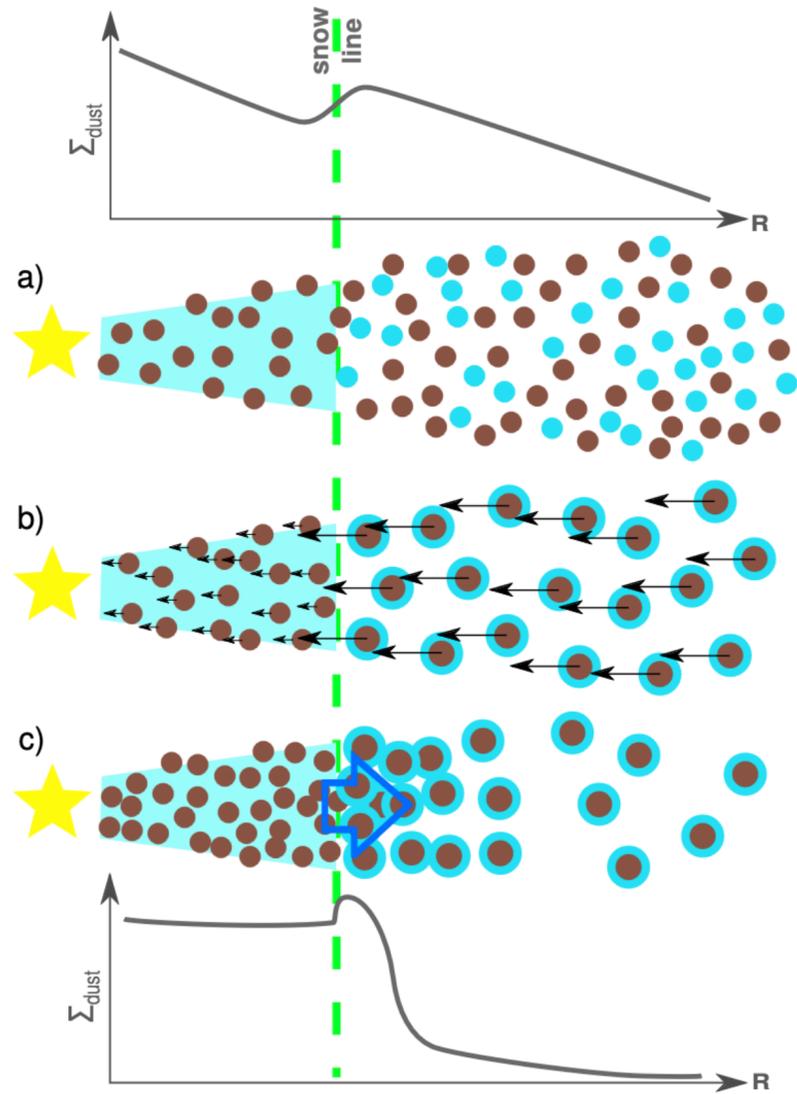
— —> Unfortunately grains could break below the barrier.

In that case, not only do we not have big grains but we also lose the small ones to radial drift... (fortunately we have seen some solutions)

*Brauer et al. (2008)*

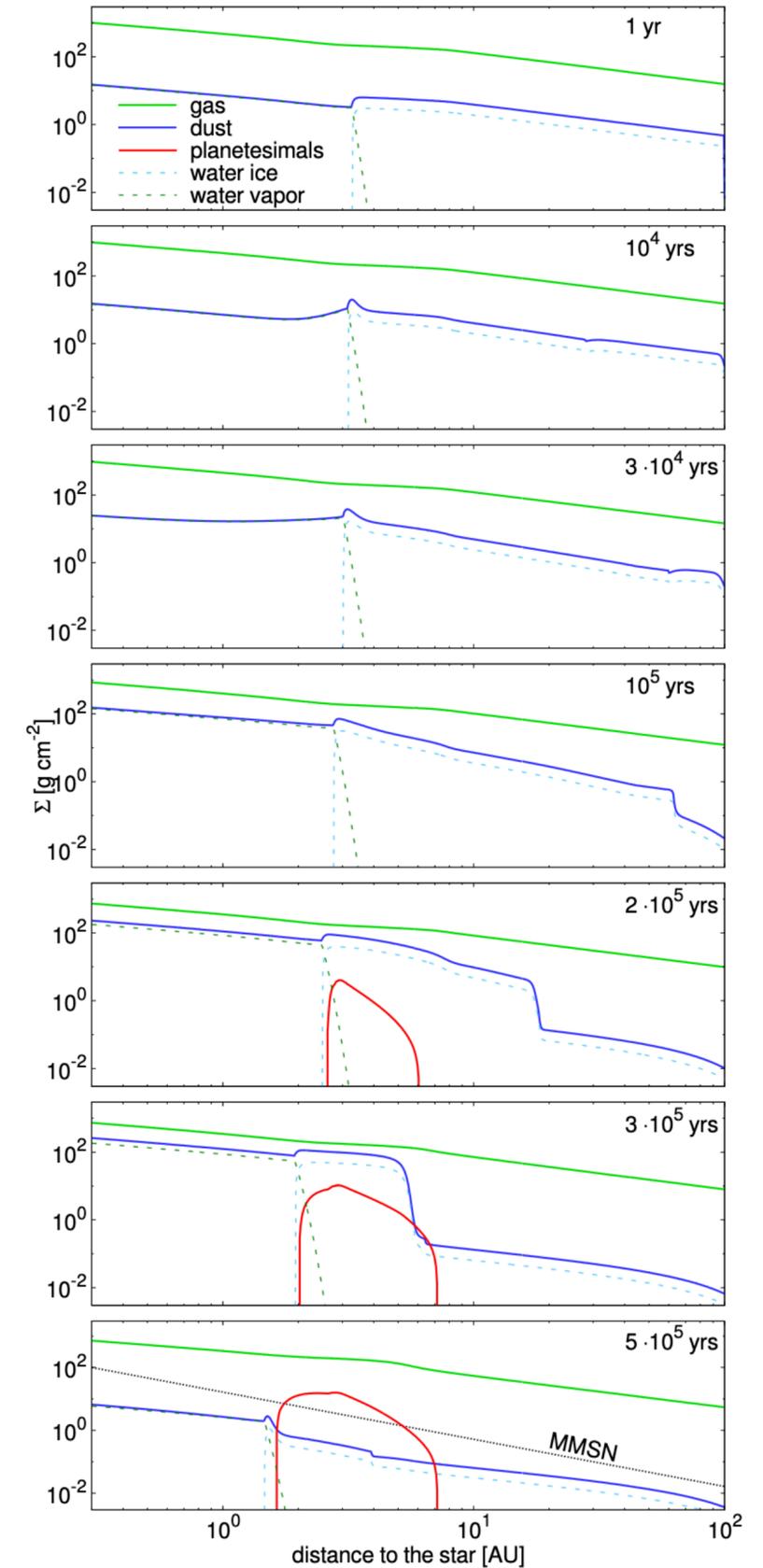
# Fragmentation: may it help ?

The change of fragmentation velocity at the snow line could trigger a traffic jam



→ increase in the dust to gas ratio

→ the streaming instability could be triggered at snow lines !



# Back to the protostellar collapse

# Fragmentation : where is the threshold ?

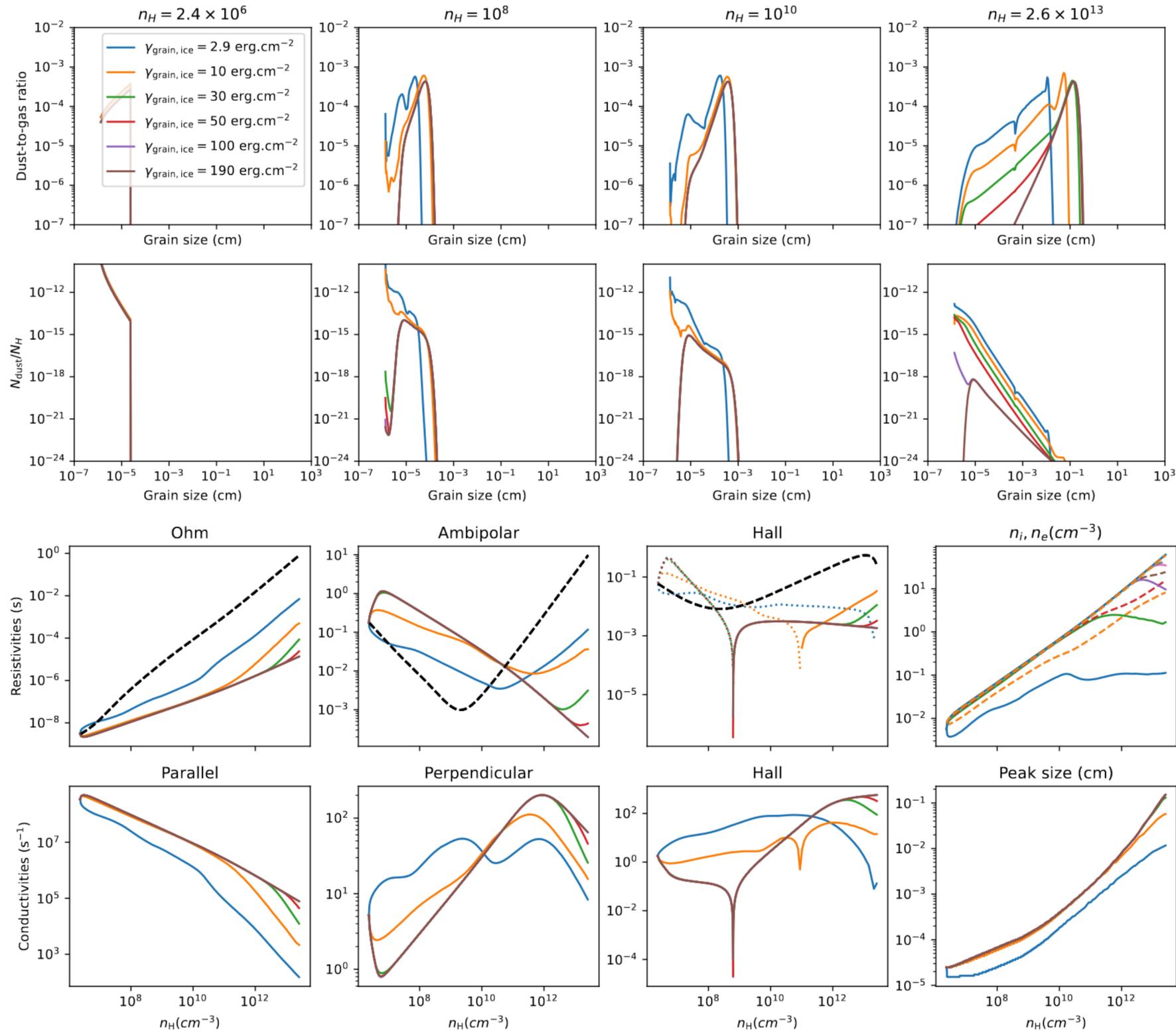
**We do not know well the elastic properties of the grains !**

The fragmentation velocity could range from as low as : 0.1m/s to 100 m/s (plus it's not supposed to be a velocity threshold)

**Big problem for planet formation because we don't know what Stokes is reached.**

Gundlach and Blum 2015; Poppe et al. 2000; Güttler et al. 2010; Musiolik et al. 2016a,b; Blum and Wurm 2008, *Musiolik & Wurm 2019, Yamamoto et al. (2014) etc..*

# Impact of coagulation/fragmentation on the magnetic resistivities



Vallucci-Goy et al. sub see also :  
 Guillet et al. 2020; Silsee et al. 2020, Lebreuilly et al. 2023

# Conclusions II

- Coagulations happens even during the protostellar collapse.
- It leads to the growth from ISM grains to grains up to mm/cm/dm? sizes
- Coagulation can either be modelled by the solving Smoluchowski equation or the Stekinski (monodisperse) equation
- Coagulation is hard to solve in 3D but new methods are on their way.
- Coagulation deeply affect the magnetic resistivities and therefore disk formation

